



FACULTAD DE MATEMÁTICAS
PONTIFICIA UNIVERSIDAD CATÓLICA DE CHILE

Pre-Publicación MATUC - 2018 - 6

**ON THE K-THEORETIC CLASSIFICATION OF DYNAMICALLY
STABLE SYSTEMS**

Giuseppe de Nittis¹, Kiyonory Gomi²

¹Facultad de Matemáticas, Pontificia Universidad Católica de Chile, Chile

²Department of Mathematical Sciences, Shinshu University, Nagano Japan

ON THE K-THEORETIC CLASSIFICATION OF DYNAMICALLY STABLE SYSTEMS

GIUSEPPE DE NITTIS AND KIYONORI GOMI

ABSTRACT. This paper deals with the construction of a suitable topological K -theory capable of classifying topological phases of dynamically stable systems described by gapped η -self-adjoint operators on a Krein space with indefinite metric η .

CONTENTS

1. Introduction	1
1.1. Motivations	2
1.2. Methodological extension of the CAZ classification scheme	4
1.3. Methodological extension of the K -theory classification scheme	7
1.4. Some open question	11
1.5. Structure of the paper	11
2. Elements of the Krein space theory	11
2.1. Indefinite metric spaces	11
2.2. Unitray vs. η -unitary operators	13
2.3. Pseudo-unitarity in indefinite metric spaces	15
3. Quantum symmetries: The Wigner's theorem	16
3.1. The Hilbert space case	17
3.2. The Krein space case	19
3.3. Symmetries of η -self-adjoint operators	20
4. \mathcal{C} -symmetry for η -self-adjoint operators	22
4.1. Definition and main properties	22
4.2. Operators with \mathcal{C} -symmetry	24
4.3. The Hilbert bundle picture	26
5. K -theory	30
5.1. Karoubi's formulation of Freed-Moore K -theory	30
5.2. Infinite-dimensional Karoubi's formulation	34
5.3. Fredholm operator formulation of Freed-Moore K -theory	35
5.4. K -theory for η -self adjoint operators	38
5.5. Reduction to the original Freed-Moore K -theory	39
5.6. Explicit computations in some special cases	41
Appendix A. Maxwell-type operators and metamaterials	43
Appendix B. Some explicit computation in dimension two	46
References	50

1. INTRODUCTION

This paper deals with the construction of a suitable K -theory capable of classifying topological phases of dynamically stable systems described by η -self-adjoint operators on a Krein space with indefinite metric (or fundamental symmetry) η . The motivations, the methodology and the main results are described below.

MSC2010 Primary: 16E20; Secondary: 81T99, 46L80, 14D21.

Keywords. K -theory, topological phases, Krein spaces, \mathcal{PT} and \mathcal{C} symmetries.

1.1. Motivations. In the standard formulation of Quantum Mechanics (QM) the physics of a (closed) system is described by a separable Hilbert space \mathcal{H} endowed with a *scalar product*¹ $\langle \cdot, \cdot \rangle$ and by a (time-independent) self-adjoint operator H , called Hamiltonian. The vectors in \mathcal{H} describe the possible *states* of the system, the scalar product $\langle \cdot, \cdot \rangle$ fixes the *probabilistic interpretation* of the theory and the Hamiltonian H generates the dynamics of the system through the Schrödinger equation. The role of the self-adjoint condition $H = H^*$ is twofold; (i) it assures the reality of the spectrum of H which is interpreted as the set of possible energies of the system; (ii) it assures that the dynamics generated by H is implemented by a unitary group $t \mapsto V_t := e^{-itH}$, called the *time propagator*, which preserves the geometry of \mathcal{H} and therefore the probabilistic interpretations. Bender and coworkers, in a long series of papers (see e. g. [BB, BBJ2, Ben1]), argued that conditions (i) and (ii) are the only crucial properties which characterize physical (quantum) systems which are *dynamically stable*. With this locution one refers to (somehow unphysical) systems which are not subjected to dissipations or losses (of matter, or energy, or information) and that are not destined to disappear in a remote future or past (cf. [PS]). Although the self-adjointness of the Hamiltonian is a sufficient condition for (i) and (ii) it is not strictly necessary (a fact already recognized at the early days of QM [Dir, Pau]) and can be relaxed somehow. A possibility to formulate rigorously the notion of dynamical stability is as follows:

Definition 1.1 (Dynamical stability). *A Hamiltonian H acting on the Hilbert space \mathcal{H} describes a dynamically stable system if there exists a bounded operator G , with bounded inverse G^{-1} , such that $\tilde{H} := GHG^{-1}$ is self-adjoint on \mathcal{H} .*

Said differently, H is dynamically stable whenever it is *similar* to a self-adjoint operator on \mathcal{H} . If this is the case condition (i) about the reality of the spectrum follows by the fact that similar operators have the same spectrum and so $\sigma(H) = \sigma(\tilde{H}) \subseteq \mathbb{R}$. The dynamical condition (ii) follows by the observation that the group $t \mapsto V_t := G^{-1}e^{-it\tilde{H}}G$ integrates the Schrödinger equation

$$i \frac{\partial}{\partial t} \psi = H\psi, \quad \psi \in \mathcal{D}(H).$$

and is unitary with respect to the modified scalar product $\langle \phi, \varphi \rangle_G := \langle \phi, G^*G\varphi \rangle$ for all $\phi, \varphi \in \mathcal{H}$. In particular, H turns out to be self-adjoint with respect to the scalar product $\langle \cdot, \cdot \rangle_Z$. For more details we refer to Corollary 4.9.

Definition 1.1 is quite strong and in principle it could be possible to set up weaker and workable reformulations of the notion of dynamical stability. However, Definition 1.1 has been inspired by physical and mathematical motivations that we will try to retrace below.

Bender and coworkers at the end of the 90's observed that certain non-self-adjoint Hamiltonians subjected to a fundamental space-time reflection symmetry (\mathcal{PT} -symmetry) seem to obey the two crucial requirements (i) and (ii) that characterize systems with a stable dynamics. This observation opened a new and fruitful line of research on \mathcal{PT} -symmetric Quantum Mechanics (\mathcal{PT} -QM). As the literature on \mathcal{PT} -QM is extremely large, and this is not the central focus of the present work, we refer to the recent review papers [Mos5, Ben1, Ben2, AK2], and references therein, for a complete overview on the subject. Nevertheless, it is worth to mention that intriguing toy examples of (one-dimensional) \mathcal{PT} -symmetric systems are provided by the family of operators

$$H_\epsilon := -\frac{d^2}{dx^2} - (if(x))^\epsilon, \quad \epsilon \in \mathbb{R} \quad (1.1)$$

where $f(-x) = -f(x)$ is an odd real-valued function. The special case $f(x) = x$ has been extensively studied in [Ben1] from a classical and a quantum point of view. A second family of examples is given by the Hamiltonians

$$m_w := w(x) \frac{d^2}{dx^2}, \quad (1.2)$$

¹In this paper the name *scalar product* is used as synonyms of *positive definite sesquilinear form*. This is the same thing as requiring that the associated quadratic form $\mathcal{Q}(\varphi) := \langle \varphi, \varphi \rangle > 0$ is strictly positive for non-zero vectors $\varphi \neq 0$.

with w a real-valued *weight* function such that $w(x)^2 = 1$. The special case $w(x) = \text{sgn}(x) := x/|x|$ has been studied for the first time in [CN] and recently in [KSu]. Moreover, models of \mathcal{PT} -QM have recently found application in a wide variety of areas like in *optics* [EMCM] (see also [Ben3] and references therein), in the study of the *atomic diffusion* in inhomogeneous magnetic fields [ZSW], in the theory of *superconducting wires* [RSM, CGBV] and *electronic circuits* [SLZEK], in *quantum information* [Cro], in the study of *nonlinear waves* [KYZ], *magnon* systems [LKS, GV] and *metamaterials* [Mos6] ... and probably others.

Just after the first works on \mathcal{PT} -QM various authors (like Mostafazadeh [Mos1, Mos2, Mos3] or Albeverio and Kuzhel [AK1, AK2] just to mention few of them) recognized that the concept of \mathcal{PT} -QM can be placed in a more general and natural mathematical framework by using the theory of *Krein spaces* [AI, Bog] (see also Section 2). In its simplest incarnation a Krein space can be thought as a separable Hilbert space \mathcal{H} along with a self-adjoint involution $\eta = \eta^* = \eta^{-1}$, called *fundamental symmetry* (or *metric* operator), such that the dimensions of the eigenspaces related to the eigenvalues ± 1 are equal. The fundamental symmetry η endows \mathcal{H} with the non-degenerate, indefinite inner product $\langle\langle \phi, \varphi \rangle\rangle_\eta := \langle \phi, \eta \varphi \rangle$ for all $\phi, \varphi \in \mathcal{H}$. The η -adjoint [Tom, Ghe] of a linear operator H acting on \mathcal{H} is then defined by

$$H^\sharp := \eta H^* \eta.$$

Operators which fulfill the conditions $H^\sharp = H$ are called *pseudo-Hermitian* [Mos1, Mos2, Mos4, Mos5, Jon, AK1, SHEK, Gho], and sometimes *quasi-Hermitia* [SGH], in the physical jargon. We prefer to refer to them as η -self-adjoint in agreement with a more mathematical tradition. It is not hard to see that both operators H_ϵ in (1.1) and m_w in (1.2) are η -self-adjoint with respect to a suitable choice of the fundamental symmetry η . For the operators H_ϵ the appropriate fundamental symmetry is given by $\eta \equiv \mathcal{P}$ where \mathcal{P} is the *parity* operator $(\mathcal{P}\phi)(x) := \phi(-x)$. In the case of the operator m_w the right fundamental symmetry is provided by $\eta \equiv w(x)$. Further examples of η -self-adjoint operators, adapted for certain physical applications, will be described in Appendix A.

In general, η -self-adjointness does not suffice to assure the dynamical stability of Definition 1.1. For that, something more is required. One says that a η -self-adjoint operator H admits a *C-symmetry* if there is an η -self-adjoint involution Ξ such that the product $\eta\Xi > 0$ is strictly positive and $H\Xi = \Xi H$ (see e. g. [Kuz, AK2, KSu] and references therein²). The theory of *C-symmetric* η -self-adjoint operators will be briefly reviewed in Section 4. Here, we need only to recall a fundamental fact [AK2, Theorem 6.3.4]: The existence of a *C-symmetry* for the η -self-adjoint operator H is equivalent to the dynamical stability according to Definition 1.1 (see Section 4.2 for more details). Consequently, on the basis of the considerations above, we feel justified to make the following:

Working hypothesis. *η -self-adjoint operators that possess a C-symmetry are good (and interesting) candidates for the description of the dynamics of certain physical (quantum) systems.*

Before continuing, it is fare to mention that to prove that a given η -self-adjoint operator admits a *C-symmetry* is not a trivial fact in general. Various conditions for the existence and the construction of a *C-symmetry* have been studied in [Kuz, AK2, KSu]. For instance, it turns out that the operator m_w of (1.2) for the case $w(x) = x/|x|$ has a *C-symmetry* described in [KSu, Proposition 3.5]. On the other hand the construction of a *C-symmetry* for the family of operators H_ϵ in (1.1) is a complicated (still open) problem. For instance, for the special case $f(x) = x$, it has been shown that $\sigma(H_\epsilon) \subseteq [0, +\infty)$ for $2 \leq \epsilon < 4$ [DDT, Appendix B] but the form of the corresponding operator implementing the *C-symmetry* (if any) is not known. Because of the complexity of the problem the majority of the available techniques for the construction of *C-symmetries* are based on perturbative approximations [Ben1].

In the last decade since the groundbreaking discovery of topologically protected states induced by strong spin-orbit interactions, tremendous progress has been made in the study and the understanding of topological properties of (quantum) matter. The investigation of *topological insulators* (we refer to

²In the standard literature a *C-symmetry* is usually denoted with the letter *C*. We prefer to use the greek letter Ξ to avoid confusion with the complex conjugation denote here with a *C*.

[HK2, AF, CTSR] for a modern overview and an updated bibliography of the subject) is doubtless a current “hot topic” in mathematical physics and condensed matter. This evidence justifies the following natural question: *Is it possible to generalize the set of results and theories available for the study and the classification of “standard” topological insulators, described by self-adjoint operators, to more general systems which are dynamically stable in the sense of Definition 1.1?* Of course the idea of investigate the topological properties of specific non-self-adjoint systems is not new in literature. The first result in this direction established the absence of topological phases for certain \mathcal{PT} -symmetric Dirac Hamiltonians [HH]. Soon after, the existence of non-trivial topological phases for non-self-adjoint generalizations of the Kane-Mele model has been proved in [ESHK]. However, these phases possess the unpleasant feature of having complex energies and this is an evidence of the fact that the corresponding Hamiltonian cannot be dynamically stable. Other examples of non-self-adjoint models with *unstable* (i. e. complex energy) topological states have been proposed in [Sch1] as a model for one-dimensional photonic crystals and in [ZLC] via a non-self-adjoint modification of the Su-Schrieffer-Heeger model. The first examples of non-self-adjoint models with *stable* (i. e. real energy) topological states have been builded in [Gho] by using Dirac-type Hamiltonians and in [Yuc] by a non-Hermitian generalization of the Aubry-Andre model. A new non-self-adjoint modification of the Su-Schrieffer-Heeger model which shows *stable* topological states has been recently proposed in [Lie]. Finally, the study of *topological edge modes* in \mathcal{PT} -symmetric models has gained interest in various physical contexts [LBHCN, KWLWL, PS], just to mention some of the most recent contributions.

An aspect common to all the works mentioned above is that they concern with very specific models. In contrast, one of the bigger success, perhaps the greatest, in the study of “standard” topological insulators was the discovery of a unified framework where to organize and classify all the possible topological phases according to the spatial dimensions. This is the so-called *periodic table* for topological insulators. After some initial partial results toward a unified classification of topological insulators [AZ, SRFL, QHZ], the mechanism underlying the periodic organization of the topological phases was finally revealed by Kitaev in the seminal paper [Kit]. Here, Kitaev showed that topological insulators can be organized and classified by using the *K-theory* [Ati2, Kar] and the periodic recurrence of the topological phases with respect to the spatial dimension was interpreted as a manifestation of the *Bott periodicity*. Starting from this moment, *K-Theory* became the principal tool in the classification of topological insulators described by self-adjoint models and the initial Kitaev’s idea has been generalized and extended in several directions [FM, Thi, Ke, Kub] producing a great advancement in the understanding of topological phases of matter. With these premises it seems natural to believe that there must be a more general classification scheme capable of including also non-self-adjoint dynamically stable systems. The present work has been inspired by this idea and its main achievement is:

Main goal. *To construct a suitable K-theory capable of classifying topological phases of systems described by η -self-adjoint operators that possess a C-symmetry.*

At a first glance, based on [AK2, Theorem 6.3.4] which proves the equivalence between η -self-adjoint operators with a C-symmetry and dynamically stable systems in the sense of Definition 1.1, one can naively think that the usual *K-theory* used for the classification of self-adjoint systems should be enough also to classify the type of systems described by the working hypothesis above. Indeed, this is not the case! In fact, as discussed below, systems described by C-symmetric η -self-adjoint operators posses more structure than simple self-adjoint systems. A very elementary example that shows this enrichment is discussed in Remark B.7.

1.2. Methodological extension of the CAZ classification scheme. In the classification scheme of “standard” topological insulators [AZ, SRFL, QHZ] enter two crucial ingredients: The spatial dimension of the system and a bunch of fundamental *involutive (pseudo-)symmetries*. Let us focus on the role of the latter.

The nature of the *quantum symmetries* has been described by Wigner in the seminal paper [Wig] (see Section 3.1 for more details). The *Wigner’s theorem* states that the physical symmetries of a

quantum systems represented in a Hilbert space \mathcal{H} can be implemented only by unitary or anti-unitary operators³. The set of all quantum symmetries on \mathcal{H} has a group structure and is denoted with $\text{QS}(\mathcal{H})$. The following notion is central in the classification of topological insulators.

Definition 1.2 (Involutive quantum symmetry). *An involutive quantum symmetry on the Hilbert space \mathcal{H} is a bounded \mathbb{R} -linear map $U : \mathcal{H} \rightarrow \mathcal{H}$ which meets:*

- (i) $U^{-1} = U^*$,
- (ii) $U(i\mathbb{1}) = \varpi i U$,
- (iii) $U = \varepsilon^{\frac{1-\varpi}{2}} U^{-1}$,

with $\varpi, \varepsilon \in \{-1, +1\}$.

Properties (i) and (ii) in Definition 1.2 are simply a way of expressing the fact that $U \in \text{QS}(\mathcal{H})$, namely that U is a unitary ($\varpi = +1$) or anti-unitary ($\varpi = -1$) operator. It is fair to observe that the specification about the boundedness of U at the beginning of Definition 1.2 is unnecessarily since it is implied by property (i). However, this is a point of distinction with respect to the Krein space case and for this reason we prefer to insist on it. Item (iii) deserves some clarifications. Usually a symmetry U is called *involutive* if $U^2 = \varepsilon\mathbb{1}$ where $\varepsilon = \pm 1$ and is a *proper* involution if $\varepsilon = +1$ and an *anti*-involution when $\varepsilon = -1$. A simple calculation shows that $(iU)^2 = i^{1+\varpi}\varepsilon\mathbb{1}$. Therefore, the transformation $U \mapsto iU$ induces the changes $\varepsilon \mapsto -\varepsilon$ in the linear case while it leaves unchanged the values of ε in the anti-linear case. This observation suggests that in the linear case an anti-involution can be always turned into a proper involution while the difference between anti-involution and proper involution is persistent in the anti-linear case. As a consequence one can use the convention $U^2 = \varepsilon^{\frac{1-\varpi}{2}}\mathbb{1}$ for the signs of the involutive quantum symmetry. The latter relation justifies property (iii). Summing up, in the linear case ($\varpi = +1$) an involutive quantum symmetry is always a self-adjoint unitary $U^{-1} = U^* = U$, independently of the second parameter ε . Conversely, in the anti-linear case ($\varpi = -1$) one has *even* symmetries ($\varepsilon = +1$) which are still specified by $U^{-1} = U^* = U$ and *odd* symmetries ($\varepsilon = -1$) which are skew-adjoint unitaries $U^{-1} = U^* = -U$.

Symmetry	Symbol	ϖ	ε	c
Proper Linear		+1	irr.	+1
Chiral	χ	+1	irr.	-1
Even Time Reversal	$T (+)$	-1	+1	+1
Odd Time Reversal	$T (-)$	-1	-1	+1
Even Particle-Hole	$P (+)$	-1	+1	-1
Odd Particle-Hole	$P (-)$	-1	-1	-1

TABLE 1.1. The six types of involutive quantum symmetries for a given self-adjoint operator $H = H^*$. In the case of the first two linear symmetries the sign ε is irrelevant (irr.) as suggested by property (iii) in Definition (1.2). The various names come from canonical examples of fundamental symmetries for electronic systems [AZ, SRFL, QHZ].

³An operator A on the Hilbert space \mathcal{H} is anti-linear when $A(a\phi + b\psi) = \overline{a}A\phi + \overline{b}A\psi$ for all $a, b \in \mathbb{C}$ and $\phi, \psi \in \mathcal{D}(A)$. For an anti-linear operator the definition of the adjoint needs to be adjusted in order to compensate for the complex conjugation. The adjoint operator of the anti-linear operator A is the anti-linear operator A^* which fulfills $\langle \phi, A\psi \rangle = \langle \psi, A^*\phi \rangle$ for all $\phi \in \mathcal{D}(A)$ and $\psi \in \mathcal{D}(A^*)$.

The interplay between quantum symmetries and a self-adjoint Hamiltonian $H = H^*$ can be expressed by the relation

$$UH = c HU, \quad c = \pm 1 \quad (\text{commuting vs. anti-commuting dichotomy}).$$

We say that a $U \in \text{QS}(\mathcal{H})$ is a *proper* ($c = +1$) quantum symmetry for H if $UH = HU$ and is a *improper* ($c = -1$) quantum symmetry if $UH = -HU$.

According to the different signs of the parameters ϖ, ε and c the involutive quantum symmetries are classified in six families as listed in Table 1.1. The latter are the building blocks for the classification of “standard” topological insulators. In the *Cartan-Altlund-Zirnbauer (CAZ) scheme* [AZ, SRFL, QHZ] gapped electronic systems are arranged in ten classes according to the fact that the (self-adjoint) Hamiltonian H that describes the system possesses or violates the *chiral* symmetry χ , the *time reversal* symmetry T or the *particle-hole* symmetry P . Linear proper symmetries do not enter in the topological classification since the commutation relation $[H, U] = 0$ assures that H can be decomposed along the eigenspaces of U providing reduced systems which can be independently classified. The ten classes of topological insulators in the CAZ classification are described in Table 1.2.

	A	AIII	AI	AII	D	C	BDI	CI	DIII	CII
χ	0	1	0	0	0	0	(1)	(1)	(1)	(1)
T	0	0	+1	-1	0	0	+1	+1	-1	-1
P	0	0	0	0	+1	-1	+1	-1	+1	-1

TABLE 1.2. The ten classes of the *Cartan-Altlund-Zirnbauer (CAZ) scheme* [AZ, SRFL, QHZ] for the topological classification of self-adjoint Hamiltonians. For the case of the chiral symmetry χ the 1 means that the symmetry is present and the 0 that the symmetry is broken. A similar convention is used for the time reversal symmetry T and the particle-hole symmetry P with the only difference that +1 or -1 means that the related symmetry is even or odd, respectively. The first two classes A and AIII share the characteristic of not having anti-linear symmetries and form the subfamily of the *complex* classes. The last eight classes have at least one anti-linear symmetry and provide the subfamily of the *real* classes. The last four classes possess simultaneously the two anti-linear symmetries T and P . In this case the commutation condition $PT = TP$ is tacitly assumed. In particular, the product PT gives rise to a linear symmetry of type χ and this fact is annotated in the table by the symbol (1).

The natural question we want to raise and analyze in this work is *how the CAZ classification scheme change in the case of η -self-adjoint operators that possess a C -symmetry*. The starting point of the analysis concerns the fate of the Wigner’s theorem in Krein spaces (see Section 3.2). This question has been investigated by Bracci, Morchio and Strocchi in [BMS] where they proved an extended version of the Wigner’s theorem adapted to spaces with an indefinite metric. The first important difference with respect to the usual Hilbert space case is that quantum symmetries in a Krein space can be implemented also by unbounded operators (see Remark 2.7). Moreover, this eventuality survives even if one focuses the attention to involutive symmetries (see Remark 4.3). For this reason, if one decides to work only with bounded quantum symmetries, such a choice has to be specified. Let $\text{QS}_\eta(\mathcal{H})$ be the set of the bounded) η -quantum symmetries acting in the Krein space \mathcal{H} with metric η . The second difference, is that η -quantum symmetries are subjected to two types of dichotomies: The usual linear vs. anti-linear dichotomy expressed by property (ii) in Definition 1.2 and a new dichotomy, absent in the Hilbert space case, described by the equation $U^{-1} = \wp U^\sharp = \wp \eta U^* \eta$ where \wp is a sign which distinguishes between η -isometries ($\wp = +1$) and η -pseudo-isometries ($\wp = -1$). The combinations of the different values of ϖ and \wp lead to four types of η -quantum symmetries as listed in Table 1.3.

By combining the η -isometric condition with the involutive requirement $U = \varepsilon U^{-1}$ one obtains $U = \wp \varepsilon U^\sharp$. Again the change $U \mapsto i U$ can be used to change the sign of ε in the linear case but it is

$U \in \text{QS}_\eta(\mathcal{H})$	η -unitary	η -anti-unitary	η -pseudo-unitary	η -pseudo-anti-unitary
\wp	+1	+1	-1	-1
ϖ	+1	-1	+1	-1

TABLE 1.3. The four types of η -quantum symmetries according to the values of ϖ and \wp .

important to notice that it preserves the sign of \wp independently of the linearity or anti-linearity of U . The last observations justify the following:

Definition 1.3 (Involutive η -quantum symmetry). *An involutive η -quantum symmetry on the Krein space \mathcal{H} with metric η is a bounded \mathbb{R} -linear map $U : \mathcal{H} \rightarrow \mathcal{H}$ which meets:*

- (i) $U^{-1} = \wp U^\sharp$,
- (ii) $U(i\mathbb{1}) = \varpi i U$,
- (iii) $U = \varepsilon^{\frac{1-\varpi}{2}} U^{-1}$,

with $\wp, \varpi, \varepsilon \in \{-1, +1\}$.

A comparison between Definition 1.2 and Definition 1.3 shows that the main difference is in the condition (i) where the extra sign \wp appears. This sign is responsible for enriching the zoology of the possible symmetries. Again, the dichotomy even vs. odd involution is relevant only in the anti-linear case while in the linear case one conventionally assumes that involutions are always of even type.

In Krein spaces the “physical” Hamiltonians are η -self adjoint $H = H^\sharp$. The interplay between Hamiltonians and η -quantum symmetries is again expressed by $UH = c HU$ where $c = \pm 1$ discriminates between proper (commuting) and improper (anti-commuting) symmetries.

According to the different signs of the parameters $\varpi, \wp, \varepsilon, c$ the involutive η -quantum symmetries are grouped in 12 families which are described in Table 1.4. A comparison between Table 1.1 and Table 1.4 shows that in the Krein space case the number of possible symmetries doubles. Along with the classical symmetries χ, T, P one has the new η -reflecting symmetry R and the reflecting versions χ_R, T_R, P_R of the classical symmetries. The set of these seven symmetries, along with the dichotomy even vs. odd in the anti-linear cases, are the building blocks for the construction of the extended classification scheme for η -self-adjoint operators. Again linear proper symmetries do not enter in the classification scheme since the commutation relations $[H, U] = 0$ and $[\eta, U] = 0$ assure that H and η can be decomposed along the eigenspaces of U providing reduced systems which can be independently classified.

More symmetries imply more symmetry classes. In the Krein space case the number of possible CAZ classes increases considerably. For instance the two complex classes A and AII described in the first two columns of Table 1.2 correspond to five classes in the Krein space case as showed in Table 1.5. The two time-reversal classes AI and AII described in the second and third columns of Table 1.2 split in six new classes in the Krein space case as displayed in Table 1.6. A similar counting exercise can also be performed for the remaining classes. As a result the total number of (possible) symmetry classes for η -self-adjoint operators is much bigger than ten (a complete counting is beyond the scope of this work). Therefore, the so called *ten-fold way* of Table 1.2 which organizes the symmetry classes of standard self-adjoint systems appears to be insufficient in the case of dynamically stable systems described by η -self-adjoint operators with a C -symmetry.

1.3. Methodological extension of the K -theory classification scheme. Each symmetry class in the CAZ scheme of Table 1.2 supports a certain number of possible topological phases. The enumeration of these phases can be performed by different incarnations of the K -theory, according to the specificities and the generality of the situation under consideration [FM, Thi, Ke, Kub]. In the case of the

Symmetry	Symbol	ϖ	\wp	ε	c
Proper Linear		+1	+1	irr.	+1
Pure Reflecting	R	+1	-1	irr.	+1
Chiral	χ	+1	+1	irr.	-1
Reflecting Chiral	χ_R	+1	-1	irr.	-1
Even Time Reversal	$T (+)$	-1	+1	+1	+1
Even Reflecting Time Reversal	$T_R (+)$	-1	-1	+1	+1
Odd Time Reversal	$T (-)$	-1	+1	-1	+1
Odd Reflecting Time Reversal	$T_R (-)$	-1	-1	-1	+1
Even Particle-Hole	$P (+)$	-1	+1	+1	-1
Even Reflecting Particle-Hole	$P_R (+)$	-1	-1	+1	-1
Odd Particle-Hole	$P (-)$	-1	+1	-1	-1
Odd Reflecting Particle-Hole	$P_R (-)$	-1	-1	-1	-1

TABLE 1.4. The 12 types of strong involutive η -quantum symmetries for a given η -self-adjoint operator $H = H^\sharp$. In the case of the first four linear symmetries the sign ε is irrelevant (irr.) as suggested by property (iii) in Definition 1.3. The various names are chosen in such a way to keep the connection with the standard nomenclature of Table 1.1.

	R	χ	χ_R	T	T_R	P	P_R
A	0	0	0	0	0	0	0
R	1	0	0	0	0	0	0
AIII	0	1	0	0	0	0	0
AIII_R	0	0	1	0	0	0	0
RIII	1	1	(1)	0	0	0	0

TABLE 1.5. The five classes which extend the two complex classes (A and AIII) of the standard CAZ scheme of Table 1.2 in the case of η -self-adjoint Hamiltonians. These are the classes defined by *linear* (strong) involutive η -quantum symmetries. As usual 1 means that the symmetry is present and 0 that the symmetry is broken. In the case of the class RIII the commutation condition $R\chi = \chi R$ is tacitly assumed.

description of systems described by self-adjoint gapped Hamiltonians the relevant K -theory has been defined and studied by Freed and Moore [FM].

	R	χ	χ_R	T	T_R	P	P_R
AI	0	0	0	+1	0	0	0
AI_R	0	0	0	0	+1	0	0
RI	1	0	0	+1	(+1)	0	0
AII	0	0	0	-1	0	0	0
AII_R	0	0	0	0	-1	0	0
RII	1	0	0	-1	(-1)	0	0

TABLE 1.6. The six classes which extend the two time reversal classes (AI and AII) of the standard CAZ scheme of Table 1.2 for the case of η -self-adjoint Hamiltonians. These are the classes defined by (strong) involutive η -quantum symmetries with $c = +1$. As usual 1 means that the symmetry is present and 0 that the symmetry is broken. In the case of the anti-linear symmetries T and T_R the signs +1 or -1 means that the related symmetry is even or odd, respectively. In the case of the classes RI and RII the commutation condition $RT = TR$ is tacitly assumed.

There are two key points in the construction of the K -theory classification proposed by Freed and Moore: One key point is that the self-adjoint Hamiltonian, describing a gapped quantum systems, plays the role of a *gradation* (namely a self-adjoint involutions) in the sense of the Karoubi's construction of K -theory [Kar]. The second key point is that one can incorporate various quantum symmetries of the systems into a modified version of the Karoubi's K -theory. For example, the incorporation of the time reversal symmetry T leads to the Atiyah's KR -theory [Ati1] in the even case and to the Dupont's KQ -theory [Du] in the odd case. The *Freed-Moore K-theory* [FM] unifies and generalizes the KR -theory and the KQ -theory to an arbitrary number of quantum symmetries.

Let us be a bit more precise on the description of the crucial ideas for the construction of the K -theory of Freed-Moore. First of all the gradation. Given a self-adjoint Hamiltonian H with a *gap* around the energy $\lambda \notin \sigma(H)$ one can construct via spectral calculus the operator

$$\Gamma_H := \text{sgn}(H) = \frac{H - \lambda \mathbb{1}}{|H - \lambda \mathbb{1}|} = \frac{H - \lambda \mathbb{1}}{\sqrt{(H - \lambda \mathbb{1})^2}}. \quad (1.3)$$

By construction Γ_H is a gradation of \mathcal{H} since $\Gamma_H = \Gamma_H^* = \Gamma_H^{-1}$ (cf. Remark 5.6). Now the symmetries. Let $T, P \in \text{QS}(\mathcal{H})$ be respectively a time reversal symmetry and a particle-hole symmetry for H , and hence for Γ_H , in the sense of Table 1.1. The operator T and P are anti-linear and by general arguments it is not necessary to consider other anti-linear symmetries [Thi, Lemma 6.1]. Moreover, one can also assume the commutativity of T and P [Thi, Proposition 6.2]. This fact allows to define also the chiral symmetry $\chi := TP = PT$. By following [Thi, Section 6.1] let us call PT -group the subgroup of $\text{QS}(\mathcal{H})$ given by $\{\mathbb{1}, T, P, \chi\} \simeq \mathbb{Z}_2 \times \mathbb{Z}_2$ where the isomorphism is generated by the signs ϖ and c . This subgroup implements different representations of appropriate Clifford algebras $C\ell^{r,s}$ in function of the signs of T^2 and P^2 . Moreover, these representations can be graded or ungraded with respect to the gradation Γ_H . The possible Clifford actions representable by the PT -group are ten and are in one-to-one correspondence with the ten classes of the CAZ scheme described in Table 1.2 (for the details see [Thi, Section 6.1]). A K -theory capable of classifying at the same time all the ten CAZ classes must be endowed with a graded structure coming from the gapped Hamiltonian H and must be compatible with (graded or ungraded) Clifford actions implemented by unitary or anti-unitary operators associated with the PT -group (the so called *PUA-representations* [Par, Thi]). This is exactly what the K -theory of

Freed and Moore does [FM]. A more precise presentation of the Freed-Moore K -theory is postponed in Section 5.1.

How the picture sketched above changes in presence of a Krein space structure given by an indefinite metric η ? In this case the role of the gapped self-adjoint operators is played by gapped η -self-adjoint operators with a C -symmetry Ξ . However, in order to extend in a straightforward way the K -theory picture sketched above we will make use of the following representation:

Theorem 1.4. *Let $(\mathcal{H}, \langle \cdot, \cdot \rangle_\eta)$ be the Krein space associated to a fundamental symmetry $\eta^* = \eta = \eta^{-1}$. Let H be a η -self-adjoint operator which admits a C -symmetry Ξ and $U \in \text{QS}_\eta(\mathcal{H})$ a bounded η -quantum symmetry of type (ϖ, \wp) . Then:*

- (1) *There is a bounded η -unitary operator $G_\Xi \in \text{QS}_\eta(\mathcal{H})$ such that the transformed operator*

$$\tilde{H} := G_\Xi H G_\Xi^{-1}$$

is self-adjoint and η -self adjoint, i. e. $\tilde{H}^ = \tilde{H} = \tilde{H}^\sharp$.*

- (2) *The transformed symmetry*

$$\tilde{U} := G_\Xi U G_\Xi^{-1} \in \text{QS}_\eta(\mathcal{H})$$

is still of type (ϖ, \wp) . Moreover

$$UH = cHU \quad \Leftrightarrow \quad \tilde{U}\tilde{H} = c\tilde{H}\tilde{U}. \quad (1.4)$$

- (3) *Assume in addition that U is an involutive η -quantum symmetry (in the sense of Definition 1.3) characterized by the signs $(\varpi, \wp, \varepsilon)$ which Ξ -compatible in the sense that*

$$U\Xi = \wp \Xi U. \quad (1.5)$$

Then \tilde{U} is an involutive η -quantum symmetry of type $(\varpi, \wp, \varepsilon)$ which fulfills

$$\tilde{U}^{-1} = \tilde{U}^*, \quad \tilde{U}\eta = \wp \eta \tilde{U}.$$

In particular \tilde{U} is a unitary ($\wp = +1$) or anti-unitary ($\wp = -1$) operator on \mathcal{H} .

Proof (sketch of). Item (1) is proved in Proposition 4.7 and implies that \tilde{H} is a self-adjoint operator such that $\tilde{H}\eta = \eta\tilde{H}$. Item (2) is quite immediate. In fact U and G_Ξ are both elements of the group $\text{QS}_\eta(\mathcal{H})$ and so also their composition \tilde{U} belongs to $\text{QS}_\eta(\mathcal{H})$. Since G_Ξ is linear, it follows that the linear or anti-linear nature of \tilde{U} depends only on U . Similar argument for the sign \wp . A direct computation provides $\tilde{U}^\sharp = G_\Xi U^\sharp G_\Xi^{-1}$ and so $\tilde{U}^\sharp = \wp \tilde{U}^{-1}$ if and only if $U^\sharp = \wp U^{-1}$. Finally equation (1.4) is a direct consequence of the definition of \tilde{U} and \tilde{H} . A general version of item (3) is proved in Proposition 4.11. Anyway, in view of item (2) one has that \tilde{U} is an η -quantum symmetry of type (ϖ, \wp) . Moreover a direct computation $\tilde{U}^2 = G_\Xi U^2 G_\Xi^{-1} = \varepsilon \mathbb{1}$ shows that \tilde{U} is an involution of parity ε . Finally let us observe that $\tilde{U}^{-1} = G_\Xi U^{-1} G_\Xi^{-1}$ and $\tilde{U}^* = G_\Xi^{-1} U^* G_\Xi$ where the condition $G_\Xi = G_\Xi^*$ has been used. Then the condition $\tilde{U}^{-1} = \tilde{U}^*$ is equivalent to $U^* G_\Xi^2 = G_\Xi^2 U^{-1}$. Since $G_\Xi^2 = \eta \Xi$ (cf. Lemma 4.4) the last equality is guaranteed by $G_\Xi^2 U^{-1} = \wp \eta U^{-1} \Xi$ which follows from (1.5) along with $U^* G_\Xi^2 = \eta U^\sharp \Xi = \wp \eta U^{-1} \Xi$ which is ensured by Definition 1.3. The condition $\tilde{U}^{-1} = \wp \eta \tilde{U}^* \eta$ along with the unitarity $\tilde{U}^{-1} = \tilde{U}^*$ implies the last relation $\tilde{U}\eta = \wp \eta \tilde{U}$. ■

The representation Theorem 1.4 allows to reconsider the problem of the construction of a K -theory for η -self-adjoint operators with a C -symmetry in the following terms: First of all, instead of *single* gapped Hamiltonians, we need to consider *pairs* (Γ_H, η) given by a gradation Γ_H (associated with a gapped Hamiltonian H according to (1.3)) and an indefinite metric η constrained by the condition $\Gamma_H \eta = \eta \Gamma_H$. Secondly, the fundamental symmetries are described by the *PTR-group* which is the subgroup of $\text{QS}_\eta(\mathcal{H}) \cap \text{QS}(\mathcal{H})$ given by $\{\mathbb{1}, T, P, \chi, R, T_R, P_R, \chi_R\} \simeq \mathbb{Z}_2 \times \mathbb{Z}_2 \times \mathbb{Z}_2$ where the isomorphism is generated by the signs ϖ, \wp and c . The *new* symmetries specified by $\wp = -1$ can be defined in terms of the fundamental reflecting symmetry R , the time reversal symmetry T and a particle-hole symmetry P according to $T_R := RT = TR$, $P_R := PR = RP$ and $\chi_R = R\chi = \chi R$. In this framework the extension of the Freed-Moore K -theory is straightforward, but not trivial. The extended K -theory

must be capable of classifying pairs (Γ_H, η) and must be compatible with (graded or ungraded) Clifford actions implemented by unitary or anti-unitary operators associated with the PTR -group. The details of this construction, which is the main goal of this work, are described in Section 5.4.

1.4. Some open question. This work contains a preliminary and foundational investigation around the problem of the K -theory classification of topological phases for η -self-adjoint operators which are dynamically stable. In the formulation developed here some strong assumptions have been formulated and there are several open questions which deserve future investigation. Among the latter, let us point out on the following questions: *What is the fate of the ten-fold way in the extended framework presented in this work? It is possible to consider the problem of the classification of topological phases induced by involutive η -quantum symmetries which are not necessarily bounded?* One last open question is presented at the end of Appendix A and concerns with the proof that certain η -self-adjoint operators admit a C -symmetry.

1.5. Structure of the paper. The paper is organized as follows: In **Section 2** the basic facts of the theory of Krein spaces are reviewed. In particular the notion of η -self-adjointness and the different types of η -isometries are presented. **Section 3** is devoted to the presentation of the extended version of the Wigner's theorem on Krein spaces. In **Section 4** the notion of C -symmetry is introduced and the most relevant properties of C -symmetric η -self-adjoint operators are proved. In **Section 5** the Freed-Moore K -theory is firstly reformulated in terms of Karoubi's gradations and then adapted for the classification of C -symmetric η -self-adjoint operators. In **Appendix A** a class of C -symmetric η -self-adjoint operators, relevant for physical applications, is described. **Appendix B** is devoted to the explicit classification of C -symmetric η -self-adjoint matrix on \mathbb{C}^2 .

Acknowledgements. GD's research is supported by the grant *Iniciación en Investigación 2015* - N° 11150143 funded by FONDECYT. KG's research is supported by the JSPS KAKENHI Grant Number 15K04871.

2. ELEMENTS OF THE KREIN SPACE THEORY

From a mathematical point of view the study of indefinite metric spaces in finite dimension started by G. Frobenius while the interest for the infinite dimension case has been pioneered by the russian mathematicians [Pon, IK1, IK2, GI, KL, Kre]. We refer to their papers (cited before), as well as to the two general monographs [AI, Bog], for a general exposition of the results obtained in that field.

2.1. Indefinite metric spaces. In this paper, the pair $(\mathcal{H}, \langle \cdot, \cdot \rangle)$ is always assumed to describe a *separable complex* Hilbert space \mathcal{H} with a *scalar product* $\langle \cdot, \cdot \rangle$ (cf. Note 1). We will use the symbol $\mathcal{B}(\mathcal{H})$ to denote the set of bounded operators on \mathcal{H} , namely the set of the linear or *anti*-linear maps $A : \mathcal{H} \rightarrow \mathcal{H}$ (cf. Note 3) such that

$$\|A\| := \sup_{\psi \in \mathcal{H} \setminus \{0\}} \left(\frac{\langle A\psi, A\psi \rangle}{\langle \psi, \psi \rangle} \right)^{\frac{1}{2}} < \infty. \quad (2.1)$$

We point out that the definition of boundedness is not affected by the dichotomy between linearity and anti-linearity. The subset of the *linear* bounded operators will be denoted by the symbol $\mathcal{B}_{\text{lin}}(\mathcal{H})$.

The notion of indefinite metric is introduced in the following way:

Definition 2.1 (Indefinite metric space). *An indefinite metric space $(\mathcal{H}, \langle \cdot, \cdot \rangle_\eta)$ is the datum of:*

- (i) *An underlying Hilbert space $(\mathcal{H}, \langle \cdot, \cdot \rangle)$.*
- (ii) *A non-degenerate, possibly indefinite, inner product $\langle \cdot, \cdot \rangle_\eta$ associated by the prescription*

$$\langle \langle \phi, \varphi \rangle \rangle_\eta := \langle \phi, \eta \varphi \rangle, \quad \forall \phi, \varphi \in \mathcal{H}$$

to a bounded linear operator $\eta \in \mathcal{B}_{\text{lin}}(\mathcal{H})$ which is self-adjoint $\eta = \eta^$, and which admits a bounded inverse $\eta^{-1} \in \mathcal{B}_{\text{lin}}(\mathcal{H})$.*

The self-adjoint condition for η assures that the quadratic form

$$\mathcal{Q}_\eta(\varphi) := \langle\langle \varphi, \varphi \rangle\rangle_\eta = \langle \varphi, \eta \varphi \rangle, \quad \varphi \in \mathcal{H}$$

is real-valued. The invertibility of η , which is equivalent to the fact that $\eta\varphi = 0$ if and only if $\varphi = 0$, assures that $\langle\langle \cdot, \cdot \rangle\rangle_\eta$ is non-degenerate.

Remark 2.2 (Fundamental symmetry). Without loss of generality an indefinite metric space can be assumed to have a *metric operator* which meets the conditions

$$\eta^* = \eta = \eta^{-1}. \quad (2.2)$$

To see this, let η' be a generic metric operator and build via functional calculus the operators $w := \sqrt{|\eta'|}$ and $\eta := \frac{\eta'}{|\eta'|}$. The map $U_w : \varphi \mapsto w\varphi$ is a unitary equivalence of the Hilbert spaces $(\mathcal{H}, \langle \cdot, \cdot \rangle)$ and $(\mathcal{H}, \langle \cdot, \cdot \rangle_{w^{-2}})$. The metric operator η' on $(\mathcal{H}, \langle \cdot, \cdot \rangle_{w^{-2}})$ gives rise to an isometry

$$\langle\langle \phi, \varphi \rangle\rangle_{\eta'} = \langle \phi, \eta' \varphi \rangle_{w^{-2}} = \langle \varphi, \eta \varphi \rangle = \langle\langle \phi, \varphi \rangle\rangle_\eta, \quad \forall \phi, \varphi \in \mathcal{H}$$

and by construction η fulfills the conditions (2.2). Operators which verify (2.2) are usually called *fundamental symmetries*. \blacktriangleleft

In view of Remark 2.2 we will always assume henceforth that the indefinite metric is generated by a fundamental symmetry $\eta^* = \eta = \eta^{-1}$. This fact allows us to introduce the corresponding self-adjoint projections

$$P_+ := \frac{\mathbb{1} + \eta}{2}, \quad P_- := \frac{\mathbb{1} - \eta}{2}$$

that determine the *fundamental* orthogonal decomposition (with respect to the scalar product) of \mathcal{H}

$$\mathcal{H} = \mathcal{H}_+ \oplus \mathcal{H}_-, \quad \mathcal{H}_\pm := \text{Ran } P_\pm = P_\pm \mathcal{H}.$$

It is not difficult to prove that the subspaces \mathcal{H}_+ and \mathcal{H}_- are also η -orthogonal. To see this, let $\phi_\pm \in \mathcal{H}_\pm$ and compute

$$\langle\langle \phi_+, \phi_- \rangle\rangle_\eta = \langle P_+ \phi_+, \eta P_- \phi_- \rangle = \langle \phi_+, P_+ \eta P_- \phi_- \rangle = 0$$

where we used the fact that P_\pm are self-adjoint projections, $P_+ \eta P_- = \pm P_+ P_- = 0$. Then we can also write the fundamental decomposition as

$$\mathcal{H} = \mathcal{H}_+ [\oplus]_\eta \mathcal{H}_-$$

where the symbol $[\oplus]_\eta$ denotes the orthogonality with respect to the inner product $\langle\langle \cdot, \cdot \rangle\rangle_\eta$.

Let $\kappa_\pm := \dim \mathcal{H}_\pm$. The *rank of indefiniteness* is defined by

$$\kappa := \min \{\kappa_+, \kappa_-\} \in \mathbb{N} \cup \{\infty\}.$$

The case $\kappa = 0$ coincides with $\eta = \mathbb{1}$ or $\eta = -\mathbb{1}$. This case is trivial since the resulting metric $\langle\langle \cdot, \cdot \rangle\rangle_\eta$ coincides, at worst up to a sign, with the initial scalar product $\langle \cdot, \cdot \rangle$. For this reason we will always assume henceforth that $\kappa > 0$.

Definition 2.3 (Krein space). *The indefinite metric space $(\mathcal{H}, \langle\langle \cdot, \cdot \rangle\rangle_\eta)$ associated to a fundamental symmetry $\eta^* = \eta = \eta^{-1}$ is called:*

- Pontryagin space if $0 < \kappa < \infty$;
- Krein space if $\kappa = \infty$.

The possibility to define Krein spaces as in Definition 2.3, namely as indefinite metric spaces with rank $\kappa = \infty$ is based on a crucial result by H. Langer [Lan1]. In this work we will use the name Krein space also for spaces of *finite* even dimension such that $\kappa_+ = \kappa_-$.

The notion of self-adjointness has a natural adaptation in the case of indefinite metric spaces. First of all let us recall that according to the existing literature (see e.g. [Tom, Ghe]) the η -adjoint of a densely defined (possible unbounded) operator A is by definition

$$A^\# := \eta A^* \eta.$$

Definition 2.4 (η -self-adjoint operators). *Let $(\mathcal{H}, \langle \langle \cdot, \cdot \rangle \rangle_\eta)$ be the indefinite metric space associated to a fundamental symmetry $\eta^* = \eta = \eta^{-1}$. A densely defined (possible unbounded) linear operator $H : \mathcal{H} \rightarrow \mathcal{H}$ is called η -symmetric if*

$$H \subseteq H^\# \quad (2.3)$$

and η -self-adjoint if

$$H = H^\# . \quad (2.4)$$

The notation

$$\mathbf{H}_\eta(\mathcal{H}) := \{H \in \mathcal{B}_{\text{lin}}(\mathcal{H}) \mid H = H^\#\}$$

is used to denote the set of bounded η -self-adjoint operators on \mathcal{H} .

From its very definition it follows that H is η -self-adjoint if and only if

$$\langle \langle \phi, H\varphi \rangle \rangle_\eta = \langle \phi, \eta H\varphi \rangle = \langle H\phi, \eta\varphi \rangle = \langle \langle H\phi, \varphi \rangle \rangle_\eta, \quad \forall \phi, \varphi \in \mathcal{D}(H)$$

which shows that H is “self-adjoint” with respect to the inner product $\langle \langle \cdot, \cdot \rangle \rangle_\eta$. Moreover, the relation $H^* = \eta H \eta$ says that H and H^* are unitarily equivalent and, in turn, that η is a bijection between the dense domains $\mathcal{D}(H)$ and $\mathcal{D}(H^*)$. Finally, from the equality

$$\langle \phi, \eta H\varphi \rangle = \langle H\phi, \eta\varphi \rangle = \langle \eta H\phi, \varphi \rangle, \quad \forall \phi, \varphi \in \mathcal{D}(H)$$

one also infers the relation

$$\eta H = (\eta H)^* . \quad (2.5)$$

2.2. Unitray vs. η -unitary operators. A densely defined linear operator U acting on the Hilbert $(\mathcal{H}, \langle \cdot, \cdot \rangle)$ is unitary if

$$\langle \phi, \varphi \rangle = \langle U\phi, U\varphi \rangle, \quad \forall \phi, \varphi \in \mathcal{D}(U) .$$

This condition immediately allows to conclude that U is bounded in norm by 1 in its domain $\mathcal{D}(U)$ and so can be extended to a bounded operator on \mathcal{H} . Moreover, U turns out to be invertible with inverse $U^{-1} = U^*$. An operator U on $(\mathcal{H}, \langle \cdot, \cdot \rangle)$ is *anti-unitary* if it is *anti-linear* and

$$\langle \phi, \varphi \rangle = \langle U\varphi, U\phi \rangle, \quad \forall \phi, \varphi \in \mathcal{H} .$$

Every anti-linear operator U is necessarily bounded and invertible. In fact, given a *complex conjugation* C on \mathcal{H} , one has that the anti-linear operator U can be represented as the product $U = CW_U$ of the complex conjugation C times the unitary $W_U := CU$. It is worth to recall that a complex conjugation is any *anti-linear* map which admits a *real* (orthonormal) basis of vectors in the sense that there exists a basis $\{\psi_j\}_{j \in \mathbb{N}} \subset \mathcal{H}$ such that $C\psi_j = \psi_j$. As a consequence one has that $C^* = C = C^{-1}$.

Things become less simple in the case of indefinite metric space $(\mathcal{H}, \langle \langle \cdot, \cdot \rangle \rangle_\eta)$. First of all it should be noted that the requirement

$$\langle \langle \phi, \varphi \rangle \rangle_\eta = \langle \langle U\phi, U\varphi \rangle \rangle_\eta, \quad \forall \phi, \varphi \in \mathcal{D}(U) . \quad (2.6)$$

does *not* allow to conclude, as in the positive metric case, that U is bounded when \mathcal{H} is infinite dimensional [AI, Shm, Tom, Ghe]. This is a fundamental difference which makes the notion of unitarity more difficult in the case of indefinite metric spaces. For these reasons, in the mathematical literature on indefinite metric spaces the boundedness property is part of the definition of the unitarity of an operator (see e. g. [AI, Chapter 2, Definition 5.1]). However, since this property is in general not justifiable on the basis of purely physical considerations and it is in general not shared by operators occurring in the physical applications (cf. Remark 2.7) we prefer to omit it, at least in this general preliminary presentation. The following definition is borrowed from [BMS, Definitions 2 & 3]:

Definition 2.5 (η -unitary and η -anti-unitary operators). *Let $(\mathcal{H}, \langle \langle \cdot, \cdot \rangle \rangle_\eta)$ be the indefinite metric space associate to a fundamental symmetry $\eta^* = \eta = \eta^{-1}$. Let $U : \mathcal{H} \rightarrow \mathcal{H}$ be an operator (not linear a priori) with domain $\mathcal{D}(U) \subseteq \mathcal{H}$ and range $\mathcal{R}(U) \subseteq \mathcal{H}$. Assume that:*

- (i) *Domain and range are dense sets, i. e. $\overline{\mathcal{D}(U)} = \mathcal{H} = \overline{\mathcal{R}(U)}$;*

The operator U is called η -unitary if:

(ii) *The condition*

$$\langle\langle\phi, \varphi\rangle\rangle_\eta = \langle\langle U\phi, U\varphi\rangle\rangle_\eta, \quad \forall \phi, \varphi \in \mathcal{D}(U) \quad (2.7)$$

holds true.

It is called η -anti-unitary if:

(ii') *The condition*

$$\langle\langle\phi, \varphi\rangle\rangle_\eta = \langle\langle U\varphi, U\phi\rangle\rangle_\eta, \quad \forall \phi, \varphi \in \mathcal{D}(U) \quad (2.8)$$

holds true.

Remark 2.6 (Comparison with the mathematical nomenclature). Possible unbounded operators which meet condition (ii) of Definition 2.5 are usually called η -isometric (see e. g. [AI, Chapter 2, Definition 5.1]). However, the notion of η -unitarity is slightly stronger in view of the condition (i) which assures the density of the domain and the range of U . Let us observe that the density of $\mathcal{D}(U)$ and condition (ii) are enough to prove that $U^{-1} \subseteq \eta U^* \eta = U^\sharp$ [BMS, Proposition 2]. Sometime, the stronger condition $U^{-1} = U^\sharp$ is used as a definition of η -unitary operator (see e. g. [Tom, Ghe]). Clearly, whenever U is a bounded operator all these definitions coincide. ◀

Remark 2.7 (Unbounded η -unitarity operators). The construction of explicit examples of unbounded η -unitarity operators is in general a non trivial task. For instance the examples presented in [AI, Chapter 2, Example 5.5] or [Ghe, Remark 5.8 (c)] are quite implicit. On the other hand is not difficult to show that η -unitarity operators can have arbitrarily large norm (cf. Lemma B.1). A relevant physical examples of unbounded η -unitarity operators is provided by the *Gupta formulation* of quantum electrodynamics. In this theory the Lorentz transformations are described by unbounded and densely defined operators (see e. g. the discussion in [Pru, Section 7.4] and references therein). ◀

Let us observe that, as for the usual Hilbert space case, one can express every η -anti-unitary operator U as the product $U = CW_U$ where C is a complex conjugation which meets

$$C\eta = \eta C \quad (2.9)$$

and $W_U := CU$ is a η -unitary operator⁴.

Since we allow η -unitary and η -anti-unitary operators to be unbounded, it becomes important to investigate the closure properties of these operators. The following result holds true.

Proposition 2.8 ([BMS, Propositions 2 & 3]). *Every η -unitary (resp. η -anti-unitary) operator is a linear (resp. anti-linear) and closable operator. Moreover, it has an inverse which is in turn an η -unitary (resp. η -anti-unitary) operator.*

According to the content of Remark 2.6 one can characterize the *group* of bounded η -unitary operators by

$$\mathcal{U}_\eta(\mathcal{H}) := \{U \in \mathcal{GL}(\mathcal{H}) \mid U^\sharp = U^{-1}\} \quad (2.10)$$

where $\mathcal{GL}(\mathcal{H}) \subset \mathcal{B}_{\text{lin}}(\mathcal{H})$ is the group of invertible bounded linear operators on \mathcal{H} and U^\sharp denotes the η -adjoint of U . We will use the notation $\mathcal{AU}_\eta(\mathcal{H})$ for the set of η -anti-unitary *bounded* operators. This set can be characterized in terms of the group $\mathcal{U}_\eta(\mathcal{H})$ once a complex conjugation C which meets (2.9) has been chosen. In this case one has

$$\mathcal{AU}_\eta(\mathcal{H}) := \{U = CW \mid W \in \mathcal{U}_\eta(\mathcal{H})\}. \quad (2.11)$$

The following result will be used several times.

Lemma 2.9. *Each operator $U \in \mathcal{U}_\eta(\mathcal{H}) \sqcup \mathcal{AU}_\eta(\mathcal{H})$ meets $\eta U^* = U^{-1}\eta$.*

⁴ A complex conjugation C which meets condition (2.9) is sometimes called an η -complex conjugation. It is not difficult to construct η -complex conjugations C . It is sufficient to select a basis of $\{\psi_j\}_{j \in \mathbb{N}} \subset \mathcal{H}$ which diagonalizes η and to define C as the anti-linear operator fixing the basis, namely $C\psi_j = \psi_j$.

Proof. For η -unitary operators the relation follows by multiplying on the left the equality $U^\sharp = U^{-1}$ with η . For the case of η -anti-unitary operators it is first necessary to observe that the relation $(AB)^* = B^*A^*$ holds true for any pair of bounded operators, regardless whether they are linear or anti-linear, provided that the adjoint for anti-linear operators is defined as in the Note 3. Then, for any $U = CW \in \text{AU}_\eta(\mathcal{H})$ one has

$$\eta U^* = \eta(CW)^* = \eta W^*C = W^{-1}\eta C = (W^{-1}C)\eta = U^{-1}\eta$$

where the equality $C^* = C = C^{-1}$ and the commutativity $\eta C = C\eta$ have been used. \blacksquare

For sake of clarity let us emphasize that the symbol \sqcup used in the claim of Lemma 2.9 stands for the *disjoint union* and is justified by the evidence $\text{U}_\eta(\mathcal{H}) \cap \text{AU}_\eta(\mathcal{H}) = \emptyset$.

2.3. Pseudo-unitarity in indefinite metric spaces. Just as the concept of anti-unitary operators arise when one uses the complex structure of the Hilbert spaces, a new class of operators naturally occurs when the inner product is allowed to be indefinite.

Definition 2.10 (η -pseudo-unitary and η -pseudo-anti-unitary operators). *Let $(\mathcal{H}, \langle\langle \cdot, \cdot \rangle\rangle_\eta)$ be the indefinite metric space associate to a fundamental symmetry $\eta^* = \eta = \eta^{-1}$. Let $U : \mathcal{H} \rightarrow \mathcal{H}$ be an operator (not linear a priori) with domain $\mathcal{D}(U) \subseteq \mathcal{H}$ and range $\mathcal{R}(U) \subseteq \mathcal{H}$. Assume that:*

- (0) *The rank of indefiniteness meets $\kappa = \kappa_+ = \kappa_-$;*
- (i) *Domain and range are dense sets, i. e. $\overline{\mathcal{D}(U)} = \mathcal{H} = \overline{\mathcal{R}(U)}$.*

The operator U is called η -pseudo-unitary if:

- (iii) *The condition*

$$\langle\langle \phi, \varphi \rangle\rangle_\eta = - \langle\langle U\phi, U\varphi, \rangle\rangle_\eta, \quad \forall \phi, \varphi \in \mathcal{D}(U) \quad (2.12)$$

holds true.

It is called η -pseudo-anti-unitary if:

- (iii') *The condition*

$$\langle\langle \phi, \varphi \rangle\rangle_\eta = - \langle\langle U\varphi, U\phi, \rangle\rangle_\eta, \quad \forall \phi, \varphi \in \mathcal{D}(U) \quad (2.13)$$

holds true.

First of all, let us remark that the occurrence of pseudo-unitary operators in Definition 2.10 is tightly bound to the indefinite metric. In fact in a Hilbert space the condition $\langle\phi, \phi\rangle = -\langle U\phi, U\phi, \rangle$ is untenable when $\phi \neq 0$. Secondly, let us observe that an η -pseudo-unitary or η -pseudo-anti-unitary operator may exist only in indefinite metric spaces in which the the eigenvalues $+1$ and -1 of the metric operator η have the same multiplicity and this is the reason for the condition (0) of Definition 2.10. This is related to the fact that an η -pseudo-unitary operator is necessarily invertible and its inverse is still an η -pseudo-unitary operator which satisfies $U^{-1} \subseteq -\eta U^* \eta$ which can be proved exactly as in the proof of [BMS, Proposition 2]. The last inclusion is the key ingredient to prove the *Cartan decomposition* which states that U can be expressed as the product $U = RV_U$ of a (positive) η -unitary operator V_U and a unitary and η -pseudo-unitary operator $R^* = R^{-1} = -R^\sharp$ (the proof of this fact is identical to that of [Tom, Theorem 2.1.1]). The last equality implies that

$$R\eta = -\eta R, \quad (2.14)$$

which can be read as the fact that η and $-\eta$ are unitarily equivalent. However this is possible only if the the condition (0) of Definition 2.10 is verified.

Remark 2.11. One of the major interests of this work is to study the consequences of the existence of pseudo-unitary symmetries on the classification of topological phases. For this reason henceforth we will assume the validity of condition (0) in Definition 2.10. Said differently, we will always work with finite or infinite dimensional Krein spaces (cf. Definition 2.3). \blacktriangleleft

Remark 2.12 (Comparison with the mathematical nomenclature). In the mathematical literature the notion of η -pseudo-unitary operators appears in different forms and with different names. For instance in [Tom, KSh] this type of operators are defined through the stronger condition $U^{-1} = -U^\sharp$ and are called \sharp -imaginary in [Tom] and J -skew-unitary in [KSh]. In this work we use definitions and terminology introduced in [BMS]. ◀

A special subclass of unitary operator which meet (2.14) are those which in addition are self-adjoint, $R^* = R = R^{-1}$. This operators are called η -reflecting operators⁵. Once an η -reflecting operator R is given one can represent any η -pseudo-unitary operator U according to the Cartan representation $U = RV_U$ where $V_U := RU$ is automatically an η -unitary operator. A similar conclusion holds true for η -pseudo-anti-unitary operators. In fact, given an η -complex conjugation C (cf. Note 4) and an η -reflecting operators R such that $CR = RC$ one has that any η -pseudo-anti-unitary operator U can be represented as $U = CRT_U$ where $T_U := RCU$ is η -unitary. By combining this last observation with the content of Proposition 2.8 and Proposition 2.13 one obtains the following description for η -pseudo-unitary and η -pseudo-anti-unitary operators:

Proposition 2.13 ([BMS, Proposition 4]). *Every η -pseudo-unitary (resp. η -pseudo-anti-unitary) operator is a linear and closable operator. Moreover, it has an inverse which is in turn an η -pseudo-unitary (resp. η -pseudo-anti-unitary) operator.*

The sets of η -pseudo-unitary bounded operators and η -pseudo-anti-unitary bounded operators can be characterized in terms of the group $U_\eta(\mathcal{H})$ by

$$PU_\eta(\mathcal{H}) := \{U = RV \mid V \in U_\eta(\mathcal{H})\} \quad (2.15)$$

and

$$PAU_\eta(\mathcal{H}) := \{U = CRT \mid T \in U_\eta(\mathcal{H})\} \quad (2.16)$$

once a commuting pair C and R has been chosen. The following result extends Lemma 2.9.

Lemma 2.14. *Each operator $U \in PU_\eta(\mathcal{H}) \sqcup PAU_\eta(\mathcal{H})$ meets $\eta U^* = -U^{-1}\eta$.*

Proof. Let $U = RV$ be an η -pseudo-unitary operator. It follows that

$$\eta U^* = \eta(RV)^* = (\eta V^*)R = (V^{-1}\eta)R = -(V^{-1}R)\eta = -U^{-1}\eta$$

where Lemma 2.9 has been used. The last relation also shows that $U \in PU_\eta(\mathcal{H})$ if and only if U is a linear operator which fulfills $U^\sharp = -U^{-1}$. In the case $U = CRT = CW$ is η -pseudo-anti-unitary one has that $W = RT$ is η -pseudo-unitary, therefore

$$\eta U^* = \eta(CW)^* = \eta W^*C = -W^{-1}\eta C = -(W^{-1}C)\eta = -U^{-1}\eta$$

just as in the final part of the proof of Lemma 2.9. ■

3. QUANTUM SYMMETRIES: THE WIGNER'S THEOREM

In this Section the classical Wigner's theorem [Wig] is revisited in a geometric language according to the presentation provided in [FM, Section 1.1]. The same description is then adapted for the extended version of the Wigner's theorem in Krein spaces proved in [BMS].

⁵ It is clear from the definition that the role of an η -reflecting operator R is to interchange an orthonormal basis of \mathcal{H}_+ with an orthonormal basis of \mathcal{H}_- and vice versa. Moreover, if one defines a η -complex structure C which preserves the bases of \mathcal{H}_+ and \mathcal{H}_- then C and R automatically commute.

3.1. The Hilbert space case. According to the ordinary formulation of Quantum Mechanics the *state* of a quantum system is described by an element of the projective space $\mathbb{P}\mathcal{H}$ of a complex separable Hilbert space $(\mathcal{H}, \langle \cdot, \cdot \rangle)$. In other words, a state is a line, or *ray*, of vectors. The transition probability between two states $[\phi], [\varphi] \in \mathbb{P}\mathcal{H}$ is given by

$$\text{Prob}([\phi], [\varphi]) := \frac{|\langle \phi, \varphi \rangle|^2}{\langle \phi, \phi \rangle \langle \varphi, \varphi \rangle}, \quad (3.1)$$

where $\phi \in [\phi]$ and $\varphi \in [\varphi]$ are suitable representatives. Equation (3.1) defines a symmetric function

$$\text{Prob} : \mathbb{P}\mathcal{H} \times \mathbb{P}\mathcal{H} \longrightarrow [0, 1].$$

Since the transition probabilities are interpreted as the possible outcomes of physical measurements, one is led to call *quantum symmetry* any invertible transformation (an automorphism) of the state space $S : \mathbb{P}\mathcal{H} \rightarrow \mathbb{P}\mathcal{H}$ preserving the symmetric function Prob . The set of these transformations will be denoted by $\text{QS}(\mathbb{P}\mathcal{H})$. A priori a quantum symmetry S is not required to have other additional structural properties. However a fundamental theorem by Wigner [Wig] states that every quantum symmetry is implemented by means of a linear map $U_S : \mathcal{H} \rightarrow \mathcal{H}$ which can be either *unitary* or *anti-unitary*.

Following [FM, Section 1.1] (see also [Thi, Kub]) one can recast the Wigner's theorem in the following terms. Let

$$\text{QS}(\mathcal{H}) := \text{U}(\mathcal{H}) \sqcup \text{AU}(\mathcal{H})$$

be the set of all unitary and anti-unitary transformations of \mathcal{H} . It is a group since the composition of two anti-unitary transformations is unitary and we refer to $\text{QS}(\mathcal{H})$ as the group of *linear* quantum symmetries. This group fits in a group extension

$$1 \longrightarrow \text{U}(\mathcal{H}) \longrightarrow \text{QS}(\mathcal{H}) \xrightarrow{\varpi} \mathbb{Z}_2 \longrightarrow 1 \quad (3.2)$$

where $\mathbb{Z}_2 := \{\pm 1\}$ is the cyclic group of order 2 and the kernel of the homomorphism ϖ is the group of unitary operators. Said differently, one has that $\varpi(U) = +1$ if U is a unitary operator and $\varpi(U) = -1$ if U is anti-unitary. The Wigner theorem asserts that there is a group extension⁶

$$1 \longrightarrow \mathbb{S}^1 \xrightarrow{\iota} \text{QS}(\mathcal{H}) \xrightarrow{\pi} \text{QS}(\mathbb{P}\mathcal{H}) \longrightarrow 1 \quad (3.3)$$

where $\mathbb{S}^1 := \{z \in \mathbb{C} \mid |z| = 1\}$ denotes the group of complex numbers of modulus one and the map ι associates to each $\lambda \in \mathbb{S}^1$ the *scaling* operator $\iota(\lambda) := \lambda \mathbb{1}$. The projection map π sends the linear quantum symmetry U into the quantum symmetry $\pi(U)$ which acts on the rays according to $\pi(U)[\psi] := [U\psi]$. Clearly $\iota(\mathbb{S}^1) \subseteq \text{QS}(\mathcal{H})$ agrees with the kernel of π and every quantum symmetry $S \in \text{QS}(\mathbb{P}\mathcal{H})$ lifts to a linear quantum symmetry $U_S := \pi^{-1}(S) \in \text{QS}(\mathcal{H})$ which is unique up to the composition with a scaling operator $\lambda \mathbb{1} \in \iota(\mathbb{S}^1)$. Even though \mathbb{S}^1 is commutative, the extension (3.3) is *not* central since $\iota(\mathbb{S}^1)$ is not in the center of $\text{QS}(\mathcal{H})$. In fact, one has that

$$U(\lambda \mathbb{1}) U^{-1} = \lambda^{\varpi(U)} \mathbb{1} = \begin{cases} \lambda \mathbb{1} & \text{if } \varpi(U) = +1 \\ \lambda^{-1} \mathbb{1} & \text{if } \varpi(U) = -1. \end{cases} \quad (3.4)$$

Hence (3.3) is not a central extension but rather a *twisted* (central) extension according to [FM, Definition 1.7] where the twisting is expressed by the map $\varpi : \text{QS}(\mathcal{H}) \rightarrow \mathbb{Z}_2$ according to (3.4).

Remark 3.1 (Choice of the topology). The group $\text{QS}(\mathcal{H})$ and its subgroup $\text{U}(\mathcal{H})$ can be endowed with different topologies. In the uniform topology given by the operator norm (2.1) $\text{U}(\mathcal{H})$ is a closed and connected space with the structure of a topological *Banach-Lie* group. Moreover, if \mathcal{H} is infinite dimensional and separable then $\text{U}(\mathcal{H})$ is contractible to the identity [Kui]. The bijection $C : \text{U}(\mathcal{H}) \rightarrow \text{AU}(\mathcal{H})$, induced by any complex conjugation, implies that also $\text{AU}(\mathcal{H})$ is connected (to C) and contractible in the infinite dimensional case. In summary $\text{QS}(\mathcal{H})$ becomes a topological group

⁶For basic facts about the notion of group extension we refer to [Bro, Chapter 4] or [TW] or [FM, Appendix A].

with two disconnected components when endowed with the uniform topology. Unfortunately the uniform topology is too strong for many purposes. For instance the one-parameter *time evolution* group $\mathbb{R} \ni t \mapsto e^{itH} \in U(\mathcal{H})$ fails to be uniformly continuous when the self-adjoint operator H is unbounded. To solve this problem one needs to weaken the topology on $QS(\mathcal{H})$ without compromising the structure of topologic group. One possible candidate which makes the time evolution group continuous is the *strong topology* generated by the semi-norms $s_\psi(A) := \|A\psi\|$ as ψ varies in \mathcal{H} . Even though the operator product is not continuous in general with respect the strong topology, it is continuous when restricted to set of equi-bounded operators. This is the basic observation that allows to show that $U(\mathcal{H})$ is a topological *Polish group*⁷ when endowed with the strong topology [Nee, Proposition II.1] or [EU, Theorem 1.2] (see also [Sch2, Proposition 1]). In one of the foundational papers for twisted equivariant K-theory [ASe], Atiyah and Segal proposed a different way to topologize $U(\mathcal{H})$. In order to deal with equivariant Hilbert bundles and its relation with its associated unitary principal equivariant bundles, they claimed that one is obliged to consider the *compact-open topology*⁸ on the structural group $U(\mathcal{H})$. A fact that seems to be remarkable is that the compact-open topology and the strong topology agree on $U(\mathcal{H})$ (see [EU, Lemma 1.8.] or [Sch2, Proposition 2]). This immediately implies that $U(\mathcal{H})$ is a topological group also with respect to the compact-open topology. This last fact has been erroneously denied in [ASe, pg. 321]. The equality between compact-open and strong topology on $U(\mathcal{H})$ along with [ASe, Proposition A2.1.] implies that $U(\mathcal{H})$ is a contractible space also with respect to the strong topology if \mathcal{H} is infinite dimensional and separable. Clearly, the strong topology can be defined also on $AU(\mathcal{H})$ and the bijection implemented by C becomes a homeomorphism. In summary $QS(\mathcal{H})$ endowed with the strong topology turns out to be a topological Polish group with two disconnected components which are contractible when \mathcal{H} is infinite dimensional and separable. The topology on the projective space $U(\mathbb{P}\mathcal{H})$ is induced as the quotient topology associated with the short exact sequence

$$1 \longrightarrow \mathbb{S}^1 \xrightarrow{i} U(\mathcal{H}) \xrightarrow{\pi} U(\mathbb{P}\mathcal{H}) \longrightarrow 1. \quad (3.5)$$

Interestingly, the strong topology on $U(\mathcal{H})$ descends to the strong topology on $U(\mathbb{P}\mathcal{H})$ and $U(\mathcal{H})$ can be seen as a principal bundle over $U(\mathbb{P}\mathcal{H})$ with typical fiber \mathbb{S}^1 [Sim, Theorem 1]. The homotopy exact sequence induced by (3.5) provides that $\pi_k(U(\mathbb{P}\mathcal{H})) = \pi_{k-1}(\mathbb{S}^1)$ showing that $U(\mathbb{P}\mathcal{H})$ is a model for the Eilenberg-MacLane space $K(\mathbb{Z}, 2)$. In much the same way the space $QS(\mathbb{P}\mathcal{H})$ can be topologized with the (strong) quotient topology inherited by $QS(\mathcal{H})$. In this way also $QS(\mathcal{H})$ can be seen as a principal bundle over $QS(\mathbb{P}\mathcal{H})$ with typical fiber \mathbb{S}^1 . As a final remark let us observe that the map $\varpi : QS(\mathcal{H}) \rightarrow \mathbb{Z}_2$ which enters in the group extension (3.2) is a continuous homomorphism with respect to the natural discrete topology on \mathbb{Z}_2 and the strong topology on $QS(\mathcal{H})$. ◀

Remark 3.2 (Twisted extensions and cohomology). It is a well known result that the central extensions of a group \mathbb{G} by an abelian group \mathbb{A} are classified by classes in the group-cohomology $H^2(\mathbb{G}, \mathbb{A})$ (see [Bro, Chapter 4, Theorem 3.12] or [TW, Proposition 3.4]). Also twisted extensions of the type of (3.3) - (3.4) can be classified by means of a cohomology theory. Let $\varpi : \mathbb{G} \rightarrow \mathbb{Z}_2$ be the homomorphism that defines the twisted extension and consider the (left) group action of \mathbb{G} over \mathbb{A} given by $g \cdot a := a^{\varpi(g)}$. The group \mathbb{A} endowed with such a \mathbb{G} -action is denoted by \mathbb{A}_ϖ and provides a system of twisted coefficients for the group cohomology of \mathbb{G} . Then, the ϖ -twisted central extensions of \mathbb{G} by \mathbb{A} are classified by the elements of $H^2(\mathbb{G}, \mathbb{A}_\varpi)$. Therefore, the twisted central extension (3.3) corresponds to an element of the cohomology group $H^2(QS(\mathbb{P}\mathcal{H}), \mathbb{S}^1_\varpi)$ where the homomorphism ϖ distinguishes between unitary and anti-unitary symmetries. ◀

⁷A Polish group is a completely metrizable topological group. To this respect, it is interesting to observe that when \mathcal{H} is infinite dimensional then $U(\mathcal{H})$ is *not* closed with respect to the strong topology (as erroneously claimed below [Sch2, Proposition 3]) even though $U(\mathcal{H})$ is closed with respect to the metric which generates the topology equivalent to the strong topology. This fact, although looking surprising at first sight, is not absurd. For instance, as proved in [Tak, Chapter II, Proposition 4.9], $U(\mathcal{H})$ is in fact closed with respect to the $*$ -strong topology even though the strong and the $*$ -strong topology coincide. As commented in [Tak, Chapter II, Remark 4.10] this phenomenon is a manifestation of the fact that the uniform structures induced by the strong and the $*$ -strong topologies are distinct.

⁸For more details about the compact-open topology on $U(\mathcal{H})$ we refer to [ASe, Appendix 1] and [FM, Appendix D].

3.2. The Krein space case. The discussion of the Wigner's theorem in Krein spaces presents some technical difficulties with respect to the Hilbert space case discussed in Section 3.1. The transition probability between two rays $[\phi], [\varphi] \in \mathbb{P}\mathcal{H}$ of a Krein space $(\mathcal{H}, \langle \cdot, \cdot \rangle_\eta)$ is naturally defined by mimicking (3.1):

$$\text{Prob}_\eta([\phi], [\varphi]) := \frac{|\langle \phi, \varphi \rangle_\eta|^2}{\langle \phi, \phi \rangle \langle \varphi, \varphi \rangle} = \frac{|\langle \phi, \eta \varphi \rangle|^2}{\langle \phi, \phi \rangle \langle \varphi, \varphi \rangle}, \quad (3.6)$$

Let us point out that in both (3.1) and (3.6) the normalization of the vectors is defined with respect to the Hilbert structure. In fact one cannot normalize the rays with respect to the indefinite inner product due to the presence of non-zero vectors of zero η -length. Also equation (3.6) defines a symmetric function

$$\text{Prob}_\eta : \mathbb{P}\mathcal{H} \times \mathbb{P}\mathcal{H} \longrightarrow [0, 1]$$

and, in analogy with the Hilbert space case, we are tempted to call η -quantum symmetry any invertible transformation of the state space $S : \mathbb{P}\mathcal{H} \rightarrow \mathbb{P}\mathcal{H}$ preserving the symmetric function Prob_η . However the situation in the framework of Krein spaces is slightly more complicated than the Hilbert space case due to the fact that maps S which preserve Prob_η in general are not defined for all the rays in $\mathbb{P}\mathcal{H}$. However, it is reasonable to require that a “good” quantum symmetry S , in addition to being invertible, has to be defined on a set of rays associated to a dense linear manifold $\mathcal{D} \subseteq \mathcal{H}$, and has to map the rays associated to \mathcal{D} onto a set of rays associated to a dense linear manifold $\mathcal{D}' \subseteq \mathcal{H}$. The requirement of working with dense manifolds is suggested by physical interesting examples (cf. Remark 2.7). If one accepts the density condition above as part of the definition of an η -quantum symmetry S then the extended version of the Wigner's theorem due to Bracci, Morchio and Strocchi [BMS] states that there exists a (possible unbounded) operator $U_S : \mathcal{H} \rightarrow \mathcal{H}$ which implements S and U_S is either η -unitary or η -anti-unitary or η -pseudo-unitary or η -pseudo-anti-unitary (the last two cases can occur only when condition (0) in Definition 2.10 holds).

However for the aims of this work we need to consider only the *bounded* version of the extended Wigner's theorem. Let us introduce the group

$$\text{QS}_\eta(\mathcal{H}) := \text{U}_\eta(\mathcal{H}) \sqcup \text{AU}_\eta(\mathcal{H}) \sqcup \text{PU}_\eta(\mathcal{H}) \sqcup \text{PAU}_\eta(\mathcal{H})$$

of the (*bounded*) linear η -quantum symmetries. This group fits in a group extension

$$1 \longrightarrow \text{U}_\eta(\mathcal{H}) \longrightarrow \text{QS}_\eta(\mathcal{H}) \xrightarrow{(\wp, \varpi)} \mathbb{Z}_2 \times \mathbb{Z}_2 \longrightarrow 1 \quad (3.7)$$

and the kernel of the homomorphism (\wp, ϖ) is the group of the η -unitary operators. More specifically, the maps \wp and ϖ determine the nature of $U \in \text{QS}_\eta(\mathcal{H})$ according to the following relation:

$$\langle \langle \phi, \varphi \rangle \rangle_\eta = \wp(U) |\langle \langle U\phi, U\varphi, \rangle \rangle_\eta|^{1-\varpi(U)} \langle \langle U\phi, U\varphi, \rangle \rangle_\eta^{\varpi(U)}, \quad \forall \phi, \varphi \in \mathcal{H}. \quad (3.8)$$

The convention on the signs, which can be extended straightforwardly to the case of unbounded symmetries, is conveniently summarized in Table 1.3 of Section 1.

The generalized Wigner's theorem proved in [BMS] asserts that the set of (*bounded*) η -quantum symmetries $\text{QS}_\eta(\mathbb{P}\mathcal{H})$ fits in the group extension

$$1 \longrightarrow \mathbb{S}^1 \xrightarrow{\iota} \text{QS}_\eta(\mathcal{H}) \xrightarrow{\pi} \text{QS}_\eta(\mathbb{P}\mathcal{H}) \longrightarrow 1 \quad (3.9)$$

where the maps ι and π have the same meaning as in (3.3). Again this extension is *not* central but *twisted* by the homomorphism $(\wp, \varpi) : \text{QS}_\eta(\mathcal{H}) \rightarrow \mathbb{Z}_2 \times \mathbb{Z}_2$ according to

$$U(\lambda \mathbb{1})U^{-1} = \wp(U)\lambda^{\varpi(U)}\mathbb{1} \quad (3.10)$$

Remark 3.3 (Choice of the topology). By making a parallel with Remark 3.3 it is natural to ask in which way one can topologize $\text{QS}_\eta(\mathcal{H})$ and $\text{QS}_\eta(\mathbb{P}\mathcal{H})$ in order to make the exact sequence (3.9) topological. First of all one can notice that the group $\text{QS}_\eta(\mathcal{H})$ is made by four disconnected components (we are tacitly assuming the validity of condition (0) in Definition 2.10) and the three components $\text{AU}_\eta(\mathcal{H})$, $\text{PU}_\eta(\mathcal{H})$, and $\text{PAU}_\eta(\mathcal{H})$ are all related to the sub-group $\text{U}_\eta(\mathcal{H})$ by the election of a η -complex

conjugation C as in (2.9) and a η -reflecting operators R as in (2.14) such that $CR = RC$. This implies that it is enough to topologize the subgroup $U_\eta(\mathcal{H})$ in order to endow the full group $AU_\eta(\mathcal{H})$ with a topology. Once a topology on $AU_\eta(\mathcal{H})$ is defined one can induce the quotient topology on $QS_\eta(\mathbb{P}\mathcal{H})$ by means of the projection π . The characterization (2.10) implies that $U_\eta(\mathcal{H})$ is an algebraic subgroup of $GL(\mathcal{H})$ which is a topological group when endowed with the uniform topology given by the operator norm. As a consequence $U_\eta(\mathcal{H})$ turns out to be a (closed) Banach-Lie group when endowed with the topology induced by the operator norm [Upm, Section 23]. Unfortunately, as mentioned in Remark 3.3, the uniform topology is generally too strong for many purposes. For this reason one could be tempted to endow $U_\eta(\mathcal{H})$ with the strong topology. However, the set $U_\eta(\mathcal{H})$ contains elements of arbitrarily large norm (cf. Lemma B.1) and the product of operators is generally not continuous with respect to the strong topology. Thus $U_\eta(\mathcal{H})$ fails to be a topological group with respect to the strong topology. A possible strategy to circumvent these problems is to use the compoact-open topology trying to adapt the results from [ASe, Appendix 1] and [FM, Appendix D]. ◀

3.3. Symmetries of η -self-adjoint operators. This Section is devoted to the study of the transformation property of η -self-adjoint operators under linear η -quantum symmetries. It is worth remembering that a densely defined (possibly unbounded) linear operator H is η -self-adjoint if and only if $H = H^\sharp = \eta H^* \eta$. This condition implies the relation $\mathcal{D}(H) = \eta[\mathcal{D}(H^*)]$ among the domains of H and H^* and assures that H is a closed operator (H^* is closed by definition of adjoint and the same holds for H^\sharp due to the boundedness and the invertibility of η).

A (possibly unbounded) densely defined linear symmetry U is said to be of type (\wp, ϖ) if it meets equation (3.8) for all $\phi, \varphi \in \mathcal{D}(U)$. It follows from Proposition 2.8 and Proposition 2.13 that any symmetry U of type (\wp, ϖ) is automatically closable and invertible. Therefore, without loss of generality, one can always consider U and U^{-1} as closed symmetries of the same type.

Although H and U are closed and densely defined operators the product UHU^{-1} may be not well defined in general. As a matter of fact the product has to be initially defined on the initial domain

$$\mathcal{D}_0 := \left\{ \varphi \in \mathcal{D}(U^{-1}) \mid U^{-1}\varphi \in \mathcal{D}(H) \text{ and } H(U^{-1}\varphi) \in \mathcal{D}(U) \right\}.$$

This domain might be non-dense or even empty. Therefore, in order to have a good definition of UHU^{-1} one must assume the density of \mathcal{D}_0 .

Theorem 3.4. *Let H be an η -self-adjoint operator with dense domain $\mathcal{D}(H)$. Let U be a linear symmetry U of type (\wp, ϖ) defined on the dense domain $\mathcal{D}(U)$. Assume that the product UHU^{-1} is initially defined on the dense domain \mathcal{D}_0 . Then the operator UHU^{-1} is closable on \mathcal{D}_0 and η -symmetric, meaning that the inclusion*

$$\eta UHU^{-1} \subseteq (UHU^{-1})^* \eta$$

holds true.

Proof. Since UHU^{-1} is densely defined it is adjointable, namely $(UHU^{-1})^*$ exists and is closed. Moreover one has $(\eta UHU^{-1})^* = (UHU^{-1})^* \eta$ since η is bounded and self-adjoint. Let $\phi, \varphi \in \mathcal{D}_0$ and observe that $UHU^{-1}\varphi \in \mathcal{R}(U) \subseteq \mathcal{D}(U^{-1})$. A straightforward computation shows that

$$\begin{aligned} \langle \phi, \eta UHU^{-1}\varphi \rangle &= \langle \phi, UHU^{-1}\varphi \rangle_\eta \\ &= \wp(U^{-1}) \left| \langle \langle U^{-1}\phi, HU^{-1}\varphi, \rangle \rangle_\eta \right|^{1-\varpi(U^{-1})} \langle \langle U^{-1}\phi, HU^{-1}\varphi, \rangle \rangle_\eta^{\varpi(U^{-1})} \\ &= \wp(U^{-1}) \left| \langle \langle HU^{-1}\phi, U^{-1}\varphi, \rangle \rangle_\eta \right|^{1-\varpi(U^{-1})} \langle \langle HU^{-1}\phi, U^{-1}\varphi, \rangle \rangle_\eta^{\varpi(U^{-1})} \end{aligned} \quad (3.11)$$

where the second equality follows from (3.8) and the third by the η -self-adjointness of H . By exploiting again (3.8) and using that $\varpi(U^{-1})\varpi(U) = 1$ in view of the fact that U and U^{-1} are of the same type one obtains

$$\langle \langle HU^{-1}\phi, U^{-1}\varphi, \rangle \rangle_\eta^{\varpi(U^{-1})} = \wp(U) \left| \langle \langle UHU^{-1}\phi, \varphi, \rangle \rangle_\eta \right|^{\varpi(U^{-1})-1} \langle \langle UHU^{-1}\phi, \varphi, \rangle \rangle_\eta. \quad (3.12)$$

By inserting (3.12) in (3.11) one gets

$$\langle\langle\phi, UHU^{-1}\varphi\rangle\rangle_\eta = \wp(U^{-1}) \wp(U) \left(\frac{|\langle\langle HU^{-1}\phi, U^{-1}\varphi, \rangle\rangle_\eta|}{|\langle\langle UHU^{-1}\phi, \varphi, \rangle\rangle_\eta|} \right)^{1-\varpi(U^{-1})} \langle\langle UHU^{-1}\phi, \varphi, \rangle\rangle_\eta. \quad (3.13)$$

Since $|\langle\langle HU^{-1}\phi, U^{-1}\varphi, \rangle\rangle_\eta| = |\langle\langle UHU^{-1}\phi, \varphi, \rangle\rangle_\eta|$ and $\wp(U^{-1})\wp(U) = 1$ one concludes that

$$\langle\langle\phi, UHU^{-1}\varphi\rangle\rangle_\eta = \langle\langle UHU^{-1}\phi, \varphi, \rangle\rangle_\eta, \quad \forall \phi, \varphi \in \mathcal{D}_0 \quad (3.14)$$

or equivalently

$$\langle\phi, \eta UHU^{-1}\varphi\rangle = \langle\eta UHU^{-1}\phi, \varphi\rangle, \quad \forall \phi, \varphi \in \mathcal{D}_0 \quad (3.15)$$

where we used again that η is self-adjoint. The last equation implies that

$$\eta UHU^{-1} \subseteq (\eta UHU^{-1})^* = (UHU^{-1})^* \eta$$

proving that UHU^{-1} is η -symmetric. Moreover, the first inequality implies that ηUHU^{-1} is symmetric and therefore closable. Since η is bounded and invertible one immediately concludes that also UHU^{-1} is closable. ■

The last result deserves a comment. In general the product of two closed operators, even when it is densely defined, can fail to be closable (see e. g. [KR, Exercise 2.8.43]). However, this is not the case for UHU^{-1} under the assumption of Theorem 3.4. Unfortunately, the conditions of Theorem 3.4 does not seems to be sufficient to prove that the closure $\overline{UHU^{-1}}$ is η -self-adjoint. For that one needs conditions which assure

$$\eta \overline{UHU^{-1}} = \overline{\eta UHU^{-1}} = (UHU^{-1})^* \eta$$

where the first equality is for free due to the boundedness and invertibility of η , while the second is the difficult part to be proven. A sufficient condition is to prove that

$$\eta[\mathcal{D}((UHU^{-1})^*)] \subseteq \mathcal{D}(\overline{UHU^{-1}})$$

but in general this is difficult to check. A big simplification comes from assuming the boundedness of the linear symmetry U .

Corollary 3.5. *Let H be an η -self-adjoint operator with dense domain $\mathcal{D}(H)$. Let $U \in \text{QS}_\eta(\mathcal{H})$ be a bounded linear symmetry U of type (\wp, ϖ) . Then the operator UHU^{-1} is η -self-adjoint.*

Proof. First of all it is worth noticing that UHU^{-1} is automatically closed and defined on the dense domain $U[\mathcal{D}(H)]$ due to the boundedness and invertibility of U . Moreover, one has

$$(UHU^{-1})^* = (U^{-1})^* H^* U^* = (U^{-1})^* \eta H \eta U^* \quad (3.16)$$

where the last equality follows from the η -self-adjointness of H . By summarizing the results of Lemma 2.9 and Lemma 2.14 one obtains that each member of $U \in \text{QS}_\eta(\mathcal{H})$ meets

$$\eta U^* = \wp(U) U^{-1} \eta. \quad (3.17)$$

It is interesting to notice that the sign $\varpi(U)$ does not enter in the (3.17). Plugging in the (3.17) and the related relation

$$(U^{-1})^* \eta = \wp(U^{-1}) \eta U$$

into (3.16) and exploiting that $\wp(U)\wp(U^{-1}) = 1$ one finally gets

$$(UHU^{-1})^* = (\eta U) H (U^{-1} \eta) = \eta (UHU^{-1}) \eta.$$

The last equation is equivalent to $UHU^{-1} = (UHU^{-1})^\sharp$ and this proves the claim. ■

Corollary 3.5 holds true, in particular, for bounded η -self-adjoint operators. The space $\mathcal{H}_\eta(\mathcal{H})$ of bounded η -self-adjoint operators has the structure of a vector space over the field \mathbb{R} of the real number. Moreover, as consequence of the continuity of the adjoint, the space $\mathcal{H}_\eta(\mathcal{H})$ is closed with respect to the operator norm. Said differently, $\mathcal{H}_\eta(\mathcal{H})$ is an \mathbb{R} -Banach space.

Corollary 3.6. *The group $QS_\eta(\mathcal{H})$ acts by automorphisms on the \mathbb{R} -Banach space $H_\eta(\mathcal{H})$ through the adjoint action*

$$QS_\eta(\mathcal{H}) \ni U \mapsto \text{Ad}_U \in \text{Aut}(H_\eta(\mathcal{H}))$$

given by

$$\text{Ad}_U(H) := UHU^{-1}$$

for any $H \in H_\eta(\mathcal{H})$.

4. C-SYMMETRY FOR η -SELF-ADJOINT OPERATORS

The notion of C -symmetry for η -self-adjoint operators dates back to the pioneering works [BBJ1, BBJ2] in the context of the \mathcal{PT} -QM. For a recent review on the subject, see [Ben1]. From a more mathematical treatment of the notion of C -symmetry we refer to [Kuz, KSu] and references therein.

4.1. Definition and main properties. A C -symmetry in a space with indefinite metric η is an η -self-adjoint involution Ξ such that the product $\eta\Xi$ is strictly positive. More precisely, following e. g. [Kuz, Definition 3.1] one has:

Definition 4.1 (C -symmetry). *Let $(\mathcal{H}, \langle \langle \cdot, \cdot \rangle \rangle_\eta)$ be the indefinite metric space associated to a fundamental symmetry $\eta^* = \eta = \eta^{-1}$. A linear bounded operator $\Xi : \mathcal{H} \rightarrow \mathcal{H}$ which fulfills*

$$\Xi^2 = \mathbb{1}, \quad \eta\Xi > 0 \quad (4.1)$$

is called a C -symmetry. The symbol

$$C_\eta(\mathcal{H}) := \{\Xi \in \mathcal{B}_{\text{lin}}(\mathcal{H}) \mid \Xi \text{ meets (4.1)}\}$$

denotes the space of the C -symmetries for the space $(\mathcal{H}, \langle \langle \cdot, \cdot \rangle \rangle_\eta)$.

Note that the positivity condition automatically implies that $\eta\Xi = (\eta\Xi)^* = \Xi^*\eta$ namely

$$\Xi = \eta\Xi^*\eta = \Xi^\# \quad (4.2)$$

which means that Ξ is η -self-adjoint. Clearly $\Xi_0 = \eta$ is a (trivial) C -symmetry showing that $C_\eta(\mathcal{H})$ is not empty. Let us introduce the space

$$R_\eta(\mathcal{H}) := \{Q \in \mathcal{B}_{\text{lin}}(\mathcal{H}) \mid Q = Q^*, \quad Q\eta + \eta Q = 0\}$$

of bounded self-adjoint operators anti-commuting with the metric η . Observe that $Q \in R_\eta(\mathcal{H})$ implies

$$Q = -\eta Q^* \eta = -Q^\#,$$

namely Q is anti- η -self-adjoint. Both the spaces $C_\eta(\mathcal{H})$ and $R_\eta(\mathcal{H})$ can be topologized with the uniform topology inherited from the operator norm on $\mathcal{B}_{\text{lin}}(\mathcal{H})$. The next result describes the relation between these two spaces.

Lemma 4.2 ([KSu, Kuz]). *Let $(\mathcal{H}, \langle \langle \cdot, \cdot \rangle \rangle_\eta)$ be the indefinite metric space associated to a fundamental symmetry $\eta^* = \eta = \eta^{-1}$. The spaces $C_\eta(\mathcal{H})$ and $R_\eta(\mathcal{H})$ are closed with respect to the operator norm topology. Moreover, the maps*

$$\begin{aligned} C_\eta(\mathcal{H}) \ni \Xi &\xrightarrow{Q} Q(\Xi) := \log(\eta\Xi) \in R_\eta(\mathcal{H}) \\ R_\eta(\mathcal{H}) \ni Q &\xrightarrow{\Xi} \Xi(Q) := \eta e^Q \in C_\eta(\mathcal{H}) \end{aligned} \quad (4.3)$$

provide a homeomorphism $C_\eta(\mathcal{H}) \simeq R_\eta(\mathcal{H})$.

Proof. Let $\Xi \in C_\eta(\mathcal{H})$. Since $\eta\Xi$ is strictly positive and bounded the functional calculus allows to define the bounded self-adjoint operator $Q \equiv Q(\Xi) = \log(\eta\Xi)$. Then $e^Q = \eta\Xi$ and $e^{-Q} = (\eta\Xi)^{-1} = \Xi\eta$. Therefore $e^{-Q} = \eta e^Q \eta = e^{\eta Q \eta}$ where the last equality is easily justified by the unitarity of η and the boundedness of Q . This proves that $Q\eta + \eta Q = 0$, namely $Q \in R_\eta(\mathcal{H})$. On the other hand let $Q \in R_\eta(\mathcal{H})$ and define $\Xi \equiv \Xi(Q) = \eta e^Q$. It follows that $\Xi\eta = \eta e^Q \eta = e^{-Q}$ and in turn $\Xi^2 = \mathbb{1}$. Moreover, $\eta\Xi = e^Q$

is strictly positive by construction, hence $\Xi \in C_\eta(\mathcal{H})$. By observing that $Q(\Xi(Q)) = \log(\eta^2 e^Q) = Q$ and $\Xi(Q(\Xi)) = \eta e^{\log(\eta \Xi)} = \Xi$ it follows that $C_\eta(\mathcal{H})$ and $R_\eta(\mathcal{H})$ are in one-to-one correspondence.

Now the topology. Let us start proving that the map $Q \mapsto \Xi(Q)$ is continuous. As for this, let $\{Q_n\} \subset R_\eta(\mathcal{H})$ be a sequence converging to Q . Given a $\delta > 0$ there exists a $N \in \mathbb{N}$ such that

$$\|Q_n - Q\| \leq \|Q_n - Q\| < \delta, \quad \forall n > N.$$

Let $r_\delta := \max\{\|Q_1\|, \dots, \|Q_N\|, \|Q\| + \delta\}$. It follows that $\|Q\| < r_\delta$ and $\|Q_n\| < r_\delta$ for all $n \in \mathbb{N}$. If $\lambda \in \mathbb{C}$ such that $|\lambda| = r_\delta$ the resolvents $R_n(\lambda) := (Q_n - \lambda \mathbb{1})^{-1}$ and $R(\lambda) := (Q - \lambda \mathbb{1})^{-1}$ are well defined and $\|R_n(\lambda) - R(\lambda)\| \rightarrow 0$ in view of the identity

$$R_n(\lambda) = R(\lambda) (\mathbb{1} - (Q_n - Q) R(\lambda))^{-1}.$$

Let $C_\delta := \{\lambda \in \mathbb{C} \mid |\lambda| = r_\delta\}$. With the help of the holomorphic functional calculus one can write

$$\eta e^{Q_n} - \eta e^Q = \eta \frac{i}{2\pi} \oint_{C_\delta} e^\lambda (R_n(\lambda) - R(\lambda)) d\lambda$$

which, in view of the unitarity of η , leads to the norm estimate

$$\|\Xi(Q_n) - \Xi(Q)\| \leq K_{r_\delta} \|R_n(\lambda) - R(\lambda)\|$$

where the constant $K_{r_\delta} > 0$ depends on r_δ . This proves that the sequence $\Xi(Q_n)$ converges in norm towards $\Xi(Q)$, namely the maps $Q \mapsto \Xi(Q)$ is continuous. In order to prove the continuity of the inverse map $\Xi \mapsto Q(\Xi)$ let us observe first that $(\eta \Xi)^{-1} = \Xi \eta = \eta(\eta \Xi) \eta$. Since η is unitary it follows that $\sigma(\eta \Xi) = \sigma((\eta \Xi)^{-1})$ and in view of the *spectral mapping theorem* one has that $\lambda \in \sigma(\eta \Xi)$ if and only if $\lambda^{-1} \in \sigma(\eta \Xi)$. This implies that $\|\Xi\| = \|\eta \Xi\| = \max \sigma(\eta \Xi) \geq 1$ and $\sigma(\eta \Xi) \subseteq [\|\Xi\|^{-1}, \|\Xi\|]$. Let $\{\Xi_n\} \subset C_\eta(\mathcal{H})$ be a sequence converging to Ξ . Given a $\delta > 0$ there exists a $N \in \mathbb{N}$ such that

$$\|\Xi_n - \Xi\| \leq \|\Xi_n - \Xi\| < \delta, \quad \forall n > N.$$

Let $b_\delta := \max\{\|\Xi_1\|, \dots, \|\Xi_N\|, \|\Xi\| + \delta\}$ and $I_\delta := (b_\delta^{-1}, b_\delta)$. It follows that $\|\Xi\| < b_\delta$ and $\|\Xi_n\| < b_\delta$ for all $n \in \mathbb{N}$ and in turn $\sigma(\eta \Xi) \subset I_\delta$ and $\sigma(\eta \Xi_n) \subset I_\delta$ for all $n \in \mathbb{N}$. Let $\tilde{C}_\delta \subset \mathbb{C}$ any closed Jordan curve that crosses the real axis only in b_δ^{-1} and b_δ . If $\lambda \in \tilde{C}_\delta$ the resolvents $\tilde{R}_n(\lambda) := (\eta \Xi_n - \lambda \mathbb{1})^{-1}$ and $\tilde{R}(\lambda) := (\eta \Xi - \lambda \mathbb{1})^{-1}$ are well defined and $\|\tilde{R}_n(\lambda) - \tilde{R}(\lambda)\| \rightarrow 0$ in view of the identity

$$\tilde{R}_n(\lambda) = \tilde{R}(\lambda) (\mathbb{1} - \eta(\Xi_n - \Xi) \tilde{R}(\lambda))^{-1}.$$

With the help of the holomorphic functional calculus one can write

$$\log(\eta \Xi_n) - \log(\eta \Xi) = \frac{i}{2\pi} \oint_{\tilde{C}_\delta} \log(\lambda) (\tilde{R}_n(\lambda) - \tilde{R}(\lambda)) d\lambda$$

where the logarithm is unambiguously defined since the curve \tilde{C}_δ does not enclose the origin of \mathbb{C} . The last equality provides

$$\|Q(\Xi_n) - Q(\Xi)\| \leq \tilde{K}_{b_\delta} \|\tilde{R}_n(\lambda) - \tilde{R}(\lambda)\|$$

where the constant $\tilde{K}_{b_\delta} > 0$ depends on b_δ . This proves that the sequence $Q(\Xi_n)$ converges in norm towards $Q(\Xi)$, namely also the maps $\Xi \mapsto Q(\Xi)$ is continuous. In summary we proved that the spaces $C_\eta(\mathcal{H})$ and $R_\eta(\mathcal{H})$ are homeomorphic.

Since the multiplication and the adjoint are continuous operations with respect to the operator norm it follows that $R_\eta(\mathcal{H})$ is a *closed* subset of $\mathcal{B}_{\text{lin}}(\mathcal{H})$. Since $C_\eta(\mathcal{H})$ is homeomorphic to $R_\eta(\mathcal{H})$ it follows that also $C_\eta(\mathcal{H})$ is closed. \blacksquare

Lemma 4.2 states that any C -symmetry Ξ can be uniquely represented by a self-adjoint and anti- η -self-adjoint operator Q . Also the set $R_\eta(\mathcal{H})$ has a *trivial* element $Q_0 := 0$. In fact this is the element associated to $\Xi_0 = \eta$ through the homeomorphism described in Lemma 4.2.

Remark 4.3 (Generalized C -symmetries). It is possible to generalize Definition 4.1 by relaxing the boundedness condition for Ξ , see e. g. [Kuz, Section 4]. Indeed, it is possible to construct unbounded involutions in infinite dimensional normed linear spaces. For instance given an orthonormal basis $\{\psi_n\}$ of the Hilbert space \mathcal{H} the operator Ξ defined by $\Xi\psi_n := 2n\psi_1 - \psi_n$ is unbounded and verifies $\Xi^2 = \mathbb{1}$ (cf. [Buc]). Moreover, also in a finite dimensional space there are C -symmetries of arbitrarily large norm (cf. Lemma B.3). Anyway, also in the unbounded case the representation described in Lemma 4.2 continues to be valid [AK2, Theorem 6.2.3]. However, in this work we will not be interested in this kind of generalization. In fact [KSu, Proposition 3.2] shows that in the context of gapped η -self-adjoint operator only the notion of bounded C -symmetry is relevant. For this reason we will focus our attention on Definition 4.1. ◀

The following result will be relevant in the next Section.

Lemma 4.4. *Let $(\mathcal{H}, \langle\langle \cdot, \cdot \rangle\rangle_\eta)$ be the indefinite metric space associated to a fundamental symmetry $\eta^* = \eta = \eta^{-1}$. Let $\Xi \in \mathcal{C}_\eta(\mathcal{H})$ be a C -symmetry and $Q = Q(\Xi) \in \mathcal{R}_\eta(\mathcal{H})$ the related self-adjoint and anti- η -self-adjoint operator associated to Ξ through the bijection described in Lemma 4.2. The operator*

$$G_\Xi := e^{\frac{1}{2}Q} = e^{\frac{1}{2}\log(\eta\Xi)} = \sqrt{\eta\Xi}$$

is η -unitary, self-adjoint and $G_\Xi \Xi G_\Xi^{-1} = \eta$.

Proof. The operator G_Ξ is defined via spectral calculus. The representation $G_\Xi = \sqrt{\eta\Xi}$ makes evident that $G_\Xi = G_\Xi^*$ is a bounded self-adjoint operator. Moreover, G_Ξ is invertible with inverse $G_\Xi^{-1} = e^{-\frac{1}{2}Q} = \eta G_\Xi \eta = G_\Xi^\sharp$. The last expression shows that $G_\Xi \in \mathcal{U}_\eta(\mathcal{H})$ is a bounded η -unitary operator. Finally the computation

$$G_\Xi^{-1} \eta G_\Xi = e^{-\frac{1}{2}Q} \eta e^{\frac{1}{2}Q} = \eta e^Q = \Xi$$

completes the proof. ■

Lemma 4.4 states that any C -symmetry Ξ can be reduced to the trivial C -symmetry $\Xi_0 = \eta$ by means of a transformation G_Ξ which preserves the geometry of the indefinite metric space $(\mathcal{H}, \langle\langle \cdot, \cdot \rangle\rangle_\eta)$.

Remark 4.5. If one renounces to the condition that η is an involution as in the original Definition 2.1, then the second condition in (4.1) for the definition of a C -symmetry must be replaced by $\frac{\eta}{|\eta|}\Xi > 0$. With this modification also the results proved in Lemma 4.2 and Lemma 4.4 still work. ◀

4.2. Operators with C -symmetry. Let us focus on the interplay between the notion of η -self-adjointness and the existence of C -symmetries. The following definition is borrowed from [AK2, Definition 6.3.3]:

Definition 4.6 (Operators with C -symmetry). *Let $(\mathcal{H}, \langle\langle \cdot, \cdot \rangle\rangle_\eta)$ be the indefinite metric space associated to a fundamental symmetry $\eta^* = \eta = \eta^{-1}$. A densely defined operator A possesses a C -symmetry if there exists an element $\Xi \in \mathcal{C}_\eta(\mathcal{H})$ such that $\Xi A = A \Xi$.*

In this work we are mainly interested in η -self-adjoint operators with C -symmetry. The following result is known in the literature [AK2, Theorem 6.3.4]:

Proposition 4.7. *Let $(\mathcal{H}, \langle\langle \cdot, \cdot \rangle\rangle_\eta)$ be the indefinite metric space associated to a fundamental symmetry $\eta^* = \eta = \eta^{-1}$. Let H be a η -self adjoint operator with dense domain $\mathcal{D}(H)$ which admits a C -symmetry $\Xi \in \mathcal{C}_\eta(\mathcal{H})$. Let G_Ξ be the associated η -unitary operator described in Lemma 4.4. Then the operator*

$$\tilde{H} := G_\Xi H G_\Xi^{-1}$$

is self-adjoint with domain $\mathcal{D}(\tilde{H}) := G_\Xi[\mathcal{D}(H)]$ and commutes with the fundamental symmetry η .

Proof. Let us start by observing that in the case of an unbounded operator H the condition $\Xi H = H \Xi$ and the invertibility of Ξ imply that $\Xi[\mathcal{D}(H)] = \mathcal{D}(H)$. The operator H_Ξ is well defined on the dense domain $\mathcal{D}(\tilde{H})$. From the identity $\Xi = G_\Xi^{-1} \eta G_\Xi$ derived in the proof of Lemma 4.4 it follows that

$$G_\Xi^{-1} \eta G_\Xi H = H G_\Xi^{-1} \eta G_\Xi \quad \Leftrightarrow \quad \eta \tilde{H} = \tilde{H} \eta$$

namely \tilde{H} commutes with η . On the other hand G_Ξ is η -unitary and therefore Corollary 3.5 ensures that \tilde{H} is η -self-adjoint. At level of domains one has that

$$\mathcal{D}(\tilde{H}^*) = \eta[\mathcal{D}(\tilde{H})] = (\eta G_\Xi)[\mathcal{D}(H)] = (\eta G_\Xi \Xi)[\mathcal{D}(H)] = G_\Xi[\mathcal{D}(H)] = \mathcal{D}(\tilde{H}) .$$

Moreover, since \tilde{H} is η -self-adjoint and commutes with η it follows that

$$\tilde{H}^* = \eta \tilde{H} \eta = \tilde{H} ,$$

namely \tilde{H} is self-adjoint. ■

Proposition 4.7 states that an η -self-adjoint operator with a C -symmetry is similar to a self-adjoint operator, or *dynamically stable* by using the jargon of Definition 1.1. Also the converse is true:

Theorem 4.8 ([AK2, Theorem 6.3.4]). *An η -self-adjoint operator is dynamically stable if and only if it has a C -symmetry.*

The next result shows that dynamically stable η -self-adjoint operators possess the characteristics for the definition of a quantum dynamics as discussed in Section 1.1.

Corollary 4.9 (Induced functional calculus). *Let $(\mathcal{H}, \langle \cdot, \cdot \rangle_\eta)$ be the indefinite metric space associated to a fundamental symmetry $\eta^* = \eta = \eta^{-1}$. Let H be an η -self adjoint operator with dense domain $\mathcal{D}(H)$ which admits a C -symmetry $\Xi \in \mathcal{C}_\eta(\mathcal{H})$.*

- (1) *The spectrum of H is real, $\sigma(H) \subseteq \mathbb{R}$.*
- (2) *There is a $*$ -algebra homomorphism $\varphi : L^\infty(\mathbb{R}) \rightarrow \mathcal{B}_{\text{lin}}(\mathcal{H})$ defined by*

$$\varphi(f) \mapsto f(H) := G_\Xi^{-1} f(\tilde{H}) G_\Xi .$$

- (3) *The map $\mathbb{R} \ni t \mapsto V_t := e^{-itH}$ provides a group of η -unitary operators that solves the Schrödinger equation*

$$i \frac{\partial}{\partial t} \psi = H \psi , \quad \psi \in \mathcal{D}(H) .$$

Proof (sketch of). Item (1) follows from the fact that H and \tilde{H} are intertwined by the invertible operator G_Ξ . This implies that $\sigma(H) = \sigma(\tilde{H}) \subseteq \mathbb{R}$ where the last inclusion follows since \tilde{H} is self-adjoint. Item (2) follows from the functional calculus for the self-adjoint operator \tilde{H} . As for item (3) one has that $V_t = G_\Xi^{-1} e^{-it\tilde{H}} G_\Xi$ and $e^{-it\tilde{H}}$ solves the Schrödinger equation for \tilde{H} , namely

$$i \frac{\partial}{\partial t} (e^{-it\tilde{H}} G_\Xi \psi) = \tilde{H} (e^{-it\tilde{H}} G_\Xi \psi) , \quad \psi \in \mathcal{D}(H) .$$

A multiplication on the left by G_Ξ^{-1} provides

$$i \frac{\partial}{\partial t} (V_t \psi) = H (V_t \psi) , \quad \psi \in \mathcal{D}(H) .$$

The η -unitarity of V_t follows since G_Ξ is η -unitarity and $e^{-it\tilde{H}}$ commutes with η . ■

Let us focus on the bounded case. Given a C -symmetry $\Xi \in \mathcal{C}_\eta(\mathcal{H})$ let

$$\mathcal{H}_{\eta, \Xi}(\mathcal{H}) := \{ H \in \mathcal{H}_\eta(\mathcal{H}) \mid \Xi H = H \Xi \}$$

be the set of bounded η -self-adjoint operators that admit Ξ as C -symmetry. More in general

$$\mathcal{CH}_\eta(\mathcal{H}) := \bigsqcup_{\Xi \in \mathcal{C}_\eta(\mathcal{H})} \mathcal{H}_{\eta, \Xi}(\mathcal{H})$$

is the set of all η -self-adjoint operators which admit a C -symmetry. The case related to the trivial C -symmetry $\Xi_0 := \eta$ is of particular relevance. One has that

$$\mathcal{H}_{\eta, 0}(\mathcal{H}) := \mathcal{H}_{\eta, \Xi_0}(\mathcal{H}) = \{ H \in \mathcal{H}_\eta(\mathcal{H}) \mid H = H^* \}$$

is the space of self-adjoint and η -self-adjoint bounded operators. As immediate consequence of Lemma 4.4 and Proposition 4.7 one has that:

Corollary 4.10. *For any $\Xi \in \mathcal{C}_\eta(\mathcal{H})$ there is an η -unitary equivalence*

$$\mathcal{H}_{\eta,\Xi}(\mathcal{H}) \simeq \mathcal{H}_{\eta,0}(\mathcal{H})$$

induced by the invertible map $\text{Ad}_\Xi : \mathcal{H}_{\eta,\Xi}(\mathcal{H}) \rightarrow \mathcal{H}_{\eta,0}(\mathcal{H})$ given by

$$\text{Ad}_\Xi(H) := G_\Xi H G_\Xi^{-1} = \tilde{H}.$$

The interplay between the η -quantum symmetries $\text{QS}_\eta(\mathcal{H})$ described in Section 3.2 and the \mathcal{C} -symmetries is provided by the following result.

Proposition 4.11. *Let $(\mathcal{H}, \langle \langle \cdot, \cdot \rangle \rangle_\eta)$ be the indefinite metric space associated to a fundamental symmetry $\eta^* = \eta = \eta^{-1}$, and $U \in \text{QS}_\eta(\mathcal{H})$ an η -quantum symmetry of type (\wp, ϖ) .*

- (1) *The map U induces an invertible transformation $\text{Ad}_U^\wp : \mathcal{C}_\eta(\mathcal{H}) \rightarrow \mathcal{C}_\eta(\mathcal{H})$ of the space of \mathcal{C} -symmetries in itself given by*

$$\text{Ad}_U^\wp(\Xi) := \wp(U) U \Xi U^{-1}.$$

Moreover, if $H \in \mathcal{H}_{\eta,\Xi}(\mathcal{H})$ then $\text{Ad}_U(H) \in \mathcal{H}_{\eta,\text{Ad}_U^\wp(\Xi)}(\mathcal{H})$.

- (2) *If $\text{Ad}_U^\wp(\Xi) = \Xi$ then the diagram*

$$\begin{array}{ccc} \mathcal{H}_{\eta,\Xi}(\mathcal{H}) & \xrightarrow{\text{Ad}_\Xi} & \mathcal{H}_{\eta,0}(\mathcal{H}) \\ \text{Ad}_U \downarrow & & \downarrow \text{Ad}_{\tilde{U}} \\ \mathcal{H}_{\eta,\Xi}(\mathcal{H}) & \xrightarrow{\text{Ad}_\Xi} & \mathcal{H}_{\eta,0}(\mathcal{H}) \end{array}$$

is commutative, with $\tilde{U} := G_\Xi U G_\Xi^{-1}$. Moreover, \tilde{U} is of type (\wp, ϖ) .

Proof. (1) The operator $\Xi' := \text{Ad}_U^\wp(\Xi)$ is clearly involutive. Moreover, for every $\varphi \neq 0$ one has

$$\begin{aligned} \langle \varphi, \eta \Xi' \varphi \rangle &= \wp(U) \langle \langle \varphi, U \Xi U^{-1} \varphi \rangle \rangle_\eta \\ &= \wp(U) \wp(U^{-1}) \left| \langle \langle U^{-1} \varphi, \Xi U^{-1} \varphi \rangle \rangle_\eta \right|^{1-\varpi(U^{-1})} \langle \langle U^{-1} \varphi, \Xi U^{-1} \varphi \rangle \rangle_\eta^{\varpi(U^{-1})} \\ &= \langle U^{-1} \varphi, \eta \Xi U^{-1} \varphi \rangle > 0 \end{aligned}$$

since $\wp(U) \wp(U^{-1}) = 1$ and $\eta \Xi > 0$. This implies that also $\eta \Xi' > 0$ is strictly positive. Finally, a straightforward computation shows that $[H, \Xi] = 0$ if and only if $[\text{Ad}_U(H), \Xi'] = 0$ and $\text{Ad}_U(H)$ is η -self-adjoint in view of Corollary 3.5. (2) The commutativity of the diagram follows from a direct computation. Since G_Ξ is η -unitary as proved in Lemma 4.4, a direct computation shows that the type of \tilde{U} only depend on the type of U . ■

4.3. The Hilbert bundle picture. The various notions introduced in Sections 2, 3 and 4 can be extended to the case of *Hilbert bundles* and complex vector bundles. The motivations of this generalization will be sketched in Remark 4.16 below. For the general theory of Hilbert bundles we refer to [DD, Dup] and references therein. For the general of vector bundles we refer to the two classic monographs [Ati2, Hus].

Let us recall that a Hilbert bundle $\pi : \mathcal{E} \rightarrow X$ is a locally trivial bundle over X such that each fiber $\mathcal{H}_x := \pi^{-1}(x)$ is a separable Hilbert space with scalar product $\langle \cdot, \cdot \rangle_x$ and the projection π is an open surjective map. When the fibers of the Hilbert bundle $\pi : \mathcal{E} \rightarrow X$ are infinite dimensional then the transition functions of \mathcal{E} take value on the structure group which is the infinite dimensional unitary group $\text{U}(\mathcal{H})$ topologized with the strong topology⁹. For the aims of this work it is sufficient to assume that:

⁹In [FM] the structure group is topologized by the compact-open topology in the sense of Atiyah and Segal. However, as discussed in Remark 3.1, the strong topology should be sufficient for all needs. The strong topology has been used also in [DD, Dup, Chapitre II.10]

Assumption 4.12. *The base space X is Hausdorff and compact and is endowed with a regular Borel measure μ .*

Under Assumption 4.12, if $\mathcal{H}_x \simeq \mathbb{C}^m$ for all $x \in X$ one has that $\pi : \mathcal{E} \rightarrow X$ is (equivalent to) a rank m complex vector bundle endowed with a Hermitian metric. In this case the structure group is the unitary group $U(\mathbb{C}^m)$ with its natural topology.

A morphism between two Hilbert bundles $\pi : \mathcal{E} \rightarrow X$ and $\pi' : \mathcal{E}' \rightarrow X$ is a continuous map $f : \mathcal{E} \rightarrow \mathcal{E}'$ which preserves the fiber structure in the sense that $\pi = \pi' \circ f$ and the linear structure of each fiber in the sense that $f|_{\mathcal{H}_x} : \mathcal{H}_x \rightarrow \mathcal{H}'_x$ is linear or anti-linear. Let $\text{End}(\mathcal{E})$ be the space of the *endomorphisms* of the Hilbert bundle $\pi : \mathcal{E} \rightarrow X$. This means that elements $A \in \text{End}(\mathcal{E})$ are continuous maps of the total space \mathcal{E} such that $\pi = \pi \circ A$ and that $A_x := A|_{\mathcal{H}_x} \in \mathcal{B}(\mathcal{H}_x)$ is fiberwise a linear or anti-linear bounded operator. Diagrammatically one has that elements $A \in \text{End}(\mathcal{E})$ make the diagrams

$$\begin{array}{ccc} \mathcal{E} & \xrightarrow{A} & \mathcal{E} \\ & \searrow \pi & \swarrow \pi \\ & X & \end{array}$$

commutative. The subset of the *fiber-linear* endomorphisms will be denoted with $\text{End}_{\text{lin}}(\mathcal{E})$. The adjoint of $A \in \text{End}(\mathcal{E})$ is the endomorphism $A^* \in \text{End}(\mathcal{E})$ defined fiberwise by $A^*|_{\mathcal{H}_x} = A_x^*$. It is self-adjoint if $A = A^*$ and it is an involution if $A = A^{-1}$.

Definition 4.13 (Krein bundle). *An indefinite metric structure on the Hilbert bundle $\pi : \mathcal{E} \rightarrow X$ is given by a self-adjoint involution $\eta \in \text{End}_{\text{lin}}(\mathcal{E})$, such that each restriction $\eta_x := \eta|_{\mathcal{H}_x}$ is a fundamental symmetry on the fiber \mathcal{H}_x , i. e. $\eta_x^* = \eta_x = \eta_x^{-1}$. If for each $x \in X$ the metric η_x provides a Krein space structure in the fiber \mathcal{H}_x (in the sense of Definition 2.3) then the pair (\mathcal{E}, η) is called a Krein bundle. In the case the fibers of \mathcal{E} are of finite even dimension then (\mathcal{E}, η) is called a Krein vector bundle.*

A morphism between two Krein bundles (\mathcal{E}, η) and (\mathcal{E}', η') is a Hilbert bundle morphism $f : \mathcal{E} \rightarrow \mathcal{E}'$ such that $f \circ \eta = \eta' \circ f$. The notion of η -self-adjointness (cf. Definition 2.4) and C -symmetry (cf. Definition 4.1) have a natural extension to the case of Krein bundles. Let (\mathcal{E}, η) be a Krein bundle. An element $H \in \text{End}_{\text{lin}}(\mathcal{E})$ is η -self-adjoint if $H = H^\sharp := \eta \circ H^* \circ \eta$. The subset of the η -self-adjoint endomorphisms will be denoted by $H_\eta(\mathcal{E}) \subset \text{End}_{\text{lin}}(\mathcal{E})$. A C -symmetry is an element $\Xi \in \text{End}_{\text{lin}}(\mathcal{E})$ such that $\Xi = \Xi^{-1}$ and the composition $\eta \circ \Xi$ is strictly positive when restricted to each fiber, i. e. $(\eta \circ \Xi)|_{\mathcal{H}_x} = \eta_x \Xi_x > 0$. The subset of the C -symmetries will be denoted by $C_\eta(\mathcal{E}) \subset \text{End}_{\text{lin}}(\mathcal{E})$. In accordance with Definition 4.1 one says that $H \in H_\eta(\mathcal{E})$ admits a C -symmetry if there exists an element $\Xi \in C_\eta(\mathcal{E})$ such that $H \circ \Xi = \Xi \circ H$. The subset of the η -self-adjoint endomorphisms with C -symmetry Ξ will be denoted with $H_{\eta, \Xi}(\mathcal{E}) \subset H_\eta(\mathcal{E})$ while $CH_\eta(\mathcal{E}) \subset H_\eta(\mathcal{E})$ will denote the subset of the C -symmetric η -self-adjoint endomorphisms. Proposition 4.7 and Corollary 4.9 (and also the more general Theorem (4.8)) can be reproved in the setting of Krein bundles with a fiberwise version of the same proofs. In particular, let $U_\eta(\mathcal{E}) \subset \text{End}_{\text{lin}}(\mathcal{E})$ be the subset of the η -unitary (fiber-linear) endomorphisms, i. e. $U \in U_\eta(\mathcal{E})$ if and only if $U^{-1} = U^\sharp$. To any C -symmetry $\Xi \in C_\eta(\mathcal{E})$ it can be associated a $G_\Xi \in U_\eta(\mathcal{E})$ defined fiberwise by $G_\Xi|_{\mathcal{H}_x} := \sqrt{\eta_x \Xi_x}$. As in Lemma 4.4 one can prove that $G_\Xi = G_\Xi^*$ and $G_\Xi \Xi G_\Xi^{-1} = \eta$. Moreover, a fiberwise adaptation of Corollary 4.10 shows that for any $H \in H_{\eta, \Xi}(\mathcal{E})$ the transformed endomorphism $\tilde{H} := G_\Xi \circ H \circ G_\Xi^{-1}$ is self-adjoint $\tilde{H} = \tilde{H}^*$ and commute with the metric structure, $\tilde{H} \circ \eta = \eta \circ \tilde{H}$. The set of self-adjoint endomorphisms commuting with η will be denoted with $H_{\eta, 0}(\mathcal{E})$. In view of the equivalence induced by G_Ξ one concludes that the classification of the C -symmetric η -self-adjoint endomorphisms is equivalent to the classification of $H_{\eta, 0}(\mathcal{E})$. This will be the main task of Section 5. Given the Krein bundle (\mathcal{E}, η) and using a local version of the construction of Note 4 and Note 5 one can define an η -reflecting endomorphism $R^* = R^{-1} = -R^\sharp$ and an anti-linear complex structure $C^* = C^\sharp = C = C^{-1}$ such that $R \circ C = C \circ R$. With these

endomorphisms and by adapting the prescriptions (2.11), (2.15) and (2.16) one can define the group

$$\text{QS}_\eta(\mathcal{E}) := \text{U}_\eta(\mathcal{E}) \sqcup \text{AU}_\eta(\mathcal{E}) \sqcup \text{PU}_\eta(\mathcal{E}) \sqcup \text{PAU}_\eta(\mathcal{E}) \subset \text{End}(\mathcal{E}).$$

We refer to $\text{QS}_\eta(\mathcal{E})$ as the group of the *fiber-preserving* η -quantum symmetries of the Krein bundle (\mathcal{E}, η) . Let us remark that any $U \in \text{QS}_\eta(\mathcal{E})$ meets the strong condition $\pi = \pi \circ U$. Indeed, by relaxing the latter condition a more general notion of symmetry can be introduced (cf. Definition 4.15). Anyway, $U \in \text{QS}_\eta(\mathcal{E})$ is a *good* symmetry for $H \in \mathcal{H}_{\eta, \Xi}(\mathcal{E})$ if: (i) $U \circ H = c(U)H \circ U$, and (ii) $U \circ \Xi = \wp(U)\Xi \circ U$. Here the signs $\wp(U), \varpi(U) \in \{-1, +1\}$ describe the nature of the symmetry U according to the convention of Table 1.3. The sign $c(U) \in \{-1, +1\}$ says if U is a proper or an improper symmetry of H according to the terminology introduced in Section 1.2. The condition (ii) expresses the *compatibility* of the symmetry U with the C -symmetry property of H . The fiberwise version of the proof of Theorem 1.4 shows that $\tilde{U} := G_\Xi \circ U \circ G_\Xi^{-1} \in \text{QS}_\eta(\mathcal{E})$ has the same nature of U (i.e. $\wp(U) = \wp(\tilde{U})$ and $\varpi(U) = \varpi(\tilde{U})$) and is *unitary* $\tilde{U}^{-1} = \tilde{U}^*$. In particular the later condition implies $\tilde{U} \circ \eta = \wp(\tilde{U})\tilde{U} \circ \eta$. Moreover, $\tilde{U} \circ \tilde{H} = c(\tilde{U})\tilde{H} \circ \tilde{U}$ with $c(\tilde{U}) = c(U)$. In conclusion we can sum up as follows:

Observation A. *Once the problem of the classification of the C -symmetric η -self-adjoint endomorphisms is reduced to the problem of the classification of $\mathcal{H}_{\eta, 0}(\mathcal{E})$ it turns out that the only relevant (fiber-preserving) η -quantum symmetries are induced by unitary or anti-unitary endomorphisms which commute or anti-commute with the metric η .*

We will see below that the framework described in Observation A deserves a further, non-trivial generalization (cf. Observation B). For the moment, let us introduce an extra structure that will be relevant for the purposes of Section 5: The action of *Clifford algebras* on Hilbert and Krein bundles. For a detailed introduction to Clifford algebras we refer to [Kar, Chapter III, Section 3] or [Le]. Let $C\ell^{r,s}$ be the Clifford algebra generated over \mathbb{R} by a collection of symbols e_1, \dots, e_{r+s} subjected to the following relations

$$\frac{e_i e_j + e_j e_i}{2} := \begin{cases} 0 & \text{if } i \neq j \\ +1 & \text{if } i = j \in \{1, \dots, r\} \\ -1 & \text{if } i = j \in \{r+1, \dots, r+s\}. \end{cases}$$

These Clifford algebras have been completely classified.

Definition 4.14 (Clifford action on Hilbert and Krein bundles). *Let $\pi : \mathcal{E} \rightarrow X$ be a Hilbert bundle and $C\ell^{r,s}$ a Clifford algebra. A (unitary) Clifford action on \mathcal{E} is given by an \mathbb{R} -algebra homomorphism $\gamma : C\ell^{r,s} \rightarrow \text{End}_{\text{lin}}(\mathcal{E})$ such that the representatives of the generators $\gamma_i := \gamma(e_i)$ are unitary, i.e. $\gamma_i^{-1} = \gamma_i^*$ for all $i = 1, \dots, r+s$. The pair (\mathcal{E}, γ) will be called a Clifford-Hilbert bundle of type (r, s) . In the case of a Krein bundle (\mathcal{E}, η) a compatible Clifford action is subjected to the extra condition $\gamma(a) \circ \eta = \eta \circ \gamma(a)$ for all $a \in C\ell^{r,s}$. The triple $(\mathcal{E}, \eta, \gamma)$ given by the Krein bundle (\mathcal{E}, η) and the compatible Clifford action γ will be called a Clifford-Krein bundle.*

A morphism between two Clifford-Hilbert bundles (resp. Clifford-Krein bundles) of the same type (\mathcal{E}, γ) and (\mathcal{E}', γ') is a Hilbert bundle (resp. Krein bundles) morphism $f : \mathcal{E} \rightarrow \mathcal{E}'$ such that $f \circ \gamma(a) = \gamma'(a) \circ f$ for all $a \in C\ell^{r,s}$. Definition 4.14 also works in the vector bundle (i.e. finite rank) case.

Up to now, all the concepts introduced above are a natural generalization of the theory and the results developed in Sections 2, 3 and 4. More specifically, the whole content of this Section reduces to the content of the previous Sections when the base space of the bundle reduces to a point $X = \{*\}$. However, the bundle picture provides more freedom for the construction of “interesting” symmetries. Let us recall that a *section* of a Hilbert bundle $\pi : \mathcal{E} \rightarrow X$ is a map $s : X \rightarrow \mathcal{E}$ such that $(\pi \circ s)(x) = x$ for all $x \in X$. The space of sections is usually denoted with $\Gamma(\mathcal{E})$. With the help of the measure structure assumed in Assumption 4.12 one can define a Hilbert structure on $\Gamma(\mathcal{E})$ by means of the scalar product

$$\langle s_1, s_2 \rangle_{\mathcal{E}} := \int_X d\mu(x) \langle s_1(x), s_2(x) \rangle_x. \quad (4.4)$$

The completion of $\Gamma(\mathcal{E})$ with respect the topology induced by the Hilbert structure (4.4) defines the Hilbert space of the L^2 -sections denoted with $\mathfrak{H}_{\mathcal{E}} := L^2(\mathcal{E}, \mu)$. Let $\mathcal{B}(\mathfrak{H}_{\mathcal{E}})$ be the space of the linear or anti-linear bounded operators acting on $\mathfrak{H}_{\mathcal{E}}$ and $\mathcal{U}(\mathfrak{H}_{\mathcal{E}})$ and $\mathcal{AU}(\mathfrak{H}_{\mathcal{E}})$ the subsets of the unitary and anti-unitary operators. Clearly the fiber-preserving endomorphisms $\text{End}(\mathcal{E})$ can be identified with operators in $\mathcal{B}(\mathfrak{H}_{\mathcal{E}})$. Similarly, one has (up to an identification) the inclusions $\mathcal{U}(\mathcal{E}) \subset \mathcal{U}(\mathfrak{H}_{\mathcal{E}})$ and $\mathcal{AU}(\mathcal{E}) \subset \mathcal{AU}(\mathfrak{H}_{\mathcal{E}})$. However not all the operator in $\mathcal{B}(\mathfrak{H}_{\mathcal{E}})$ come from fiber-preserving endomorphisms. An interesting class of operator is generated by the *covariant endomorphisms* of \mathcal{E} .

Definition 4.15 (Covariant endomorphism). *Let $\pi : \mathcal{E} \rightarrow X$ be a Hilbert bundle. A covariant endomorphism of \mathcal{E} is a pair (A, g_A) with $A : \mathcal{E} \rightarrow \mathcal{E}$ is a continuous map on the total space, $g_A : X \rightarrow X$ is a homeomorphism on the base space, such that the following diagram commutes*

$$\begin{array}{ccc} \mathcal{E} & \xrightarrow{A} & \mathcal{E} \\ \pi \downarrow & & \downarrow \pi \\ X & \xrightarrow{g_A} & X \end{array}$$

The covariant endomorphism (A, g_A) is linear if $A(\lambda p) = \lambda A(p)$ for all $\lambda \in \mathbb{C}$ and $p \in \mathcal{E}$. It is called anti-linear when $A(\lambda p) = \bar{\lambda} A(p)$. The space of linear or anti-linear covariant endomorphisms will be denoted with $\text{Cov}(\mathcal{E})$ while $\text{Cov}_{\text{lin}}(\mathcal{E})$ will be used for the subspace of the linear ones.

Clearly from Definition 4.15 one deduces the following inclusions $\text{End}(\mathcal{E}) \subset \text{Cov}(\mathcal{E})$ and $\text{Cov}_{\text{lin}}(\mathcal{E}) \subset \text{Cov}(\mathcal{E})$ obtained when $g_A = \text{Id}_X$. An element $(A, g_A) \in \text{Cov}(\mathcal{E})$ induces a map $\hat{A} : \Gamma(\mathcal{E}) \rightarrow \Gamma(\mathcal{E})$ defined by

$$\hat{A}s := A \circ s \circ g_A^{-1}.$$

The map \hat{A} is linear or anti-linear according to the nature of A and by density it defines an element $\hat{A} \in \mathcal{B}(\mathfrak{H}_{\mathcal{E}})$. The covariant endomorphism (A, g_A) is called unitary (resp. anti-unitary) if the associated operator \hat{A} is unitary (resp. anti-unitary). It is an (even or odd) involution if $\hat{A}^2 = \varepsilon \mathbb{1}$ with $\varepsilon = \pm 1$. In this case it is necessary that $g_A^2 = \text{Id}_X$, namely an *involution* of the base space.

In the case of a Krein bundle (\mathcal{E}, η) the map $\eta \in \text{End}_{\text{lin}}(\mathcal{E})$ induces a Krein structure on the Hilbert space $\mathfrak{H}_{\mathcal{E}}$ by the indefinite inner product

$$\langle\langle s_1, s_2 \rangle\rangle_{\mathcal{E}, \eta} := \int_X d\mu(x) \langle s_1(x), \eta_x s_2(x) \rangle_x. \quad (4.5)$$

This allows to define the full group of the (non necessarily fiber-preserving) η -quantum symmetries

$$\text{QS}_{\eta}(\mathfrak{H}_{\mathcal{E}}) := \mathcal{U}_{\eta}(\mathfrak{H}_{\mathcal{E}}) \sqcup \mathcal{AU}_{\eta}(\mathfrak{H}_{\mathcal{E}}) \sqcup \mathcal{PU}_{\eta}(\mathfrak{H}_{\mathcal{E}}) \sqcup \mathcal{PAU}_{\eta}(\mathfrak{H}_{\mathcal{E}}) \subset \mathcal{B}(\mathfrak{H}_{\mathcal{E}}).$$

The main goal of this work is to provide a topological classification of the space $\mathcal{H}_{\eta, \Xi}(\mathcal{E})$ of dynamically stable (fiber-preserving) operators under the possible action of compatible η -quantum symmetries. However the group $\text{QS}_{\eta}(\mathfrak{H}_{\mathcal{E}})$ is too vast and there are physical reasons to restrict the analysis only to *covariant* η -quantum symmetries $\text{QS}_{\eta}(\mathfrak{H}_{\mathcal{E}}) \cap \text{Cov}(\mathcal{E})$. Clearly $\text{QS}_{\eta}(\mathcal{E}) \subset \text{QS}_{\eta}(\mathfrak{H}_{\mathcal{E}}) \cap \text{Cov}(\mathcal{E})$ and in this sense we are generalizing the framework described before Observation A allowing also the action of covariant η -quantum symmetries in addition to the fiber-preserving symmetry. In view of the equivalence established by Theorem 1.4 we can generalize Observation A as follows:

Observation B. *Once the problem of the classification of the C-symmetric η -self-adjoint endomorphisms is reduced to the problem of the classification of $\mathcal{H}_{\eta, 0}(\mathcal{E})$ it turns out that the only relevant η -quantum symmetries are induced by element of $\text{QS}_{\eta}(\mathfrak{H}_{\mathcal{E}}) \cap \text{Cov}(\mathcal{E})$ which commute or anti-commute with the metric η .*

Remark 4.16 (Motivations for the bundle point of view: Topological quantum systems). Hilbert bundles arise in condensed matter physics when one considers electronic systems in a periodic crystal background. In this case the invariance under translations allows to use the *Bloch-Floquet transform* [Kuc] which transform the “physical” Hilbert space $L^2(\mathbb{R}^d)$ in the L^2 -space of an Hilbert bundle \mathcal{E}

over the torus $\mathbb{T}^d \simeq \mathbb{R}^d/\mathbb{Z}^d$ which is usually called the Brillouin zone. Under some regularity condition the Hamiltonian H of the system is mapped into an operator \hat{H} associated to a fiber-preserving endomorphism. Translation-invariant *unitary* operators are mapped into *fiber-preserving* unitary operators while translation-invariant *anti-unitary* operators are mapped into *covariant* unitary operators associated with the involution $\tau : \mathbb{T}^d \rightarrow \mathbb{T}^d$ given by $\tau(k) := -k$. For more details on this construction we refer to [DG1, Section 2] and reference therein. However, the bundle picture is also typical of more general physical systems, not necessarily related with periodic electronic structure, called Topological Quantum Systems (see e. g. [DG2, Definition 1.1]). ◀

5. K -THEORY

The main goal of this Section is to construct a K -theory capable of classifying gapped η -self adjoint operators with a C -symmetry. In order to do that we will use a re-formulation of the Freed-Moore K -theory introduced in [FM] in the spirit of the Karoubi's presentation of the K -theory [Kar]. In the process we will also touch the Fredholm version of the K -theory as initially proposed in [Ati3]. Section 5.1 and Section 5.3 provide a soft overview of the results contained in [Gom]. The rest of the material is instead original.

5.1. Karoubi's formulation of Freed-Moore K -theory. The twisted equivariant K -theory introduced by Freed and Moore in [FM] unifies the “conventional” twisted equivariant complex K -theory [DK, Ros, FHT] and Atiyah's KR -theory [Ati1]. The relevance of this K -theory is that it handles complex anti-linear symmetry. In this Section we provide a re-formulation of the Freed-Moore K -theory using the Karoubi's construction.

A sufficiently general setup for the Freed-Moore K -theory, at least for applications to condensed matter physics, requires the following ingredients:

Assumption 5.1 (Standard framework). *Let $X, \mathbb{G}, \varpi, c, \tau$ be as follows:*

- (a) X is a compact Hausdorff space;
- (b) \mathbb{G} is a finite group acting on X from the left;
- (c) $\varpi : \mathbb{G} \rightarrow \mathbb{Z}_2$ and $c : \mathbb{G} \rightarrow \mathbb{Z}_2$ are homomorphisms;
- (d) $\tau \in Z_{\text{group}}^2(\mathbb{G}; C(X, \mathbb{U}(1))_{\varpi})$.

Property (b) says that X is a (left) \mathbb{G} -space. In (c) the symbol $\mathbb{Z}_2 = \{\pm 1\}$ denotes the cyclic group of order 2. Condition (d) needs some elucidation. Let $C(X, \mathbb{U}(1))$ the group of the continuous functions on X with values in the unitary group $\mathbb{U}(1) := \mathbb{U}(\mathbb{C})$. This group admits a left-action and a right-action of \mathbb{G} . The left-action of $g \in \mathbb{G}$ on $f \in C(X, \mathbb{U}(1))$ is given by $f \mapsto f^{\varpi(g)}$. The right-action is given by the pull-back $f \mapsto g^*f$ where $g^*f(x) := f(gx)$. The group $C(X, \mathbb{U}(1))$ endowed with this \mathbb{G} -bimodule structure is denoted with $C(X, \mathbb{U}(1))_{\varpi}$. A 2-cocycle τ is a map $\tau : \mathbb{G} \times \mathbb{G} \rightarrow C(X, \mathbb{U}(1))_{\varpi}$ satisfying

$$\tau(g_2, g_3)^{\varpi(g_1)} = \tau(g_1 g_2, g_3) \tau(g_1, g_2 g_3)^{-1} g_3^* \tau(g_1, g_2), \quad \forall g_1, g_2, g_3 \in \mathbb{G} \quad (5.1)$$

and the set of the 2-cocycles is denoted with $Z_{\text{group}}^2(\mathbb{G}; C(X, \mathbb{U}(1))_{\varpi})$. Then, condition (d) can be rephrased by saying that τ is a group 2-cocycle of \mathbb{G} with values in the \mathbb{G} -bimodule $C(X, \mathbb{U}(1))_{\varpi}$. Without loss of generality one can assume that the 2-cocycle τ is *normalized*, i. e. $\tau(e, e) = 1$ where e is the unit of \mathbb{G} . By combining together 2-cocycle condition (5.1) and the normalization condition one gets

$$\tau(g, e) = 1 = \tau(e, g), \quad g \in \mathbb{G}. \quad (5.2)$$

The following definition is adapted from [FM, Definition 7.23].

Definition 5.2 (Twisted equivariant Hilbert bundle). *Let $X, \mathbb{G}, \varpi, c, \tau$ be as in Assumption 5.1. A (ϖ, c, τ) -twisted \mathbb{G} -equivariant (ungraded) Hilbert bundle on X with $C\ell^{r,s}$ -action (or a twisted bundle for short) is a Hilbert bundle $\pi : \mathcal{E} \rightarrow X$ equipped with the following data:*

- (a) A twisted \mathbb{G} -action provided by a representation ρ of \mathbb{G} on the total space \mathcal{E} which covers the left action of \mathbb{G} on X according to the following diagram

$$\begin{array}{ccc} \mathcal{E} & \xrightarrow{\rho(g)} & \mathcal{E} \\ \pi \downarrow & & \downarrow \pi \\ X & \xrightarrow{g} & X \end{array} \quad \forall g \in \mathbb{G}$$

and which satisfies

$$\rho(g)^* = \rho(g)^{-1} \quad (\text{isometry})$$

$$\rho(g) \circ \mathbb{1} = \varpi(g) \circ \rho(g) \quad (\text{linearity vs. anti-linearity})$$

$$\rho(g_1)\rho(g_2) = \tau(g_1, g_2)\rho(g_1g_2) \quad (\text{projective representation})$$

for all $g, g_1, g_2 \in \mathbb{G}$.

- (b) A (unitary) action γ of the Clifford algebra $C\ell^{r,s}$ (in the sense of Definition 4.14) which meets the compatibility condition

$$\gamma(a)\rho(g) = c(g)\rho(g)\gamma(a) \quad (\text{Koszul sign rule})$$

for all $a \in C\ell^{r,s}$ and $g \in \mathbb{G}$.

A homomorphism $f : (\mathcal{E}, \rho, \gamma) \rightarrow (\mathcal{E}', \rho', \gamma')$ of twisted bundles is a complex linear map $f : \mathcal{E} \rightarrow \mathcal{E}'$ which covers the identity of X and satisfies

$$f \circ \rho(g) = \rho'(g) \circ f$$

$$f \circ \gamma(a) = \gamma'(a) \circ f$$

for all $a \in C\ell^{r,s}$ and $g \in \mathbb{G}$. When f is bijective $(\mathcal{E}, \rho, \gamma)$ and $(\mathcal{E}', \rho', \gamma')$ are called isomorphic.

Definition 5.2 works simultaneously for finite and infinite rank bundles. In the finite rank case the underlying Hilbert bundle reduces to nothing more than a finite rank Hermitian vector bundle. In both cases the structures (a) and (b) are defined in the same way. The infinite rank case will be relevant in Section 5.2.

Remark 5.3 (PUA-representation). Condition (a) in Definition 5.2 can be summarized by saying that the Hilbert bundle $\pi : \mathcal{E} \rightarrow X$ is endowed with the action of a (finite) group \mathbb{G} induced by the representation $\rho : \mathbb{G} \rightarrow \text{QS}(\mathfrak{H}_{\mathcal{E}}) \cap \text{Cov}(\mathcal{E})$ which associates to any $g \in \mathbb{G}$ the covariant quantum symmetry $\rho(g)$. More precisely, according to Definition 4.15 the operator $\rho(g)$ is induced by a covariant endomorphism of the bundle \mathcal{E} and it is unitary or anti-unitary. The presence of the 2-cocycle τ implies that the representation ρ is projective. In summary ρ is a covarianta PUA-representation¹⁰ in the sense of [Par, Thi]. ◀

The separation between PUA-representation and Clifford action in Definition 2 is in some sense artificial although useful for the K -theory construction described below. In fact the Clifford action can always be seen as a factor of an extended PUA-representation or, on the contrary, a PUA-representation can contain a Clifford action.

Proposition 5.4 (PUA-representation vs. Clifford action). *There is a categorical correspondence between twisted equivariant Hilbert bundles with Clifford action and twisted equivariant Hilbert bundles without Clifford action.*

Proof (sketch of). Let us start with a (ϖ, c, τ) -twisted \mathbb{G} -equivariant Hilbert bundle $\pi : \mathcal{E} \rightarrow X$ equipped with a $C\ell^{r,s}$ -action. Consider the extended group $\mathbb{G}' := \mathbb{G} \times \mathbb{Z}_2$ with $\mathbb{Z}_2 = \{\pm 1\}$ and let $p : \mathbb{G}' \rightarrow \mathbb{G}$ be the first factor projection given by $p((g, \epsilon)) = g$ for all $g \in \mathbb{G}$ and $\epsilon \in \mathbb{Z}_2$. The action of \mathbb{G}' on X is defined as follows: $g'x = p(g')x$ for all $g' := (g, \epsilon) \in \mathbb{G}'$ and $x \in X$. The

¹⁰PUA stands for projective, unitary, anti-unitary.

homomorphisms ϖ and c can be extended to homomorphisms $\varpi' : \mathbb{G}' \rightarrow \mathbb{Z}_2$ and $c' : \mathbb{G}' \rightarrow \mathbb{Z}_2$ as follows: $\varpi' := \varpi \circ p$ and $c'((g, \pm 1)) = \pm c(g)$ for all $g \in \mathbb{G}$. Also the 2-cocycle τ can be extended to a 2-cocycle $\tau' \in Z_{\text{group}}^2(\mathbb{G}'; C(X, \mathbb{U}(1))_{\varpi'})$ by following the prescription of Table 5.1. The

τ'	$(g_2, +1)$	$(g_2, -1)$
$(g_1, +1)$	$\tau(g_1, g_2)$	$\tau(g_1, g_2)$
$(g_1, -1)$	$c(g_1) \tau(g_1, g_2)$	$c(g_2) \tau(g_1, g_2)$

TABLE 5.1. Prescription for the extended 2-cocycle τ' on $\mathbb{G} \times \mathbb{Z}_2$.

map $\rho' : \mathbb{G}' \rightarrow \text{QS}(\mathfrak{H}_{\mathcal{E}}) \cap \text{Cov}(\mathcal{E})$ defined by $\rho'((g, +1)) := \rho(g)$ and $\rho'((g, -1)) := \rho(g)\gamma_{r+s}$ for all $g \in \mathbb{G}$ provides a PUA-representation of \mathbb{G}' on the underlying Hilbert bundle \mathcal{E} . Therefore $\pi : \mathcal{E} \rightarrow X$ can be equivalently seen as a (ϖ', c', τ') -twisted \mathbb{G}' -equivariant Hilbert bundle $\pi : \mathcal{E} \rightarrow X$ equipped with a $C\ell^{r,s-1}$ -action. By repeating this construction $r + s$ times one finally gets a twisted Hilbert bundle without an explicit Clifford action but with an extended PUA-representation which encodes the original Clifford action. ■

As in the case of standard untwisted bundles, one can take the direct sum of twisted bundles, and pull a τ -twisted bundle on X back to a $(\varphi^*\tau)$ -twisted bundle on Y by a \mathbb{G} -equivariant map $\varphi : Y \rightarrow X$. Notice that the so-called homotopy property holds true for twisted bundles, and homotopy equivalent maps induce isomorphic twisted bundles by pullback (up to a canonical isomorphism matching the induced twists). This is because the argument to prove the homotopy property in [Ati2] can be applied to twisted bundles, in view of the fact that for any $(\mathcal{E}, \rho, \gamma)$ and $(\mathcal{E}', \rho', \gamma')$ the complex vector bundle $\text{Hom}(\mathcal{E}, \mathcal{E}') \simeq \mathcal{E}^* \otimes \mathcal{E}'$ gives rise to a usual \mathbb{G} -equivariant vector bundle with commuting $C\ell^{r,s}$ -action whose invariant sections are in bijective correspondence with the homomorphisms of the twisted bundles.

Definition 5.5 (Gradation). *Let $(\mathcal{E}, \rho, \gamma)$ be a (ϖ, c, τ) -twisted \mathbb{G} -equivariant (ungraded) Hilbert bundle on X as in Definition 5.2. A gradation (or \mathbb{Z}_2 -grading) of $(\mathcal{E}, \rho, \gamma)$ is a $\Gamma \in \text{End}_{\text{lin}}(\mathcal{E})$ such that:*

- (a) Γ is a self-adjoint involution, i. e. $\Gamma = \Gamma^*$ and $\Gamma^2 = \mathbb{1}$;
- (b) The relations

$$\begin{aligned} \Gamma \rho(g) &= c(g) \rho(g) \Gamma \\ \Gamma \gamma(a) &= -\gamma(a) \Gamma \end{aligned}$$

hold for all $a \in C\ell^{r,s}$ and $g \in \mathbb{G}$.

As explained in Section 4.3, the condition $\Gamma \in \text{End}_{\text{lin}}(\mathcal{E})$ means that Γ can be seen as a continuous map that associates to each $x \in X$ a linear operator Γ_x on the fiber \mathcal{E}_x , such that $\Gamma_x = \Gamma_x^*$ and $\Gamma_x^2 = \mathbb{1}_x$. Then a gradation is nothing but a continuous family of Hermitian operators (or Hermitian matrices) squaring to the identity. Notice that the Clifford action γ is required to be *odd* with respect to the gradation Γ .

Remark 5.6 (Gapped Hamiltonians and gradations). A (non-trivial) *gapped Hamiltonian* is a self-adjoint element $H \in \text{End}_{\text{lin}}(\mathcal{E})$ which admits a continuous function $\lambda : X \rightarrow \mathbb{R}$ such that $H_x - \lambda(x) \mathbb{1}_x$ is invertible and compact in each fiber $\mathcal{H}_x := \pi^{-1}(x)$ and

$$\min \sigma(H_x) < \lambda(x) < \max \sigma(H_x) \quad (5.3)$$

where $\sigma(H_x) \subset \mathbb{R}$ is the spectrum of H_x . In this case the self-adjoint operator

$$\Gamma_H|_{\mathcal{H}_x} := \frac{H_x - \lambda(x) \mathbb{1}_x}{|H_x - \lambda(x) \mathbb{1}_x|}$$

is fiberwise well-defined, $\Gamma_H|_{\mathcal{H}_x} \neq \mathbb{1}_x$ in view of (5.3) and $(\Gamma_H|_{\mathcal{H}_x})^2 = \mathbb{1}_x$. The collection of the $\Gamma_H|_{\mathcal{H}_x}$ defines a gradation $\Gamma_H : \mathcal{E} \rightarrow \mathcal{E}$ in the sense of Definition 5.5 (a). Moreover, condition (b) is satisfied if and only if $H\rho(g) = c(g)\rho(g)H$ and $H\gamma(a) = -\gamma(a)H$ and this requires certain symmetries of the spectra $\sigma(H_x)$. The transformation $H \rightarrow \Gamma_H$ is induced by a homotopy as showed in Remark B.7. \blacktriangleleft

Two gradations Γ_0 and Γ_1 on $(\mathcal{E}, \rho, \gamma)$ are said to be *homotopic* if there is a gradation $\hat{\Gamma}$ on $p^*(\mathcal{E}, \rho, \gamma)$ such that $\hat{\Gamma}|_{X \times \{j\}} = \Gamma_j$ for $j = 0, 1$. Here $p : X \times [0, 1] \rightarrow X$ is the equivariant projection from the \mathbb{G} -space $X \times [0, 1]$ (with trivial action on the second component) onto the \mathbb{G} -space X . It is clear from this definition that a homotopy of gradations $\hat{\Gamma}$ amounts to a continuous family of gradations $\Gamma(t)$ on \mathcal{E} such that $\Gamma(0) = \Gamma_0$ and $\Gamma(1) = \Gamma_1$.

Let us focus on *triples* $(\mathcal{E}, \Gamma_0, \Gamma_1)$ given by a (ϖ, c, τ) -twisted \mathbb{G} -equivariant (finit rank) vector bundle $(\mathcal{E}, \rho, \gamma)$ on X as in Definition 5.2 and two gradations Γ_0 and Γ_1 as in Definition 5.5. One has the following two structure:

- *Isomorphism:* Two triples $(\mathcal{E}, \Gamma_0, \Gamma_1)$ and $(\mathcal{E}', \Gamma'_0, \Gamma'_1)$ are isomorphic if there is an isomorphism $f : \mathcal{E} \rightarrow \mathcal{E}'$ of twisted bundles such that $f \circ \Gamma_i = \Gamma'_i \circ f$ for both $i = 0, 1$;
- *Direct sum:* The triple

$$(\mathcal{E}, \Gamma_0, \Gamma_1) \oplus (\mathcal{E}', \Gamma'_0, \Gamma'_1) := (\mathcal{E} \oplus \mathcal{E}', \Gamma_0 \oplus \Gamma'_0, \Gamma_1 \oplus \Gamma'_1)$$

is called the direct sum of $(\mathcal{E}, \Gamma_0, \Gamma_1)$ and $(\mathcal{E}', \Gamma'_0, \Gamma'_1)$.

Clearly, the direct sum $\mathcal{E} \oplus \mathcal{E}'$ inherits the natural twisted structure induced by the direct sum of the PUA-representations $\rho \oplus \rho'$ and by the direct sum of the Clifford actions $\gamma \oplus \gamma'$.

Definition 5.7 (*K-groups - finite rank case*). *Let $X, \mathbb{G}, \varpi, c, \tau$ be as in Assumption 5.1. Then:*

- $\varpi \mathcal{M}_{\mathbb{G}}^{(\tau, c) + (r, s)}(X)$ is the abelian monoid of isomorphism classes of triples $(\mathcal{E}, \Gamma_0, \Gamma_1)$ with addition given by the direct sum of triples;
- $\varpi \mathcal{Z}_{\mathbb{G}}^{(\tau, c) + (r, s)}(X)$ is the submonoid consisting of isomorphism classes of triples $(\mathcal{E}, \Gamma_0, \Gamma_1)$ such that Γ_0 and Γ_1 are homotopy equivalent;
- $\varpi \mathcal{K}_{\mathbb{G}}^{(\tau, c) + (r, s)}(X) := \varpi \mathcal{M}_{\mathbb{G}}^{(\tau, c) + (r, s)}(X) / \varpi \mathcal{Z}_{\mathbb{G}}^{(\tau, c) + (r, s)}(X)$ is the quotient monoid.

The following result is crucial and is based on [Kar, Chap. III, Lemma 4.16].

Lemma 5.8. *The monoid $\varpi \mathcal{K}_{\mathbb{G}}^{(\tau, c) + (r, s)}(X)$ is an abelian group in which the additive inverse of $[(\mathcal{E}, \Gamma_0, \Gamma_1)]$ is given by $[(\mathcal{E}, \Gamma_1, \Gamma_0)]$.*

Proof. It suffices to prove that the gradations $\Gamma_0 \oplus \Gamma_1$ and $\Gamma_1 \oplus \Gamma_0$ on $\mathcal{E} \oplus \mathcal{E}$ are homotopy equivalent. This fact follows by observing that such a homotopy is realized by the map

$$[0, \pi/2] \ni \theta \mapsto \Gamma(\theta) := T_\theta \begin{pmatrix} \Gamma_0 & 0 \\ 0 & \Gamma_1 \end{pmatrix} T_\theta$$

where the matrix T_θ is defined in (B.8). \blacksquare

The groups $\varpi \mathcal{K}_{\mathbb{G}}^{(\tau, c) + (r, s)}(X)$ have the following periodicity [Gom].

Proposition 5.9 (Periodicity). *Let $X, \mathbb{G}, \varpi, c, \tau$ be as in Assumption 5.1 and $\varpi \mathcal{K}_{\mathbb{G}}^{(\tau, c) + (r, s)}(X)$ the abelian group introduced in Definition 5.7. Then $\varpi \mathcal{K}_{\mathbb{G}}^{(\tau, c) + (r, s)}(X)$ is subject to the following periodicities:*

$$\begin{aligned} \varpi \mathcal{K}_{\mathbb{G}}^{(\tau, c) + (r, s)}(X) &\simeq \varpi \mathcal{K}_{\mathbb{G}}^{(\tau, c) + (r+1, s+1)}(X) \\ &\simeq \varpi \mathcal{K}_{\mathbb{G}}^{(\tau, c) + (r+8, s)}(X) \\ &\simeq \varpi \mathcal{K}_{\mathbb{G}}^{(\tau, c) + (r, s+8)}(X) . \end{aligned}$$

In the case $\varpi \equiv +1$ is trivial, the group $\mathcal{K}_G^{(\tau,c)+(r,s)}(X) := {}^{+1}\mathcal{K}_G^{(\tau,c)+(r,s)}(X)$ is subject to the extra reduced periodicities:

$$\begin{aligned} \mathcal{K}_G^{(\tau,c)+(r,s)}(X) &\simeq \mathcal{K}_G^{(\tau,c)+(r+2,s)}(X) \\ &\simeq \mathcal{K}_G^{(\tau,c)+(r,s+2)}(X). \end{aligned}$$

Some details about the proof of Proposition 5.9 will be provided in Remark 5.18.

Remark 5.10 (Generalized cohomology). Proposition 5.9 shows that the groups ${}^\varpi\mathcal{K}_G^{(\tau,c)+(r,s)}(X)$ depend only on the difference $s - r \bmod 8$ and this justify the following notation

$${}^\varpi\mathcal{K}_G^{(\tau,c)+n}(X) := {}^\varpi\mathcal{K}_G^{(\tau,c)+(0,n)}(X) \simeq {}^\varpi\mathcal{K}_G^{(\tau,c)+(r,s)}(X), \quad n = s - r \bmod 8$$

The groups ${}^\varpi\mathcal{K}_G^{(\tau,c)+n}(X)$ form a \mathbb{G} -equivariant generalized cohomology theory, so that the homotopy axiom, the excision axiom, etc. hold true (see e. g. [FHT, Gom]). \blacktriangleleft

5.2. Infinite-dimensional Karoubi's formulation. Definition 5.7 is modeled on finite rank twisted vector bundles. In this Section we will focus on the case of twisted Hilbert bundles $\pi : \mathcal{E} \rightarrow X$ of infinite rank. This means that each fiber of \mathcal{E} is isomorphic to a separable Hilbert space \mathcal{H} of infinite dimension. The infinite dimensionality allows the existence of a *universal* twisted bundle of infinite rank [Gom, Section 2.6.]:

Lemma 5.11 (Universal twisted bundle). *Let $X, \mathbb{G}, \varpi, c, \tau$ be as in Assumption 5.1. There exists a (ϖ, c, τ) -twisted \mathbb{G} -equivariant Hilbert bundle $\mathcal{E}_{\text{univ}}$ on X with $C\ell^{r,s}$ -action and a gradation Γ_{univ} which are locally universal in the following sense: For any \mathbb{G} -invariant closed subspace $Y \subseteq X$ and any $(\varpi, c, \tau|_Y)$ -twisted \mathbb{G} -equivariant Hilbert bundle $\pi : \mathcal{E} \rightarrow Y$ with $C\ell^{r,s}$ -action and gradation Γ on it, there is an embedding $\mathcal{E} \rightarrow \mathcal{E}_{\text{univ}}|_Y$ of twisted bundles which preserves the gradations.*

The universal bundle $\mathcal{E}_{\text{univ}}$ in Lemma 5.11 has infinite rank, and also the eigenspaces of Γ_{univ} associated to the eigenvalues ± 1 are infinite dimensional. This allows the twisted bundle $\pi : \mathcal{E} \rightarrow Y$ to have infinite rank. It is known [Gom, Section 2.6.] that the space of embeddings $\mathcal{E} \rightarrow \mathcal{E}_{\text{univ}}|_Y$ is contractible with respect to an appropriate topology, and the pair $(\mathcal{E}_{\text{univ}}, \Gamma_{\text{univ}})$ that implements the local universality in Lemma 5.11 is essentially unique.

Definition 5.12. *Let $X, \mathbb{G}, \varpi, c, \tau$ be as in Assumption 5.1, and $(\mathcal{E}_{\text{univ}}, \Gamma_{\text{univ}})$ the locally universal pair of Lemma 5.11. Let $\mathcal{K}_x := \mathcal{K}(\mathcal{E}_{\text{univ},x})$ be the algebra of compact operators on the fiber $\mathcal{E}_{\text{univ},x}$. We denote with*

$$\mathcal{G}(\mathcal{E}_{\text{univ}}) := \{\Gamma : \mathcal{E}_{\text{univ}} \rightarrow \mathcal{E}_{\text{univ}} \mid \Gamma \text{ is gradation and } (\Gamma - \Gamma_{\text{univ}})_x \in \mathcal{K}_x \quad \forall x \in X\}$$

the set of gradations Γ on $\mathcal{E}_{\text{univ}}$ which differ from Γ_{univ} by a compact operator on each fiber.

In view of Definition 5.12 any $\Gamma \in \mathcal{G}(\mathcal{E}_{\text{univ}})$ is a *compact perturbation* of the universal gradation Γ_{univ} . This implies that also the eigenspaces of Γ associated to the eigenvalues ± 1 are infinite dimensional.

A homotopy on $\mathcal{G}(\mathcal{E}_{\text{univ}})$ is a family of gradations which is continuous with respect to the operator norm on the fibers of $\mathcal{E}_{\text{univ}}$. Such homotopies introduce an equivalence relation in $\mathcal{G}(\mathcal{E}_{\text{univ}})$ denoted with \sim . The following result is proved in [Gom, Lemma 4.16].

Proposition 5.13. *Let $X, \mathbb{G}, \varpi, c, \tau$ be as in Assumption 5.1. There is a natural bijection*

$${}^\varpi\mathcal{K}_G^{(\tau,c)+(r,s)}(X) \simeq \mathcal{G}(\mathcal{E}_{\text{univ}}) / \sim.$$

The bijection in Proposition 5.13 is constructed as follows. Let $(\mathcal{E}, \Gamma_0, \Gamma_1)$ be a triple consisting of a finite rank twisted vector bundle \mathcal{E} on X and its gradations Γ_0 and Γ_1 . The nature of the locally universal structure described in Lemma 5.11 allows us to embed the pair (\mathcal{E}, Γ_1) into $(\mathcal{E}_{\text{univ}}, \Gamma_{\text{univ}})$. This embedding induces an orthogonal decomposition $\mathcal{E}_{\text{univ}} \simeq \mathcal{E} \oplus \mathcal{E}^\perp$ of twisted bundles which respects the gradations in the sense that $\Gamma_{\text{univ}} = \Gamma_1 \oplus \Gamma^\perp$ where Γ^\perp is a gradation on the complement \mathcal{E}^\perp .

Now, from the second gradation Γ_0 on \mathcal{E} one has a second gradation $\Gamma_0 \oplus \Gamma^\perp$ on $\mathcal{E}_{\text{univ}}$, which differs from Γ_{univ} by $(\Gamma_0 - \Gamma_1) \oplus 0$ which is a compact operator on each fiber of $\mathcal{E}_{\text{univ}}$. In view of the definition of $\varpi \mathcal{K}_G^{(\tau,c)+(r,s)}(X) = \varpi \mathcal{M}_G^{(\tau,c)+(r,s)}(X) / \varpi \mathcal{Z}_G^{(\tau,c)+(r,s)}(X)$ and the property of $\mathcal{E}_{\text{univ}}$, the assignment $(\mathcal{E}, \Gamma_0, \Gamma_1) \mapsto \Gamma_0 \oplus \Gamma^\perp$ induces a well-defined natural map

$$J : \varpi \mathcal{K}_G^{(\tau,c)+(r,s)}(X) \longrightarrow \mathcal{G}(\mathcal{E}_{\text{univ}}) / \sim$$

which indeed realizes the bijection in Proposition 5.13. The proof of the bijectivity of J requires the construction of the inverse to J by a finite dimensional “approximation” of the triple $(\mathcal{E}_{\text{univ}}, \Gamma, \Gamma_{\text{univ}})$ formed by an element $\Gamma \in \mathcal{G}(\mathcal{E}_{\text{univ}})$. We refer to [Gom, Section 4.3] for the detail.

In view of Proposition 5.13 we can generalize Definition 5.7 for the infinite rank case:

Definition 5.14 (*K-groups - infinite rank case*). *Let $X, \mathbb{G}, \varpi, c, \tau$ be as in Assumption 5.1. Then:*

- $\varpi \mathcal{M}_G^{(\tau,c)+(r,s)}(X)$ is the abelian monoid of isomorphism classes of triples $(\mathcal{E}, \Gamma_0, \Gamma_1)$ where \mathcal{E} is a (ϖ, c, τ) -twisted \mathbb{G} -equivariant Hilbert bundle over X and the two gradations verify the compactness condition

$$(\Gamma_0 - \Gamma_1)_x \in \mathcal{K}_x \quad \forall x \in X.$$

Isomorphisms and addition are defined as in Section 5.1;

- $\varpi \mathcal{Z}_G^{(\tau,c)+(r,s)}(X)$ and $\varpi \mathcal{K}_G^{(\tau,c)+(r,s)}(X)$ are defined as in Definition 5.7.

The only difference between Definition 5.7 and Definition 5.14 is the addition of the compactness condition. Clearly, whenever \mathcal{E} is a finite rank vector bundle Definition 5.14 reduces to Definition 5.7.

5.3. Fredholm operator formulation of Freed-Moore K-theory. Topological K -theory has various formulations. The original formulation of the Freed-Moore K -theory is based on finite rank twisted bundles [FM, Definition 7.33]. However, a possible formulation based on Fredholm operators is also mentioned [FM, Remark 7.35]. In this formulation the basic objects are Fredholm families.

Definition 5.15 (Self-adjoint Fredholm family). *Let $(\mathcal{E}, \rho, \gamma)$ be a (ϖ, c, τ) -twisted \mathbb{G} -equivariant Hilbert bundle on X with $C\ell^{r,s}$ -action, and Γ a gradation on $(\mathcal{E}, \rho, \gamma)$. A self-adjoint Fredholm family is a self-adjoint element $A \in \text{End}_{\text{lin}}(\mathcal{E})$. This means that A restricts to a linear bounded operator $A_x = A|_{\mathcal{H}_x}$ on each fiber Hilbert space $\mathcal{H}_x := \pi^{-1}(x)$ and*

- (a) A_x is self-adjoint, $A_x = A_x^*$.

In addition one requires that:

- (b) $A_x^2 - \mathbb{1}_x$ is compact;
- (c) $\sigma(A_x) \subset [-1, +1]$;

Finally the following compatibility relations are required:

- (d) For all $a \in C\ell^{r,s}$ and $g \in \mathbb{G}$ it holds that

$$\begin{aligned} A\rho(g) &= c(g) \rho(g)A \\ A\gamma(a) &= -\gamma(a)A \\ A\Gamma &= -\Gamma A. \end{aligned}$$

Property (a) in Definition 5.15 says that A_x is a (bounded) self-adjoint operator for each $x \in X$ and this justifies the reality of the spectrum $\sigma(A_x) \subset \mathbb{R}$. Property (b) says that A_x is an involution (hence an invertible) modulo compact operators. In view of the Atkinson characterization [Atk] this implies that A_x is a self-adjoint Fredholm operator for each $x \in X$. The following result shows the strong link between self-adjoint Fredholm families and gradations.

Lemma 5.16. *Let A be a self-adjoint Fredholm family of the (ϖ, c, τ) -twisted \mathbb{G} -equivariant Hilbert bundle $(\mathcal{E}, \rho, \gamma)$ on X , and Γ a gradation on $(\mathcal{E}, \rho, \gamma)$. Then the element*

$$\vartheta(A) := -e^{\pi A \Gamma} \Gamma$$

provides a second gradation of $(\mathcal{E}, \rho, \gamma)$ and the difference $\Gamma - \vartheta(A)$ is compact in each fiber.

Proof. Since A and Γ are by definition elements in $\text{End}_{\text{lin}}(\mathcal{E})$ then also the product $A\Gamma$ is a fiber preserving linear endomorphism of \mathcal{E} . Fiberwise one has that $i(A\Gamma)_x = iA_x\Gamma_x$ is a self-adjoint operator in view of the anti-commutativity of A and Γ . Therefore $e^{\pi(A\Gamma)_x} = e^{-i(i\pi A_x\Gamma_x)}$ is a well defined unitary operator in each fiber. The element $\vartheta(A) \in \text{End}_{\text{lin}}(\mathcal{E})$ is then defined fiberwise by $\vartheta(A)_x := -e^{\pi(A\Gamma)_x}\Gamma_x$. Direct computations show that $\vartheta(A)$ is a self-adjoint involution. Indeed

$$\vartheta(A)^* = -\Gamma^* e^{\pi(A\Gamma)^*} = -\Gamma e^{-\pi A \Gamma} = -e^{\pi A \Gamma} \Gamma = \vartheta(A)$$

and

$$\vartheta(A)^2 = e^{\pi A \Gamma} \Gamma e^{\pi A \Gamma} \Gamma = \Gamma e^{-\pi A \Gamma} e^{\pi A \Gamma} \Gamma = \mathbb{1}$$

where the computations can be understood fiberwise, if necessary. Since $A\Gamma$ commutes with $\gamma(a)$ and $\rho(g)$ for all $a \in C\ell^{r,s}$ and $g \in \mathbb{G}$ it follows that the commutation relations between $\vartheta(A)$ and $\gamma(a)$ or $\rho(g)$ are the same commutation relations of Γ . This proves that $\vartheta(A)$ is a gradation according to Definition 5.5. To finish the proof one has to prove that $\Gamma - \vartheta(A) = (\mathbb{1} + e^{\pi A \Gamma})\Gamma$ is compact in each fiber. For that it is enough to show that $\mathbb{1}_x + e^{\pi A_x \Gamma_x}$ is a compact operator for each $x \in X$. Observe that the compactness of $\mathbb{1}_x - A_x^2$ implies that the spectrum of A_x^2 is pure point with 1 as the only possible accumulation point. Moreover, the only eigenvalue with possible infinite multiplicity is indeed 1. From the spectral mapping theorem one concludes that $\sigma(A_x)$ is pure point and the only eigenvalues with possible infinite multiplicity are ± 1 . Let us anticipate the following result:

$$\lambda \in \sigma(A_x) \quad \Leftrightarrow \quad i\lambda \in \sigma(A_x \Gamma_x). \quad (5.4)$$

Condition (5.4) implies that also $\sigma(A_x \Gamma_x)$ is pure point with $\pm i$ as the only possible eigenvalues of infinite multiplicity. As a consequence the spectrum of $\mathbb{1}_x + e^{\pi A_x \Gamma_x}$ is pure point with 0 as the only possible eigenvalue of infinite multiplicity. This is equivalent to the compactness of $\mathbb{1}_x + e^{\pi A_x \Gamma_x}$. Finally, the proof of (5.4). Let Π_x^\pm be the spectral projections of Γ_x related to the eigenvalues ± 1 respectively. Then $\Gamma_x = \Pi_x^+ - \Pi_x^-$. As a consequence of the anti-commutation between A and Γ one has that $A_x = A_x^{-,+} + A_x^{+,-}$ where $A_x^{\pm,\mp} := \Pi_x^\pm A_x \Pi_x^\mp$. Said differently A_x is completely off-diagonal with respect to the grading induced by Γ_x . As a consequence $A_x \Gamma_x = A_x^{-,+} - A_x^{+,-}$. Let $\lambda \in \text{Res}(A_x)$ be a point of the resolvent set of A_x . Then the resolvent $R_x(\lambda) := (A_x - \lambda \mathbb{1}_x)^{-1}$ exists and can be decomposed along Π_x^\pm as follows: $R_x(\lambda) = R_x^{+,+} + R_x^{-,-} + R_x^{+,-} + R_x^{-,+}$. The conditions $(A_x - \lambda \mathbb{1}_x)R_x(\lambda) = \mathbb{1}_x = R_x(\lambda)(A_x - \lambda \mathbb{1}_x)$ imply

$$\begin{aligned} A_x^{-,+} R_x^{+,+} &= \lambda R_x^{-,+} = R_x^{-,-} A_x^{-,+} \\ A_x^{+,-} R_x^{-,-} &= \lambda R_x^{+,-} = R_x^{+,+} A_x^{+,-} \\ A_x^{-,+} R_x^{+,-} &= \lambda R_x^{-,-} + \Pi_x^- = R_x^{-,+} A_x^{+,-} \\ A_x^{+,-} R_x^{-,+} &= \lambda R_x^{+,+} + \Pi_x^+ = R_x^{+,-} A_x^{-,+}. \end{aligned} \quad (5.5)$$

Consider now the operator $\tilde{R}_x(\lambda) := -i(R_x^{+,+} + R_x^{-,-}) + R_x^{+,-} - R_x^{-,+}$. By using the relations (5.5) a direct computation shows that $(A_x \Gamma_x - i\lambda \mathbb{1}_x)\tilde{R}_x(\lambda) = \mathbb{1}_x = \tilde{R}_x(\lambda)(A_x \Gamma_x - i\lambda \mathbb{1}_x)$, hence $i\lambda \in \text{Res}(A_x \Gamma_x)$. On the other hand if $i\lambda \in \text{Res}(A_x \Gamma_x)$ then $\lambda \in \text{Res}(-iA_x \Gamma_x)$ and in view of the same argument above this implies $i\lambda \in \text{Res}(-iA_x)$ or equivalently $\lambda \in \text{Res}(-A_x) = \text{Res}(\Gamma_x A_x \Gamma_x) = \text{Res}(A_x)$. In summary we proved that $\lambda \in \text{Res}(A_x)$ if and only if $i\lambda \in \text{Res}(A_x \Gamma_x)$ and therefore (5.4) follows. ■

As it has been pointed out in Section 5.2, any finite rank twisted bundle $(\mathcal{E}, \rho, \gamma)$ with a gradation Γ can be embedded into the locally universal bundle $(\mathcal{E}_{\text{univ}}, \rho_{\text{univ}}, \gamma_{\text{univ}})$ with its universal gradation Γ_{univ} . A Fredholm family A on \mathcal{E} can be extended to one on $\mathcal{E}_{\text{univ}}$ in an essentially unique way. Thus, it turns out that it is sufficient to consider Fredholm families on $\mathcal{E}_{\text{univ}}$. Let $\mathcal{F}(\mathcal{E}_{\text{univ}})$ be the space of Fredholm families on $\mathcal{E}_{\text{univ}}$. A notion of a homotopy between Fredholm families can be defined in a natural

way [Gom, Section 3.1] and the space $\mathcal{F}(\mathcal{E}_{\text{univ}})$ turns out to be endowed with an induced equivalence relation (still) denoted with \sim .

Proposition 5.17. *There is a natural bijection between $\mathcal{G}(\mathcal{E}_{\text{univ}})/\sim$ and $\mathcal{F}(\mathcal{E}_{\text{univ}})/\sim$ which, in view of Proposition 5.13, implies*

$$\varpi \mathcal{K}_G^{(\tau,c)+(r,s)}(X) \simeq \mathcal{F}(\mathcal{E}_{\text{univ}})/\sim.$$

Proposition 5.17 is essentially proved in [Gom, Section 4.2 & Section 4.2] and the proof boils down to the approach developed in [ASi]. The basic result is provided by Lemma 5.16 which assures that for a given self-adjoint Fredholm family $A \in \mathcal{F}(\mathcal{E}_{\text{univ}})$ the operator $\vartheta(A) := e^{\pi A \Gamma_{\text{univ}}} \Gamma_{\text{univ}}$ is a gradation such that $\vartheta(A) - \Gamma_{\text{univ}}$ is fiberwise compact. Then ϑ defines a map $\vartheta : \mathcal{F}(\mathcal{E}_{\text{univ}}) \rightarrow \mathcal{G}(\mathcal{E}_{\text{univ}})$ which descends to a map on the set of homotopy classes. By means a suitable generalization of the argument used in [ASi], and the previous isomorphisms proved in Proposition 5.13, one can show that the induced map ϑ realizes indeed the bijection in Proposition 5.17.

Remark 5.18 (Recalibration of the twist). In [Gom, Definition 3.1] the notion of the Fredholm family is given in terms of skew-adjoint operators and the self-adjoint version given in Definition 5.15 is mentioned in [Gom, Remark 4.12]. These Fredholm families lead to the K -theories $\varpi \mathbb{K}_G^{(\tau,c)+(r,s)}(X)$ and $\varpi \mathbb{K}_G^{(\tau,c)+(r,s)}(X)$, respectively. The former is the K -theory mainly studied in [Gom] while the latter is nothing but $\mathcal{F}(\mathcal{E}_{\text{univ}})/\sim$ in Proposition 5.17. Usually, one has a one-to-one correspondence between self-adjoint operators and skew-adjoint operators given by the multiplication with the imaginary unit i . However, in the presence of a non-trivial ϖ , this correspondence violates the compatibility relations with the $\rho(g)$ (cf. Definition 5.2), and consequently does not provide an isomorphism between $\varpi \mathbb{K}_G^{(\tau,c)+(r,s)}(X)$ and $\varpi \mathbb{K}_G^{(\tau,c)+(r,s)}(X)$. Another correspondence between self-adjoint operators and skew-adjoint operators is provided by $A \mapsto A\Gamma$ whenever Γ is a gradation such that $A\Gamma = -\Gamma A$. The correspondence $A \mapsto A\Gamma$ together with a simultaneous modification of the symmetries $\rho \mapsto \hat{\rho}$ and $\gamma \mapsto \hat{\gamma}$ where

$$\hat{\gamma}(a) := \gamma(a)\Gamma, \quad \hat{\rho}(g) := \begin{cases} \rho(g), & \text{if } c(g) = +1 \\ \rho(g)\Gamma & \text{if } c(g) = -1 \end{cases},$$

leads to the isomorphism $\varpi \mathbb{K}_G^{(\tau,c)+(r,s)}(X) \rightarrow \varpi \mathbb{K}_G^{(\hat{\tau},\hat{c})+(s,r)}(X)$, where the new cocycle $\hat{\tau}$ is related to τ by

$$\hat{\tau}(g_1, g_2) := \begin{cases} +\tau(g_1, g_2), & \text{if } c(g_1) = 1 \text{ or } c(g_2) = 1 \\ -\tau(g_1, g_2) & \text{if } c(g_1) = c(g_2) = -1 \end{cases}$$

The cocycles τ and $\hat{\tau}$ are equivalent if ϖ or c are trivial, so that their difference is essential only if both ϖ and c are non-trivial. Notice also that $\hat{\hat{\tau}} = \tau$. The isomorphisms

$$\varpi \mathcal{K}_G^{(\tau,c)+(r,s)}(X) \simeq \varpi \mathbb{K}_G^{(\tau,c)+(r,s)}(X) \simeq \varpi \mathbb{K}_G^{(\hat{\tau},\hat{c})+(s,r)}(X) \quad (5.6)$$

combined with the periodicity in [Gom, Lemma 3.5] provides a proof of Proposition 5.9. In particular one has that

$$\varpi \mathcal{K}_G^{(\tau,c)+n}(X) \simeq \varpi \mathbb{K}_G^{(\hat{\tau},\hat{c})+n}(X) \quad (5.7)$$

where $n := s - r \bmod 8$. The latter isomorphism will be used in the computation of Section 5.6. \blacktriangleleft

Remark 5.19. Let $\varpi \text{Vec}_G^{(\tau,c)+(r,s)}(X)$ be the monoid of isomorphism classes of finite rank (ϖ, c, τ) -twisted G -equivariant vector bundles on X with $C\ell^{r,s}$ -action and gradations. There is a forgetful functor which induces a homomorphism $f : \varpi \text{Vec}_G^{(\tau,c)+(r,s+1)}(X) \rightarrow \varpi \text{Vec}_G^{(\tau,c)+(r,s)}(X)$. We write $\varpi \text{Triv}_G^{(\tau,c)+(r,s)}(X)$ for the submonoid given by the image of f in $\varpi \text{Vec}_G^{(\tau,c)+(r,s)}(X)$ and

$$\varpi \mathbb{K}_G^{(\tau,c)+(r,s)}(X)_{\text{fin}} := \varpi \text{Vec}_G^{(\tau,c)+(r,s)}(X) / \varpi \text{Triv}_G^{(\tau,c)+(r,s)}(X)$$

for the quotient monoid. In the case that c is trivial and $r = s = 0$, one can identify $\varpi \mathbb{K}_G^{\tau+(0,0)}(X)_{\text{fin}}$ with the Grothendieck construction applied to the monoid of isomorphism classes of finite rank (ϖ, τ) -twisted G -equivariant (ungraded) vector bundles on X without Clifford action. Any (ϖ, τ) -twisted

vector bundle (\mathcal{E}, ρ) on X with a gradation Γ admits the trivial (skew-adjoint) Fredholm family $A \equiv 0$. Thus, by the universality, we have a natural homomorphism $\iota : {}^\varpi \mathbb{K}_G^{\tau+(0,0)}(X)_{\text{fin}} \rightarrow {}^\varpi \mathbb{K}_G^{\tau+(0,0)}(X)$. It is known [FM] that the homomorphism ι is bijective. This fact justifies the original formulation of the Freed-Moore K -theory based on finite rank twisted vector bundles. We refer the reader to [Gom, Section 4.4] for more details about the finite rank realizability of the K -theory. \blacktriangleleft

5.4. K -theory for η -self adjoint operators. The framework described in Section 5.1 admits a suitable generalization adapted for the case of gapped η -self adjoint operators with a C -symmetry. Let us recall two fundamental facts. In view of Corollary 4.10 every η -selfadjoint operator with a C -symmetry can be η -unitarily reduced to a self-adjoint operator $H = H^*$ commuting with the fundamental symmetry η . Moreover, Theorem 1.4 shows that this reduction transforms η -quantum symmetries compatible with the C -symmetry in unitary or anti-unitary operators which commute or anti-commute with the metric η according to the type of the original symmetry. The bundle version of these facts has been summarized in Observation B in Section 4.3. The latter justifies the introduction of the following structures.

Assumption 5.20 (Framework for η -self adjoint systems). *Let $X, \mathbb{G}, \varpi, \wp, c, \tau$ be as follows:*

- (a) X is a compact Hausdorff space;
- (b) \mathbb{G} is a finite group acting on X from the left.
- (c) $\varpi : \mathbb{G} \rightarrow \mathbb{Z}_2, \wp : \mathbb{G} \rightarrow \mathbb{Z}_2$ and $c : \mathbb{G} \rightarrow \mathbb{Z}_2$ are homomorphisms.
- (d) $\tau \in Z_{\text{group}}^2(\mathbb{G}; C(X, \mathbb{U}(1))_{\varpi})$

Definition 5.21 (Twisted equivariant Krein bundle). *Let $X, \mathbb{G}, \varpi, \wp, c, \tau$ be as in Assumption 5.20. A (ϖ, \wp, c, τ) -twisted \mathbb{G} -equivariant (ungraded) Krein bundle on X with $C\ell^{r,s}$ -action (or a twisted Krein bundle for short) is a Krein vector bundle $\pi : \mathcal{E} \rightarrow X$ with fundamental symmetry η (cf. Definition 4.13) equipped with the following data:*

- (a) A twisted \mathbb{G} -action provided by a representation ρ of \mathbb{G} on the total space \mathcal{E} which covers the left action of \mathbb{G} on X according to the following diagram

$$\begin{array}{ccc} \mathcal{E} & \xrightarrow{\rho(g)} & \mathcal{E} \\ \pi \downarrow & & \downarrow \pi \\ X & \xrightarrow{g} & X \end{array} \quad \forall g \in \mathbb{G}$$

and which satisfies

$$\begin{aligned} \rho(g)^* &= \rho(g)^{-1} && \text{(isometry)} \\ \rho(g)\eta &= \wp(g)\eta\rho(g) && \text{(proper-isometry vs. pseudo-isometry)} \\ \rho(g) \circ \mathbb{1} &= \varpi(g) \circ \rho(g) && \text{(linearity vs. anti-linearity)} \\ \rho(g_1)\rho(g_2) &= \tau(g_1, g_2)\rho(g_1g_2) && \text{(projective representation)} \end{aligned}$$

for all $g, g_1, g_2 \in \mathbb{G}$.

- (b) A (unitary) action γ of the Clifford algebra $C\ell^{r,s}$ (in the sense of Definition 4.14) which meets the structural conditions

$$\begin{aligned} \gamma(a)\eta &= \eta\gamma(a) && \text{(compatibility Krein-Clifford structure)} \\ \gamma(a)\rho(g) &= c(g)\rho(g)\gamma(a) && \text{(Koszul sign rule)} \end{aligned}$$

for all $a \in C\ell^{r,s}$ and $g \in \mathbb{G}$.

A homomorphism $f : (\mathcal{E}, \eta, \rho, \gamma) \rightarrow (\mathcal{E}', \eta', \rho', \gamma')$ of twisted Krein bundles is a complex linear map $f : \mathcal{E} \rightarrow \mathcal{E}'$ which covers the identity of X and satisfies

$$\begin{aligned} f \circ \eta &= \eta' \circ f \\ f \circ \rho(g) &= \rho'(g) \circ f \\ f \circ \gamma(a) &= \gamma'(a) \circ f \end{aligned}$$

for all $a \in C\ell^{r,s}$ and $g \in \mathbb{G}$. When f is bijective $(\mathcal{E}, \eta, \rho, \gamma)$ and $(\mathcal{E}', \eta', \rho', \gamma')$ are called isomorphic.

Remark 5.22 (η -PUA-representation). The only difference between Assumption 5.1 and Assumption 5.20 is in item (c). In the latter case the presence of the map \wp allows to discriminate between (anti)unitarity and pseudo-(anti)unitarity symmetries. Condition (a) in Definition 5.21 says that on the bundle $\pi : \mathcal{E} \rightarrow X$ acts the (finite) group \mathbb{G} through the representation $\rho : \mathbb{G} \rightarrow \text{QS}_\eta(\mathfrak{H}_\mathcal{E}) \cap \text{Cov}(\mathcal{E})$ which associates to any $g \in \mathbb{G}$ the covariant η -quantum symmetry $\rho(g)$. More precisely any $\rho(g)$ is induced by a covariant endomorphism of the bundle \mathcal{E} which is unitary or anti-unitary and which commute or anti-commute with the metric η . The combination of these properties is modeled out of Observation B in Section 4.3. The presence of the 2-cocycle τ suggests that the representation ρ is projective, in general. By adapting the jargon of Remark 5.3, we can refer to ρ as a covariant η -PUA-representation of the group \mathbb{G} . ◀

Gapped η -self-adjoint Hamiltonians with a C -symmetry are reducible to gapped self-adjoint Hamiltonians commuting with η . At this point the “spectral flattening” described in Remark 5.6 can be applied as well. This observation suggests that gapped η -self-adjoint operators with a C -symmetry can be classified by η -gradations:

Definition 5.23 (η -gradation). Let $(\mathcal{E}, \eta, \rho, \gamma)$ be a (ϖ, \wp, c, τ) -twisted \mathbb{G} -equivariant (ungraded) Krein bundle on X as in Definition 5.21. An η -gradation of $(\mathcal{E}, \eta, \rho, \gamma)$ is a $\Gamma \in \text{End}_{\text{lin}}(\mathcal{E})$ such that:

- (a1) Γ is a self-adjoint involution, i. e. $\Gamma = \Gamma^*$ and $\Gamma^2 = \mathbb{1}$;
- (a2) $\Gamma\eta = \eta\Gamma$;
- (b) The relations

$$\begin{aligned}\Gamma\rho(g) &= c(g)\rho(g)\Gamma \\ \Gamma\gamma(a) &= -\gamma(a)\Gamma\end{aligned}$$

hold for all $a \in C\ell^{r,s}$ and $g \in \mathbb{G}$.

By mimicking the construction of the Freed-Moore K -theory discussed in Section 5.1 one can introduce the notions of isomorphism of twisted Krein bundles and of homotopy equivalence of η -gradations. By focusing on triples $(\mathcal{E}, \Gamma_0, \Gamma_1)$ given by a (ϖ, \wp, c, τ) -twisted \mathbb{G} -equivariant Krein bundle $(\mathcal{E}, \eta, \rho, \gamma)$ on X and two η -gradations Γ_0 and Γ_1 , along with the notions of equivalence relation and direct sum introduced in Section 5.1, one defines the following objects:

Definition 5.24 (ηK -groups - finite rank case). Let $X, \mathbb{G}, \varpi, \wp, c, \tau$ be as in Assumption 5.20. Then:

- $\varpi_\eta \mathcal{M}_\mathbb{G}^{(\tau, c, \wp) + (r, s)}(X)$ is the abelian monoid of isomorphism classes of triples $(\mathcal{E}, \Gamma_0, \Gamma_1)$ with addition given by the direct sum of triples;
- $\varpi_\eta \mathcal{Z}_\mathbb{G}^{(\tau, c, \wp) + (r, s)}(X)$ is the submonoid consisting of isomorphism classes of triples $(\mathcal{E}, \Gamma_0, \Gamma_1)$ such that Γ_0 and Γ_1 are homotopy equivalent;
- $\varpi_\eta \mathcal{K}_\mathbb{G}^{(\tau, c, \wp) + (r, s)}(X) := \varpi_\eta \mathcal{M}_\mathbb{G}^{(\tau, c, \wp) + (r, s)}(X) / \varpi_\eta \mathcal{Z}_\mathbb{G}^{(\tau, c, \wp) + (r, s)}(X)$ is the quotient monoid.

The properties of the K -theory defined above will be investigated in the last part of this work by a reduction of the monoids $\varpi_\eta \mathcal{K}_\mathbb{G}^{(\tau, c, \wp) + (r, s)}(X)$ to the Freed-Moore K -theory $\varpi_\eta \mathcal{K}_\mathbb{G}^{(\tau, c) + (r, s)}(X)$ described in Section 5.1. Once this reduction will be proved, the extension to the infinite rank case and the relation with the Fredholm operator formulation will follow straightforwardly as in Section 5.2 and Section 5.3 respectively.

5.5. Reduction to the original Freed-Moore K -theory. Let us start by considering the case of a trivial $\wp \equiv +1$. which is the case, for instance, when \mathbb{G} reduces to the trivial group. In this case any (ϖ, \wp, c, τ) -twisted $(\mathcal{E}, \eta, \rho, \gamma)$ bundle admits a decomposition

$$(\mathcal{E}, \eta, \rho, \gamma) = (\mathcal{E}^+, \eta^+ \equiv +\mathbb{1}, \rho^+, \gamma^+) \oplus (\mathcal{E}^-, \eta^- \equiv -\mathbb{1}, \rho^-, \gamma^-)$$

in which the sectors $(\mathcal{E}^\pm, \eta^\pm, \rho^\pm, \gamma^\pm)$, generated by projecting on the the eigenspaces of η are (ϖ, τ, c) -twisted bundles in the sense of Definition 5.2. Also the η -gradations have a corresponding decomposition and one finds that the ηK -theory in Definition 5.24 splits in the sum of two copies of the Freed-Moore K -theory of Definition 5.7, namely

$${}_{\eta}\mathcal{K}_{\mathbb{G}}^{(\tau, c, \wp \equiv +1) + (r, s)}(X) \simeq {}_{\varpi}\mathcal{K}_{\mathbb{G}}^{(\tau, c) + (r, s)}(X) \oplus {}_{\varpi}\mathcal{K}_{\mathbb{G}}^{(\tau, c) + (r, s)}(X). \quad (5.8)$$

The splitting phenomenon described by (5.8) is illustrated in the concrete (simple) example discussed in Appendix B. However, a reduction similar to (5.8) can be systematically extended to the general case.

To reduce the ηK -theory in Definition 5.24 to Freed-Moore K -theory of Definition 5.7, let us start with a recipe to produce data in Assumption 5.1 from those in Assumption 5.20. This recipe partly mimics the strategy of the proof of Proposition 5.4. In particular, let us recall that given the group \mathbb{G} which acts on the left on the space X one can form the extended group $\mathbb{G}' := \mathbb{G} \times \mathbb{Z}_2$. Also this group acts on X in the following way: Let $p : \mathbb{G}' \rightarrow \mathbb{G}$ be the first factor projection given by $p((g, \epsilon)) = g$ for all $g \in \mathbb{G}$ and $\epsilon \in \mathbb{Z}_2$, then $g'x = p(g')x$ for all $g' := (g, \epsilon) \in \mathbb{G}'$ and $x \in X$.

Definition 5.25 (Reduction of data). *Let $X, \mathbb{G}, \varpi, \wp, c, \tau$ be a set of data as in Assumption 5.20. Consider a new set of data $X', \mathbb{G}', \varpi', c', \tau'$ defined as follows:*

- *The base space is unchanged: $X' = X$;*
- *The new group is $\mathbb{G}' := \mathbb{G} \times \mathbb{Z}_2$ and acts on $X' = X$ via the first factor projection $p : \mathbb{G}' \rightarrow \mathbb{G}$;*
- *The new homomorphisms $\varpi' : \mathbb{G}' \rightarrow \mathbb{Z}_2$ and $c' : \mathbb{G}' \rightarrow \mathbb{Z}_2$ are given by composing ϕ and c with the projection p , i. e.*

$$\varpi' = \varpi \circ p, \quad c' = c \circ p;$$

- *The new 2-cocycle $\tau' \in Z_{\text{group}}^2(\mathbb{G}; C(X', \mathbb{U}(1))_{\varpi'})$ is given by*

$$\tau'((g_1, \epsilon_1), (g_2, \epsilon_2)) = \wp(g_2)^{\frac{1-\epsilon_1}{2}} \tau(g_1, g_2)$$

for all $g_1, g_2 \in \mathbb{G}$ and $\epsilon_1, \epsilon_2 \in \mathbb{Z}_2 = \{\pm 1\}$.

The new data $X', \mathbb{G}', \varpi', c', \tau'$ fulfill all the requirements of Assumption 5.1.

The reduction of data described in Definition 5.25 essentially consists in eliminating the map \wp which takes into account the presence of a Krein structure induced by η . As will be clarified in the next result, this reduction is compensated by the factor \mathbb{Z}_2 in the extended group $\mathbb{G} \times \mathbb{Z}_2$ that will be responsible for the specification of an indefinite inner product. In order to properly enunciate the Lemma below it is necessary to introduce a bit of terminology. Let $\rho' : \mathbb{G} \times \mathbb{Z}_2 \rightarrow \text{QS}(\mathfrak{H}_{\mathcal{E}}) \cap \text{Cov}(\mathcal{E})$ be a PUA-representation of a Hilbert bundle \mathcal{E} over X . Observe that in view of the group law one has that

$$\rho'((g, \epsilon)) = \rho'((g, +1))\rho'((e, \epsilon)), \quad g \in \mathbb{G}, \quad \epsilon \in \mathbb{Z}_2$$

where e is the unit of \mathbb{G} . One says that ρ' is a *reducible* PUA-representation if $\rho'((e, \epsilon)) = \pm \mathbb{1}$. If this is the case one has that $\rho'((g, \epsilon)) = \pm \rho'((g, +1)) =: \rho(g)$ where ρ is a PUA-representation of the group \mathbb{G} . It follows that the PUA-representation ρ' is called *non-reducible* when $\eta' := \rho'((e, -1)) \neq \pm \mathbb{1}$. Assume that the extended group $\mathbb{G} \times \mathbb{Z}_2$ acts on X only via the factor \mathbb{G} . In this case η' preserves the fibers of \mathcal{E} , namely $\eta' \in \text{Cov}(\mathcal{E})$. Moreover, the group law provides $(\eta')^2 = \rho'((e, +1)) = \mathbb{1}$. Then η' is a unitary involution, hence self-adjoint. The *non-reducible* PUA-representation ρ' is called *balanced* if in each fiber of \mathcal{E} the dimensions of the eigenspaces of η' related to the eigenvalues ± 1 have same (possibly infinite) dimension.

Lemma 5.26. *Let $X', \mathbb{G}', \varpi', c', \tau'$ be the set of data obtained from $X, \mathbb{G}, \varpi, \wp, c, \tau$ according to the recipe in Definition 5.25. Then, there is a natural one-to-one correspondence between (ϖ, \wp, c, τ) -twisted \mathbb{G} -equivariant Krein bundles on X with $C\ell^{r,s}$ -action and (ϖ', c', τ') -twisted \mathbb{G}' -equivariant*

Hilbert bundle on X (endowed with a non-reducible balanced PUA-representation ρ') with $\mathcal{C}l^{r,s}$ -action. Moreover, under this correspondence also gradations and η -gradations are in one-to-one correspondence.

Proof. (\Rightarrow) Let $(\mathcal{E}, \eta, \rho, \gamma)$ be a (ϖ, \wp, c, τ) -twisted \mathbb{G} -equivariant Krein bundles on X . Consider the triple $(\mathcal{E}', \rho', \gamma')$ defined as follows: $\mathcal{E}' := \mathcal{E}$, $\gamma' := \gamma$ and

$$\rho'((g, \epsilon)) := \rho(g)\eta^{\frac{1-\epsilon}{2}} = \begin{cases} \rho(g) & \text{if } \epsilon = +1 \\ \rho(g)\eta & \text{if } \epsilon = -1. \end{cases}$$

A direct check shows that $(\mathcal{E}', \rho', \gamma')$ is a (ϖ', c', τ') -twisted \mathbb{G}' -equivariant Hilbert bundle on X . Moreover $(\mathcal{E}, \eta, \rho, \gamma)$ and $(\mathcal{E}', \rho', \gamma')$ have by construction the same Clifford action. If Γ is an η -gradation of $(\mathcal{E}, \eta, \rho, \gamma)$ then the condition $\eta\Gamma = \Gamma\eta$ assures that

$$\rho'((g, \epsilon)) = \Gamma\rho(g)\eta^{\frac{1-\epsilon}{2}} = c(g)\Gamma\eta^{\frac{1-\epsilon}{2}} = c'((g, \epsilon))\rho'((g, \epsilon))\Gamma$$

and this proves that Γ is automatically a gradation of $(\mathcal{E}', \rho', \gamma')$.

(\Leftarrow) Now the converse implication. Let $(\mathcal{E}, \rho', \gamma)$ be a (ϖ', c', τ') -twisted $\mathbb{G} \times \mathbb{Z}_2$ -equivariant Hilbert bundle on X such that the action of the factor \mathbb{Z}_2 is trivial on X and the maps (ϖ', c', τ') are related to the maps (ϖ, c, τ) defined on \mathbb{G} according to the prescription in Definition 5.25. Let $\eta := \rho'((e, -1))$. Since the factor \mathbb{Z}_2 acts trivially on X it follows that η preserves the fibers of \mathcal{E} . Since the PUA-representation ρ' is non-reducible and balanced the operator η endows \mathcal{E} with a Krein structure. The PUA-representation ρ' can be represented as

$$\rho'((g, \epsilon)) = \rho(g)\eta^{\frac{1-\epsilon}{2}}$$

where $\rho(g) := \rho'((g, +1))$ provides a PUA-representation of \mathbb{G} . A direct check shows that $(\mathcal{E}, \eta, \rho, \gamma)$ is a (ϖ, \wp, c, τ) -twisted \mathbb{G} -equivariant Krein bundles on X . If Γ is a gradation of $(\mathcal{E}, \rho', \gamma)$ then

$$\Gamma\rho'((g, \epsilon)) = c'((g, \epsilon))\rho'((g, \epsilon))\Gamma = c(g)\rho(g)\eta^{\frac{1-\epsilon}{2}}\Gamma.$$

The same computation for the group element $(e, -1)$ provides $\Gamma\eta = \eta\Gamma$ which in turn implies $\Gamma\rho(g) = c(g)\rho(g)\Gamma$. This proves that Γ behaves like an η -gradation of $(\mathcal{E}, \eta, \rho, \gamma)$. ■

The correspondence established in Lemma 5.26 immediately leads to the reduction anticipated at the beginning of this Section.

Theorem 5.27 (*K-theory reduction*). *For a given data set $\mathbb{G}, \varpi, \wp, c, \tau$ as in Assumption 5.20 there is a natural isomorphism of K-theories*

$$\varpi_{\eta} \mathcal{K}_{\mathbb{G}}^{(\tau, c, \wp) + (r, s)}(X) \simeq \varpi' \mathcal{K}_{\mathbb{G}'}^{(\tau', c') + (r, s)}(X),$$

where the group \mathbb{G}' , the homomorphisms ϖ', c' and the 2-cocycle τ' are obtained from the original data according to the prescription described in Definition 5.25.

Proof. In view of Lemma 5.26, one has the identification of monoids

$$\varpi_{\eta} \mathcal{M}_{\mathbb{G}}^{(\tau, c, \wp) + (r, s)}(X) \simeq \varpi' \mathcal{M}_{\mathbb{G}'}^{(\tau', c') + (r, s)}(X), \quad \varpi_{\eta} \mathcal{Z}_{\mathbb{G}}^{(\tau, c, \wp) + (r, s)}(X) \simeq \varpi' \mathcal{Z}_{\mathbb{G}'}^{(\tau', c') + (r, s)}(X).$$

This leads to the isomorphism in the theorem. ■

5.6. Explicit computations in some special cases. The aim of this Section is to provide the explicit computation of the groups $\varpi_{\eta} \mathcal{K}_{\mathbb{G}}^{(\tau, c, \wp) + (r, s)}(X)$ in some simple cases making use of the reduction described in Theorem 5.27. Let us assume that: (i) $X = \{*\}$ consists of a single point, (ii) $\mathbb{G} = \mathbb{Z}_2$ and (iii) $\varpi \equiv +1$ is the trivial homomorphism. Under these three assumptions the group cohomology turns out to be trivial, i. e. $H_{\text{group}}^2(\mathbb{Z}_2; \mathbb{U}(1)) = 0$, so one can assume that $\tau \equiv +1$ is also trivial. The three assumptions above describe the situation in which one wants to classify \mathcal{C} -symmetric η -self-adjoint gapped operators defined on a given (possibly finite dimensional) Hilbert space \mathcal{H} on which the group action of \mathbb{Z}_2 is implemented by a single linear operator $\rho(-1)$. In the simplest possible case $\mathcal{H} = \mathbb{C}^2$ and the group-action of \mathbb{Z}_2 is trivial, i. e. $\rho(-1) = \mathbb{1}$. This situation is described explicitly in Appendix

B. The case of a general Hilbert space \mathcal{H} with a possible non-trivial group-action but with a trivial homomorphism $\wp \equiv +1$ has been described through the isomorphism (5.8). In this Section we focus on the case of a non-trivial group-action with $\wp : \mathbb{Z}_2 \rightarrow \mathbb{Z}_2$ the identity map, i. e. $\wp(\pm 1) = \pm 1$. This means that the action of \mathbb{Z}_2 on \mathcal{H} is implemented by a linear unitary operator $R := \rho(-1)$ such that $R^2 = \mathbb{1}$ and $R\eta = -\eta R$, namely by an η -reflecting operator. The only remaining choice is about c .

Under the assumptions above one is left with the computation of the groups ${}^{+1}_\eta \mathcal{K}_{\mathbb{Z}_2}^{(+1,c,\text{Id})+(r,s)}(\{*\})$ where the constraints $\varpi \equiv +1, \tau \equiv +1$ and $\wp \equiv \text{Id}$ have been explicitly indicated. By effect of the reduction in Theorem 5.27 one obtains the isomorphisms

$${}^{+1}_\eta \mathcal{K}_{\mathbb{Z}_2}^{(+1,c,\text{Id})+(r,s)}(\{*\}) \simeq {}^{+1} \mathcal{K}_{\mathbb{Z}_2 \times \mathbb{Z}_2}^{(\tau',c',\cdot)+(r,s)}(\{*\}) \simeq \mathcal{K}_{\mathbb{Z}_2 \times \mathbb{Z}_2}^{(\tau',c')^n}(\{*\}), \quad n = s - r \pmod{2}$$

where in the last isomorphism the periodicity and the notation of Proposition 5.9 have been used. The maps τ' and c' are defined according to the prescription of Definition 5.25. More precisely one has that $c' : \mathbb{Z}_2 \times \mathbb{Z}_2 \rightarrow \mathbb{Z}_2$ is given by

$$c'((\zeta, \epsilon)) = c(\zeta), \quad (\zeta, \epsilon) \in \mathbb{Z}_2 \times \mathbb{Z}_2. \quad (5.9)$$

The 2-cocycle $\tau' : (\mathbb{Z}_2 \times \mathbb{Z}_2) \times (\mathbb{Z}_2 \times \mathbb{Z}_2) \rightarrow \mathbb{U}(1)$ is described by the following formula:

$$\tau'((\zeta_1, \epsilon_1), (\zeta_2, \epsilon_2)) = \wp(\zeta_2)^{\frac{1-\epsilon_1}{2}} \tau(\zeta_1, \zeta_2) = \zeta_2^{\frac{1-\epsilon_1}{2}}, \quad (\zeta_j, \epsilon_j) \in \mathbb{Z}_2 \times \mathbb{Z}_2.$$

This 2-cocycle is non-trivial and its values are summarized in Table 5.2.

τ'	(+1, +1)	(-1, +1)	(+1, -1)	(-1, -1)
(+1, +1)	+1	+1	+1	+1
(-1, +1)	+1	+1	+1	+1
(+1, -1)	+1	-1	+1	-1
(-1, -1)	+1	-1	+1	-1

TABLE 5.2. Values of the 2-cocycle τ' .

Interestingly, the presence of the Krein metric η translates in the effects of the non-trivial twist τ' in the computation of the Freed-Moore K -theory $\mathcal{K}_{\mathbb{Z}_2 \times \mathbb{Z}_2}^{(\tau',c')^n}(\{*\})$.

In view of Remark 5.18 the Freed-Moore K -groups can be expressed in terms of the twisted K -theory for skew-adjoint Fredholm operators. More precisely one has that

$$\mathcal{K}_{\mathbb{Z}_2 \times \mathbb{Z}_2}^{(\tau',c')^n}(\{*\}) \simeq \mathbb{K}_{\mathbb{Z}_2 \times \mathbb{Z}_2}^{(\tau',c')^n}(\{*\})$$

where one has tacitly used that $\hat{\tau}' = \tau'$. Let us compute the latter groups in the two cases given by the different choices of c . Definition (5.9) shows that c' is trivial if and only if also c is trivial.

Lemma 5.28. *Regardless of the choice of c' it holds true that*

$$\mathbb{K}_{\mathbb{Z}_2 \times \mathbb{Z}_2}^{(\tau',c')^{-n}}(\{*\}) \simeq K_{\mathbb{C}}^{-n}(\{*\})$$

where

$$K_{\mathbb{C}}^{-n}(\{*\}) \simeq \begin{cases} \mathbb{Z} & \text{if } n \text{ even} \\ 0 & \text{if } n \text{ odd} \end{cases}$$

is the usual complex K -theory of the single-point space $\{*\}$.

Proof (sketch of). ($c' \equiv +1$). Assume that c' is trivial. As a consequence $\mathbb{K}_{\mathbb{Z}_2 \times \mathbb{Z}_2}^{(\tau', +1) - n}(\{*\})$ is a usual twisted K -theory group. Then one can apply the known decomposition

$$\mathbb{K}_{\mathbb{Z}_2 \times \mathbb{Z}_2}^{(\tau', +1) - n}(\{*\}) \simeq \mathcal{R}_{\tau'}(\mathbb{Z}_2 \times \mathbb{Z}_2) \otimes K_{\mathbb{C}}^{-n}(\{*\})$$

where $\mathcal{R}_{\tau'}(\mathbb{Z}_2 \times \mathbb{Z}_2)$ is the free abelian group generated by the τ' -projective complex representations of $\mathbb{Z}_2 \times \mathbb{Z}_2$. Such projective representations are unique up to isomorphisms, since they are in one-to-one correspondence with ungraded complex representations of the complexified Clifford algebra $\mathbb{C}\ell^2 := C\ell^{0,2} \otimes \mathbb{C}$. Hence $\mathcal{R}_{\tau'}(\mathbb{Z}_2 \times \mathbb{Z}_2) \simeq \mathbb{Z}$ and one finally gets

$$\mathbb{K}_{\mathbb{Z}_2 \times \mathbb{Z}_2}^{(\tau', +1) - n}(\{*\}) \simeq K_{\mathbb{C}}^{-n}(\{*\}) .$$

($c' \not\equiv +1$). If c' is non-trivial then the PUA-representation of $\mathbb{Z}_2 \times \mathbb{Z}_2$ is generated by $R := \rho((-1, +1))$ which is odd and by $\eta := \rho((+1, -1))$ which is even. Let us introduce a different pair of generators $\tilde{\gamma}_1 := R$ and $\tilde{\gamma}_2 := R\eta$ which are both odd. In view of the twisting τ' one has that $R\eta = -\eta R$ and this implies that

$$\tilde{\gamma}_1^2 = \mathbb{1} , \quad \tilde{\gamma}_2^2 = -\mathbb{1} , \quad \tilde{\gamma}_1 \tilde{\gamma}_2 = -\tilde{\gamma}_2 \tilde{\gamma}_1 .$$

Therefore a (τ', c') -twisted representation of $\mathbb{Z}_2 \times \mathbb{Z}_2$ with $C\ell^{r,s}$ -action amounts to a \mathbb{Z}_2 -graded representation of the complexified Clifford algebra $\mathbb{C}\ell^{r+s+2} := C\ell^{r+1, s+1} \otimes \mathbb{C}$. The K -theory $\mathbb{K}_{\mathbb{Z}_2 \times \mathbb{Z}_2}^{(\tau', c') - (s-r)}(\{*\})$ is the quotient of the free abelian monoid generated by \mathbb{Z}_2 -graded $\mathbb{C}\ell^{r+s+2}$ -modules by the submonoid generated by those which can be extended to \mathbb{Z}_2 -graded $\mathbb{C}\ell^{r+s+3}$ -modules. This group is exactly $K_{\mathbb{C}}^{r+s+2}(\{*\})$. Thus, by using the Bott periodicity one gets

$$\mathbb{K}_{\mathbb{Z}_2 \times \mathbb{Z}_2}^{(\tau', c') - n}(\{*\}) := \mathbb{K}_{\mathbb{Z}_2 \times \mathbb{Z}_2}^{(\tau', c') - (s-r)}(\{*\}) \simeq K_{\mathbb{C}}^{r+s+2}(\{*\}) \simeq K_{\mathbb{C}}^{(r-s)+2(s+1)}(\{*\}) \simeq K_{\mathbb{C}}^{-n}(\{*\}) .$$

The proof is completed. ■

In conclusion, under the simplified hypotheses (i), (ii) and (iii) above, and using the various isomorphisms and the computation in Lemma 5.28 one obtains that

$${}^{+1}_{\eta} \mathcal{K}_{\mathbb{Z}_2}^{(+1, c, \text{Id}) + (r, s)}(\{*\}) \simeq K_{\mathbb{C}}^{-(s-r)}(\{*\}) , \quad \wp \equiv \text{Id} . \quad (5.10)$$

This result must be compared with the splitting (5.8) which provides

$${}^{+1}_{\eta} \mathcal{K}_{\mathbb{Z}_2}^{(+1, c, +1) + (r, s)}(\{*\}) \simeq \mathcal{K}_{\mathbb{Z}_2}^{(+1, c) + (r, s)}(\{*\}) \oplus \mathcal{K}_{\mathbb{Z}_2}^{(+1, c) + (r, s)}(\{*\})$$

in the case of a trivial homomorphism \wp . Since $\mathcal{K}_{\mathbb{Z}_2}^{(+1, c) + (r, s)}(\{*\}) \simeq \mathbb{K}_{\mathbb{Z}_2}^{(+1, c) - (s-r)}(\{*\})$ and

$$\mathbb{K}_{\mathbb{Z}_2}^{(+1, c) - (s-r)}(\{*\}) \simeq K_{\mathbb{C}}^{-(s-r) - \delta(c)}(\{*\})$$

where

$$\delta(c) := \begin{cases} 0 & \text{if } c \equiv +1 \\ 1 & \text{if } c \equiv \text{Id} , \end{cases}$$

one concludes that

$${}^{+1}_{\eta} \mathcal{K}_{\mathbb{Z}_2}^{(+1, c, +1) + (r, s)}(\{*\}) \simeq K_{\mathbb{C}}^{-(s-r) - \delta(c)}(\{*\}) \oplus K_{\mathbb{C}}^{-(s-r) - \delta(c)}(\{*\}) , \quad \wp \equiv +1$$

which is valid when \wp acts trivially.

APPENDIX A. MAXWELL-TYPE OPERATORS AND METAMATERIALS

This Section is devoted to the presentation of a particular family of η -self-adjoint models that is finding use in the recent physical literature. The toy model m_w described in equation (1.2) is a particularly simple representative of this family.

Let M_0 be a (possibly unbounded) self-adjoint operator on the complex Hilbert space \mathcal{H} with dense domain \mathcal{D}_0 and spectrum $\sigma(M_0) \subseteq \mathbb{R}$. In many physical situations it can be useful to endow \mathcal{H} with a complex structure given by the anti-unitary involution C . In this case it is natural to require that M_0 respect the given complex structure:

Assumption A.1. *The complex structure preserves the domain of M_0 , namely $C[\mathcal{D}_0] \subseteq \mathcal{D}_0$, and M_0 is even ($CM_0C = M_0$) or odd ($CM_0C = -M_0$) with respect to C .*

Assumption A.1 has consequences on the form of the spectrum of M_0 . Let us focus on the odd case:

Proposition A.2. *Under the odd case of Assumption A.1 the spectrum of M_0 is symmetric around zero, namely $\omega \in \sigma(M_0)$ if and only if $-\omega \in \sigma(M_0)$.*

Proof. Let $\omega \in \rho(M_0) = \mathbb{C} \setminus \sigma(M_0)$ be a point in the resolvent set. Then

$$R_0(\omega) := \frac{1}{M_0 - \omega \mathbb{1}}$$

exists as a bounded operator. Since $CR_0(\omega)C = -R_0(-\omega)$ one concludes that also $-\omega \in \rho(M_0)$. The proof is concluded by observing that $\sigma(M_0) = \mathbb{C} \setminus \rho(M_0)$. ■

Definition A.3 (Metamaterial weight). *A metamaterial weight on \mathcal{H} is a bounded self-adjoint operator W that satisfies $0 \notin \sigma(W)$. The weight is called real if $CWC = W$.*

Proposition A.4. *Let W be a metamaterial weight. Then there is a pair of bounded self-adjoint operators (W_+, W_-) and non negative constants $0 \leq a_{\pm} < +\infty$ such that:*

- (a) $W = W_+ + W_-$ and $W_{\pm}W_{\mp} = 0$;
- (b) $0 < \pm W_{\pm} \leq a_{\pm} \mathbb{1}$.

Proof. Since the spectrum of a self-adjoint operator is closed there are positive constants $0 < b_{\pm} \leq a_{\pm} < +\infty$ such that

$$\sigma(W) \subseteq [-a_-, -b_-] \cup [b_+, a_+].$$

By spectral calculus one can define the spectral projections $\Pi_- := \chi_{[-a_-, -b_-]}(W)$ and $\Pi_+ := \chi_{[b_+, a_+]}(W)$. Then the claim follows by setting $W_{\pm} := \Pi_{\pm} W \Pi_{\pm}$. ■

From Definition A.3 it follows that W^{-1} exists as a bounded self-adjoint operator. Therefore one can associate with W the operator

$$\eta_W := \frac{W}{|W|} = \frac{W_+}{|W|} + \frac{W_-}{|W|}. \quad (\text{A.1})$$

It turns out that η_W has the properties of a fundamental symmetry according to Remark 2.2.

Let us focus the attention on the family of operators of the following type:

Definition A.5 (Metamaterial Maxwell-type operator). *Let M_0 be a self-adjoint operator with dense domain \mathcal{D}_0 which meets Assumption A.1. Let W be a metamaterial weight in the sense of Definition A.3. The operator*

$$M := W M_0 = \eta_W \tilde{M} \quad \text{with} \quad \tilde{M} := |W| M_0 \quad (\text{A.2})$$

is called a metamaterial Maxwell-type operator.

From its very definition it follows that M is densely defined with domain

$$\mathcal{D}(M) = \mathcal{D}(M_0) = \mathcal{D}_0.$$

Remark A.6 (Maxwell-type operator). The general notion of Maxwell-type operator has been discussed in [DL3, Section 6]. The name comes from the fact that the Maxwell's equations for the propagation of the light in a medium can be recast in the form (A.2) [Wil, Kat, BS] (see also [DL1, DL2, DL3]). In this specific case the *free* operator, acting the Hilbert space $L^2(\mathbb{R}^3, \mathbb{C}^6)$, is

$$M_0 := \begin{pmatrix} 0 & +i \nabla^{\times} \\ -i \nabla^{\times} & 0 \end{pmatrix} \quad (\text{A.3})$$

where ∇^{\times} denotes the *curl* operator. The material weight is described by the matrix-valued function

$$W(x) := \begin{pmatrix} \varepsilon^{-1}(x) & 0 \\ 0 & \mu^{-1}(x) \end{pmatrix}$$

where $\varepsilon(x)$ is the electric permittivity tensor and $\mu(x)$ is the magnetic permeability tensor. In the “standard” definition of Maxwell-type operator the operator W is required to be positive (cf. [DL3, Definition 6.1]). On the contrary, Definition A.3 allows W to have also a negative part. For instance this is the case which models *single-negative metamaterials* [FR, AE, JCHZ3] in which the electric permittivity tensor is positive $\varepsilon > 0$ and the magnetic permeability tensor is negative $\mu < 0$ (or viceversa). This justifies the use of the expression *metamaterial* in Definition A.3 and Definition A.5. Finally, let us notice that the operator m_w in (1.2) is an example of metamaterial Maxwell-type operator. ◀

Let \mathcal{H}_* be the Hilbert space obtained by endowing \mathcal{H} with the weighted scalar product

$$\langle \phi, \varphi \rangle_* := \langle \phi, |W|^{-1} \varphi \rangle. \quad (\text{A.4})$$

The equation

$$\langle \phi, \eta_W \varphi \rangle_* = \langle \phi, W|W|^{-2} \varphi \rangle = \langle \eta_W \phi, |W|^{-1} \varphi \rangle = \langle \eta_W \phi, \varphi \rangle_* \quad (\text{A.5})$$

shows that η_W is self-adjoint, hence a fundamental symmetries, also in the weighted Hilbert space \mathcal{H}_* . Therefore η_W induces a Krein space structure $(\mathcal{H}_*, \langle \langle \cdot, \cdot \rangle \rangle_{\eta_W})$ by means of the indefinite inner product

$$\langle \langle \phi, \varphi \rangle \rangle_{\eta_W} := \langle \phi, \eta_W \varphi \rangle_* = \langle \phi, W^{-1} \varphi \rangle, \quad \forall \phi, \varphi \in \mathcal{H}_*.$$

Proposition A.7. *Let $M = WM_0$ be a metamaterial Maxwell-type operator in the sense of Definition A.5. Then:*

- (1) *The operator $\tilde{M} := |W|M_0$ is self-adjoint on the weighted Hilbert space \mathcal{H}_* with dense domain \mathcal{D}_0 .*
- (2) *The operator $M = \eta_W \tilde{M}$ is η_W -self-adjoint on the Krein space $(\mathcal{H}_*, \langle \langle \cdot, \cdot \rangle \rangle_{\eta_W})$.*

Proof. (1) is a particular case of [DL3, Proposition 6.2]. (2) follows from the self-adjointness of \tilde{M} and η_W with respect to the Hilbert structure of \mathcal{H}_* . This implies that the adjoint of M is $\tilde{M}\eta_W = \eta_W M\eta_W$. The last relation is equivalent to the η_W -self-adjointness of M . ■

The last result is preparatory for the following natural question: *When a metamaterial Maxwell-type operator admits a C-symmetry Ξ ?* A set of sufficient conditions, borrowed from [KSu, Proposition], is described in the following:

Theorem A.8. *Let $M = WM_0$ be a metamaterial Maxwell-type operator in the sense of Definition A.5. Assume that: (a) M_0 is a non negative operator on \mathcal{H} and $0 \notin \text{Ker}(M_0)$; (b) $\rho(M) \neq \emptyset$; (c) 0 and ∞ are not singular critical points of M . Then, the η_W -self-adjoint operator M admits the C-symmetry $\Xi := E_M(\mathbb{R}_+) - E_M(\mathbb{R}_-)$ where E_M denotes the projection-valued spectral measure of M .*

Proof. Condition (a) is equivalent to $\langle \langle \varphi, M\varphi \rangle \rangle_{\eta_W} \geq 0$ for all $\varphi \in \mathcal{H}_*$ meaning that M is η_W -non negative. Together with condition (b) this is enough to prove that $\sigma(M) \subseteq \mathbb{R}$ [AI, Chapter 2, Theorem 3.27]. Moreover, conditions (a) and (b) imply that M is *definitizable* in the sense of [Lan2, Section I.3]. Moreover the only possible critical points of M are 0 and ∞ [Lan2, Section II.6]. The second part of condition (a) along with the invertibility of W assures that 0 is not an eigenvalue of M . If 0 and ∞ are not singular critical points then [Lan2, Theorem 5.7] assures the existence of a spectral functional calculus for M and the spectral projections $E_M(\mathbb{R}_+)$ and $E_M(\mathbb{R}_-)$ are well defined. The operator $\Xi := E_M(\mathbb{R}_+) - E_M(\mathbb{R}_-)$ verifies all the properties of a C-symmetry for M . ■

Theorem A.8 provides a set of criterions for a metamaterial Maxwell-type operator to describe a dynamically stable system in the sense of Definition 1.1. However condition (a) is quite unpleasant since it excludes the case of the Maxwell equations described in Remark A.6. Indeed, the operator M_0 described by (A.3) has spectrum $\sigma(M_0) = \mathbb{R}$ and does not meet condition (a). However, it is odd in the sense of Assumption A.1. This leads to the following:

Open problem. *Is it possible to prove that a metamaterial Maxwell-type operator $M = WM_0$ admits a C-symmetry Ξ under the odd case of Assumption A.1? If so, what other conditions are necessary?*

APPENDIX B. SOME EXPLICIT COMPUTATION IN DIMENSION TWO

In this Section we focus on the simplest (finite dimensional) Krein space given by the Hilbert space \mathbb{C}^2 endowed with the metric operator

$$\eta := \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}. \quad (\text{B.1})$$

The canonical basis $v_1 := (1, 0)$ and $v_2 := (0, 1)$ diagonalizes η and the usual complex conjugation C is compatible with η in view of the fact that $Cv_j = v_j$ for $j = 1, 2$. The η -reflecting symmetry is given by

$$R := \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}.$$

The following relations hold true: $C\eta = \eta C$, $R\eta = \eta R$ and $R\eta = -\eta R$. In order to determine

$$\text{QS}_\eta(\mathbb{C}^2) := \text{U}_\eta(\mathbb{C}^2) \sqcup \text{AU}_\eta(\mathbb{C}^2) \sqcup \text{PU}_\eta(\mathbb{C}^2) \sqcup \text{PAU}_\eta(\mathbb{C}^2)$$

it is enough to compute $\text{U}_\eta(\mathbb{C}^2)$ in view of the relations (2.11), (2.15) and (2.16).

Lemma B.1. *Any element of $\text{U}_\eta(\mathbb{C}^2)$ can be parametrized as*

$$U(r, \alpha, \beta, \delta) := \begin{pmatrix} e^{i\alpha} \sqrt{1+r^2} & e^{i(\beta+\delta)} r \\ e^{i(\alpha-\delta)} r & e^{i\beta} \sqrt{1+r^2} \end{pmatrix}, \quad r \in [0, +\infty), \quad \alpha, \beta, \delta \in [-\pi, +\pi].$$

Moreover, $\|U(r, \alpha, \beta, \delta)\| = r + \sqrt{1+r^2}$.

Proof. Given a complex matrix

$$U := \begin{pmatrix} a & b \\ c & d \end{pmatrix}, \quad a, b, c, d \in \mathbb{C}$$

the condition $U^* \eta U = \eta$, which ensures $U \in \text{U}_\eta(\mathbb{C}^2)$, reads as

$$\begin{pmatrix} |a|^2 - |c|^2 & \bar{a}b - \bar{c}d \\ \bar{a}b - \bar{c}d & |b|^2 - |d|^2 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

From $\bar{a}b = \bar{c}d$ one gets $|a|^2|b|^2 = |c|^2|d|^2$. By inserting $|a|^2 = |c|^2 + 1$ and $|d|^2 = |b|^2 + 1$ in the previous equation one finally gets $|b| = |c| =: r \geq 0$. This implies that $|a| = |d| = \sqrt{1+r^2}$ and so there are two angles α and β such that

$$a := e^{i\alpha} \sqrt{1+r^2}, \quad \text{and} \quad d := e^{i\beta} \sqrt{1+r^2}.$$

Let $b = e^{i\delta_b} r$ and $c = e^{i\delta_c} r$. The condition $\bar{a}b = \bar{c}d$ implies that $\delta_b + \delta_c = \alpha + \beta$. The last condition is satisfied by $\delta_b := \beta + \delta$ and $\delta_c := \alpha - \delta$ for some angle δ . This completes the proof of the parametrization of $U(r, \alpha, \beta, \delta)$. From the parametrization it follows that

$$U(r, \alpha, \beta, \delta) = H(r, \delta) D(\alpha, \beta)$$

where

$$H(r, \delta) := \begin{pmatrix} \sqrt{1+r^2} & e^{+i\delta} r \\ e^{-i\delta} r & \sqrt{1+r^2} \end{pmatrix}, \quad \text{and} \quad D(\alpha, \beta) := \begin{pmatrix} e^{i\alpha} & 0 \\ 0 & e^{i\beta} \end{pmatrix}.$$

The matrix $H(r, \delta)$ is self-adjoint while $D(\alpha, \beta)$ is unitary. This implies that

$$\|U(r, \alpha, \beta, \delta)\| = \|H(r, \delta)D(\alpha, \beta)\| = \|H(r, \delta)\| = r + \sqrt{1+r^2}$$

where $r + \sqrt{1+r^2}$ is the maximum eigenvalue of $H(r, \delta)$. ■

In the next result the \mathbb{R} -Banach space $\text{H}_\eta(\mathbb{C}^2)$ of the η -self-adjoint operators is described.

Lemma B.2. *Any element of $\text{H}_\eta(\mathbb{C}^2)$ can be parametrized as*

$$H(x_1, x_2, y, z) := \begin{pmatrix} x_1 & y + iz \\ -y + iz & x_2 \end{pmatrix}, \quad x_1, x_2, y, z \in \mathbb{R}.$$

Proof. Given a complex matrix

$$H := \begin{pmatrix} a & b \\ c & d \end{pmatrix}, \quad a, b, c, d \in \mathbb{C} \quad (\text{B.2})$$

the condition $\eta H^* \eta = H$, which ensures $H \in \mathbb{H}_\eta(\mathbb{C}^2)$, reads as

$$\begin{pmatrix} \bar{a} & -\bar{c} \\ -\bar{b} & \bar{d} \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix}.$$

The three relations $\bar{a} = a$, $\bar{d} = d$ and $\bar{b} = -c$ can be solved in terms of four real parameters $x_1, x_2, y, z \in \mathbb{R}$ as follows

$$a := x_1 \quad d := x_2 \quad b := y + iz \quad c := -y + iz$$

This completes the proof. \blacksquare

Also the space $\mathbb{C}_\eta(\mathbb{C}^2)$ of the C -symmetries and the space $\mathbb{R}_\eta(\mathbb{C}^2)$ of the self-adjoint operators anti-commuting with the metric η can be easily described.

Lemma B.3. *Any element of $\mathbb{C}_\eta(\mathbb{C}^2)$ can be parametrized as*

$$\Xi(r, \theta) := \begin{pmatrix} \sqrt{1+r^2} & re^{i\theta} \\ -re^{-i\theta} & -\sqrt{1+r^2} \end{pmatrix}, \quad r \in [0, +\infty), \quad \theta \in [-\pi, +\pi]. \quad (\text{B.3})$$

Moreover, $\|\Xi(r, \theta)\| = r + \sqrt{1+r^2}$. Similarly, any element of $\mathbb{R}_\eta(\mathbb{C}^2)$ can be parametrized as

$$Q(r, \theta) := \begin{pmatrix} 0 & re^{i\theta} \\ re^{-i\theta} & 0 \end{pmatrix}, \quad r \in [0, +\infty), \quad \theta \in [-\pi, +\pi] \quad (\text{B.4})$$

and $\|Q(r, \theta)\| = r$.

Proof. Let $\Xi \in \mathbb{C}_\eta(\mathbb{C}^2)$. In view of (4.2) one has that Ξ is η -self-adjoint and so it can be represented as

$$\Xi := \begin{pmatrix} x_1 & y + iz \\ -y + iz & x_2 \end{pmatrix}, \quad x_1, x_2, y, z \in \mathbb{R}.$$

according to Lemma B.2. The involutive property $\Xi^2 = \mathbb{1}$ reads as

$$\begin{pmatrix} x_1^2 - y^2 - z^2 & (x_1 + x_2)(y + iz) \\ (x_1 + x_2)(-y + iz) & x_2^2 - y^2 - z^2 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}.$$

The relations $x_1^2 - y^2 - z^2 = 1 = x_2^2 - y^2 - z^2$ imply that $|x_1| = |x_2| \neq 0$. Let us set $y + iz := re^{i\theta}$ and $-y + iz := -re^{-i\theta}$. The conditions $(x_1 + x_2)(y + iz) = 0 = (x_1 + x_2)(-y + iz)$ imply $|(x_1 + x_2)r| = 0$. In the case $r = 0$ the matrix Ξ is diagonal and $|x_1| = 1 = |x_2|$. The condition $\eta\Xi > 0$ implies $x_1 > 0$ and $-x_2 > 0$. As a result the condition $r = 0$ implies $\Xi = \eta$. Let us assume now that $r \neq 0$. This implies that $x_2 = -x_1$ and $x_1 = \sqrt{1+r^2}$. The matrix $\eta\Xi$ reads

$$\eta\Xi := \begin{pmatrix} \sqrt{1+r^2} & re^{i\theta} \\ re^{-i\theta} & -\sqrt{1+r^2} \end{pmatrix}.$$

The eigenvalues of the self-adjoint matrix $\eta\Xi$ are $\sqrt{1+r^2} \pm r > 0$ showing that $\eta\Xi$ is positive. Moreover, the unitarity of η implies $\|\Xi\| = \|\eta\Xi\| = r + \sqrt{1+r^2}$.

Since the elements of $\mathbb{R}_\eta(\mathbb{C}^2)$ are self-adjoint one has

$$Q := \begin{pmatrix} x_1 & re^{i\theta} \\ re^{-i\theta} & x_2 \end{pmatrix}$$

with $x_1, x_2 \in \mathbb{R}$. The condition $\eta Q = -Q\eta$ forces $x_1 = x_2 = 0$. The eigenvalues of Q are $\pm r$ and this shows that $\|Q\| = r$. \blacksquare

Starting from (B.4) one can compute

$$\eta e^Q = \begin{pmatrix} \cosh(r) & e^{i\theta} \sinh(r) \\ -e^{-i\theta} \sinh(r) & -\cosh(r) \end{pmatrix} = \begin{pmatrix} \sqrt{1 + \sinh(r)^2} & e^{i\theta} \sinh(r) \\ -e^{-i\theta} \sinh(r) & -\sqrt{1 + \sinh(r)^2} \end{pmatrix}$$

where $\cosh(r) := \frac{1}{2}(e^r + e^{-r})$ and $\sinh(r) := \frac{1}{2}(e^r - e^{-r})$ are the hyperbolic functions that satisfy $\cosh(r)^2 - \sinh(r)^2 = 1$. The change of variable $r' := \sinh(r)$ and the comparison with (B.3) shows that $\eta e^Q \in C_\eta(\mathbb{C}^2)$ in accordance with the general result proved in Lemma 4.2. Moreover, as a consequence of Lemma B.3 one has that $C_\eta(\mathbb{C}^2) \simeq \mathbb{C} \simeq R_\eta(\mathbb{C}^2)$.

We are now in position to study the set of η -self adjoint operators which admit a C -symmetry.

Lemma B.4. *Let $\Xi \in C_\eta(\mathbb{C}^2)$ be a C -symmetry and $H \in H_\eta(\mathbb{C}^2)$ an η -self adjoint operator. The condition $H\Xi = \Xi H$ holds true if and only if H is of the form*

$$H = u \mathbb{1} + v \Xi, \quad u, v \in \mathbb{R}. \quad (\text{B.5})$$

Proof. Clearly operators of the form (B.5) are η -self-adjoint since $\mathbb{1}$ and Ξ are η -self-adjoint and commute with Ξ . To prove the converse let us use the parametrization (B.2) and (B.3) for H and Ξ , respectively. A straightforward computation shows that

$$H\Xi = \begin{pmatrix} x_1 \sqrt{1+r^2} - (y+iz)re^{-i\theta} & x_1 re^{i\theta} - (y+iz)\sqrt{1+r^2} \\ (-y+iz)\sqrt{1+r^2} - x_2 re^{-i\theta} & -x_2 \sqrt{1+r^2} + (-y+iz)re^{i\theta} \end{pmatrix}$$

and

$$\Xi H = \begin{pmatrix} x_1 \sqrt{1+r^2} + (-y+iz)re^{i\theta} & x_2 re^{i\theta} + (y+iz)\sqrt{1+r^2} \\ -(-y+iz)\sqrt{1+r^2} - x_1 re^{-i\theta} & -x_2 \sqrt{1+r^2} - (y+iz)re^{-i\theta} \end{pmatrix}.$$

The commutation relations between H and Ξ imply the following conditions

$$(y+iz) = (x_1 - x_2) \frac{r}{2\sqrt{1+r^2}} e^{i\theta}, \quad (-y+iz) = -(y+iz)e^{-i2\theta} = (x_2 - x_1) \frac{r}{2\sqrt{1+r^2}} e^{-i\theta}.$$

Let us define $u := \frac{x_1+x_2}{2}$ and $v := \frac{x_1-x_2}{2\sqrt{1+r^2}}$. With this notation one has that $(y+iz) = vre^{i\theta}$, $x_1 = u + v\sqrt{1+r^2}$ and $x_2 = u - v\sqrt{1+r^2}$. After plugging in these expressions in (B.2) one obtains the representation (B.5). ■

Let us compare the \mathbb{R} -Banach space $H_\eta(\mathbb{C}^2)$ of the η -self-adjoint operators with its subspace $H_{\eta,\Xi}(\mathbb{C}^2)$ of the η -self-adjoint operators commuting with the C -symmetry Ξ . Lemma B.2 shows that $H_\eta(\mathbb{C}^2) \simeq \mathbb{R}^4$ while Lemma B.4 proves that $H_{\eta,\Xi}(\mathbb{C}^2) \simeq \mathbb{R}^2$. The (total) space of all η -self-adjoint operators which admit a C -symmetry is given by

$$\text{CH}_\eta(\mathbb{C}^2) := \bigsqcup_{\Xi \in C_\eta(\mathbb{C}^2)} H_{\eta,\Xi}(\mathbb{C}^2).$$

This is a real vector bundle over $C_\eta(\mathbb{C}^2) \simeq \mathbb{C}$ with typical fiber \mathbb{R}^2 . This bundle is trivial since \mathbb{C} is contractible.

The metric operator η has the property of a C -symmetry. It turns out that $H_{\eta,0}(\mathbb{C}^2) := H_{\eta,\eta}(\mathbb{C}^2)$ is the space of the self-adjoint and η -self-adjoint operators. As a consequence of Lemma (B.4) one has that:

Corollary B.5. *Any $H \in H_{\eta,0}(\mathbb{C}^2)$ can be uniquely represented as*

$$H = \begin{pmatrix} u+v & 0 \\ 0 & u-v \end{pmatrix}, \quad u, v \in \mathbb{R}. \quad (\text{B.6})$$

Let $H = H(u, v)$ the element of $H_{\eta,\Xi}(\mathbb{C}^2)$ given by the parametrization (B.5). Then the following facts are easily verifiable:

- (1) $\det H(u, v) = u^2 - v^2$ and $\text{Tr } H(u, v) = 2u$;

(2) The eigenvalues of $H(u, v)$ are $\lambda_{\pm} := u \pm v$ and the related eigenvectors are

$$\psi_+ := \begin{pmatrix} r e^{i\theta} \\ 1 - \sqrt{1+r^2} \end{pmatrix}, \quad \psi_- := \begin{pmatrix} -r e^{i\theta} \\ 1 + \sqrt{1+r^2} \end{pmatrix}$$

where r and θ are the parameters which parametrize the C -symmetry Ξ according to (B.3). Moreover $\Xi\psi_{\pm} = \pm\psi_{\pm}$.

An element $H \in \mathbf{H}_{\eta, \Xi}(\mathbb{C}^2)$ is called *gapped* (in zero) if it is invertible. The subset of the gapped element of $\mathbf{H}_{\eta, \Xi}(\mathbb{C}^2)$ is denoted with

$$\mathbf{H}_{\eta, \Xi}^G(\mathbb{C}^2) := \{H \in \mathbf{H}_{\eta, \Xi}(\mathbb{C}^2) \mid \det H \neq 0\}.$$

Since $H \in \mathbf{H}_{\eta, \Xi}^G(\mathbb{C}^2)$ implies that $\lambda_{\pm} \neq 0$ it follows that one can decompose $\mathbf{H}_{\eta, \Xi}^G(\mathbb{C}^2)$ into four disjoint components according to the signs of λ_+ and λ_- . Let introduce the subset

$$\begin{aligned} \mathbf{H}_{\eta, \Xi}^{(+,+)}(\mathbb{C}^2) &:= \{H \in \mathbf{H}_{\eta, \Xi}(\mathbb{C}^2) \mid \lambda_+ > 0, \lambda_- > 0\} \simeq \{(u, v) \in \mathbb{R}^2 \mid u > |v| \geq 0\} \\ \mathbf{H}_{\eta, \Xi}^{(+,-)}(\mathbb{C}^2) &:= \{H \in \mathbf{H}_{\eta, \Xi}(\mathbb{C}^2) \mid \lambda_+ > 0, \lambda_- < 0\} \simeq \{(u, v) \in \mathbb{R}^2 \mid v > |u| \geq 0\} \\ \mathbf{H}_{\eta, \Xi}^{(-,+)}(\mathbb{C}^2) &:= \{H \in \mathbf{H}_{\eta, \Xi}(\mathbb{C}^2) \mid \lambda_+ < 0, \lambda_- > 0\} \simeq \{(u, v) \in \mathbb{R}^2 \mid v > -|u| \leq 0\} \\ \mathbf{H}_{\eta, \Xi}^{(-,-)}(\mathbb{C}^2) &:= \{H \in \mathbf{H}_{\eta, \Xi}(\mathbb{C}^2) \mid \lambda_+ < 0, \lambda_- < 0\} \simeq \{(u, v) \in \mathbb{R}^2 \mid u < -|v| \leq 0\} \end{aligned}$$

where the last isomorphisms are obtained in terms of the parametrization (B.5) and the related description of the eigenvalues $\lambda_{\pm} = u \pm v$. It results that these four components are connected and contractible. Summing up, we proved that:

Theorem B.6. *The space $\mathbf{H}_{\eta, \Xi}^G(\mathbb{C}^2)$ is the disjoint union of four connected and contractible component*

$$\mathbf{H}_{\eta, \Xi}^G(\mathbb{C}^2) = \mathbf{H}_{\eta, \Xi}^{(+,+)}(\mathbb{C}^2) \sqcup \mathbf{H}_{\eta, \Xi}^{(+,-)}(\mathbb{C}^2) \sqcup \mathbf{H}_{\eta, \Xi}^{(-,+)}(\mathbb{C}^2) \sqcup \mathbf{H}_{\eta, \Xi}^{(-,-)}(\mathbb{C}^2).$$

Then $\mathbf{H}_{\eta, \Xi}^G(\mathbb{C}^2)$ has the topology of a space of four points. The representative of each component can be chosen as follow:

$$\mathbf{H}_{\eta, \Xi}^{(\pm, \pm)}(\mathbb{C}^2) \simeq \{\pm \mathbb{1}\}, \quad \mathbf{H}_{\eta, \Xi}^{(\pm, \mp)}(\mathbb{C}^2) \simeq \{\pm \Xi\}.$$

The two components $\mathbf{H}_{\eta, \Xi}^{(\pm, \pm)}(\mathbb{C}^2)$ are called *trivial* since describe gapped operators which can be always deformed to a multiple of the identity without violating the gap condition. As a consequence the two components $\mathbf{H}_{\eta, \Xi}^{(\pm, \mp)}(\mathbb{C}^2)$ are called *non trivial*.

Remark B.7 (Comparison with the self-adjoint case). Let H be a self-adjoint matrix on \mathbb{C}^2 . Let λ_1 and λ_2 be the eigenvalues of H and assume the *non-trivial* zero gap condition $\lambda_1 \lambda_2 < 0$. Let P_1 and P_2 be the spectral projections associated with the eigenvectors ψ_1 and ψ_2 . Consider the continuous path

$$H_t := \sum_{j=1}^2 \left[(1-t) \lambda_j + t \operatorname{sgn}(\lambda_j) \right] P_j, \quad t \in [0, 1]$$

where $\operatorname{sgn}(\lambda_j) := \frac{\lambda_j}{|\lambda_j|}$. Evidently $H_0 = H$ and $\Gamma := H_1$ verifies $\Gamma^2 = \mathbb{1}$, namely it is a gradation. With respect the basis $\{\psi_1, \psi_2\}$ the matrix Γ reads

$$\Gamma = \begin{pmatrix} \operatorname{sgn}(\lambda_1) & 0 \\ 0 & \operatorname{sgn}(\lambda_2) \end{pmatrix}. \quad (\text{B.7})$$

Consider the unitari matrix

$$T_{\theta} := \begin{pmatrix} \cos \theta & \sin \theta \\ \sin \theta & -\cos \theta \end{pmatrix}. \quad (\text{B.8})$$

A direct computation shows that $T_\theta = T_\theta^* = T_\theta^{-1}$. Moreover, one has that

$$\Gamma_\theta := T_\theta \Gamma T_\theta := \begin{pmatrix} \operatorname{sgn}(\lambda_1) \cos^2 \theta + \operatorname{sgn}(\lambda_2) \sin^2 \theta & [\operatorname{sgn}(\lambda_1) - \operatorname{sgn}(\lambda_2)] \cos \theta \sin \theta \\ [\operatorname{sgn}(\lambda_1) - \operatorname{sgn}(\lambda_2)] \cos \theta \sin \theta & \operatorname{sgn}(\lambda_1) \sin^2 \theta + \operatorname{sgn}(\lambda_2) \cos^2 \theta \end{pmatrix}.$$

If one consider the continuous path of Γ_θ with $\theta \in [0, \pi/2]$ one obtains that $\Gamma_0 = \Gamma$ while

$$\Gamma_{\frac{\pi}{2}} = \begin{pmatrix} \operatorname{sgn}(\lambda_2) & 0 \\ 0 & \operatorname{sgn}(\lambda_1) \end{pmatrix}. \quad (\text{B.9})$$

This computation shows that the *non-trivial* component of the self-adjoint matrix gapped in zero is contractible to a single point represented by the standard gradation η . On the contrary, Theorem (B.6) proves that the set $H_{\eta,0}$ has two non-trivial components $H_{\eta,0}^{(\pm, \mp)}(\mathbb{C}^2)$ represented by $\pm\eta$ respectively. In this case it is not possible to use the matrix T_θ to intertwine between η and $-\eta$ since $\eta T_\theta \neq T_\theta \eta$ and the commutation with η is a defining property for $H_{\eta,0}$. This example shows that the indefinite metric structure induced by η induces more structure with respect to the usual self-adjoint case. The difference between the self-adjoint case and the η -self-adjoint case is a manifestation of the splitting phenomenon (5.8). ◀

REFERENCES

- [AE] Alù, A.; Engheta, N.: *Pairing an Epsilon-Negative Slab With a Mu-Negative Slab: Resonance, Tunneling and Transparency*. IEEE Trans. Antennas Propag. **51**, 2558-2571 (2003)
- [AF] Ando, Y.; Fu, L.: *Topological crystalline insulators and topological superconductors: from concepts to materials*. Annu. Rev. Cond. Matt. Phys. **6**, 361-381 (2015)
- [AI] Azizov, T. Ya.; Iokhvidov, I. S.: *Linear operators in spaces with an indefinite metric*. John Wiley & Sons Inc., Chichester, 1989
- [AK1] Albeverio, S.; Kuzhel, S.: *Pseudo-Hermiticity and Theory of Singular Perturbations*. Lett. Math. Phys. **67**, 223-238 (2004)
- [AK2] Albeverio, S.; Kuzhel, S.: *\mathcal{PT} – Symmetric Operators in Quantum Mechanics : Krein Spaces Methods*. In *Non-Selfadjoint Operators in Quantum Physics: Mathematical Aspects* (Eds. F. Bagarello, J. P. Gazeau, F. H. Szafraniec & M. Znojil), pg. 293-343. Wiley & Sons, Inc. 2015
- [ASe] Atiyah, M. F.; Segal, G.: *Twisted K-theory*. Ukr. Mat. Visn. **1**, 287-330 (2004)
- [ASi] Atiyah, M. F.; Singer, I. M.: *Index theory for skew-adjoint Fredholm operators*. Inst. Hautes Études Sci. Publ. Math. **37**, 5-26 (1969)
- [Ati1] Atiyah, M. F.: *K-theory and reality*. Quart. J. Math. Oxford Ser. (2) **17**, 367-386 (1966)
- [Ati2] Atiyah, M. F.: *K-theory*. W. A. Benjamin, New York, 1967
- [Ati3] Atiyah, M. F.: *Algebraic topology and operators in Hilbert space*. In *Lectures in Modern Analysis and Applications. I* **103**, pg 101-121, Springer, Berlin, 1969
- [Atk] Atkinson, F. V.: *The normal solvability of linear equations in normed spaces*. Mat. Sb. **28**, 3-14 (1951)
- [AZ] Altland, A.; Zirnbauer, M.: *Non-standard symmetry classes in mesoscopic normal-superconducting hybrid structures*. Phys. Rev.B **55**, 1142-1161 (1997)
- [BB] Bender, C. M.; Boettcher S.: *Real Spectra in Non-Hermitian Hamiltonians Having \mathcal{PT} Symmetry*. Phys. Rev. Lett. **80**, 5243 (1998)
- [BBJ1] Bender, C. M.; Brody, D. C.; Jones, H. F.: *Complex Extension of Quantum Mechanics*. Phys. Rev. Lett. **89**, 270401 (2002)
- [BBJ2] Bender, C. M.; Brody, D. C.; Jones, H. F.: *Must a Hamiltonian be Hermitian?*. Am. J. Phys. **71**, 1905-1102 (2003)
- [Ben1] Bender C. M.: *Making sense of non-Hermitian Hamiltonians*. Rep. Prog. Phys. **70**, 947-1018 (2007)
- [Ben2] Bender C. M.: *\mathcal{PT} -symmetric quantum theory*. J. Phys.: Conf. Ser. **631**, 012002 (2015)
- [Ben3] Bender C. M.: *\mathcal{PT} symmetry in quantum physics: From a mathematical curiosity to optical experiments*. Euro. News **47**, 17-20 (2016)
- [BMS] Bracci, L.; Morchio, G.; Strocchi, F.: *Wigner's theorem on symmetries in indefinite metric spaces*. Commun. Math. Phys. **41**, 289-299 (1975)
- [Bog] Bognár, J.: *Indefinite Inner Product Spaces*. Springer-Verlag, Berlin, 1974
- [Bro] Brown, K. S.: *Cohomology of Groups*. Springer-Verlag, Berlin, 1982
- [BS] Birman, M. Sh.; Solomyak, M. Z.: *L_2 -Theory of the Maxwell operator in arbitrary domains*. Russ. Math. Surv. **42**, 75-96 (1987)

- [Buc] Buckholtz, D.: *Hilbert Space Idempotents and Involutions*. Proc. Amer. Math. Soc. **128**, 1415-1418 (2000)
- [CGBV] Chtchelkatchev, N. M.; Golubov, A. A.; Baturina, T. I.; Vinokur, V. M.: *Stimulation of the Fluctuation Superconductivity by \mathcal{PT} Symmetry*. Phys. Rev. Lett. **109**, 1504058 (2012)
- [CN] Ćurgus, B.; Najman, B.: *The Operator $(\operatorname{sgn} x) d^2/dx^2$ is Similar to a Selfadjoint Operator in $L^2(\mathbb{R})$* . Proc. Amer. Math. Soc. **123**, 1125-1128 (1995)
- [Cro] Croke, S.: *\mathcal{PT} -symmetric Hamiltonians and their application in quantum information*. Phys. Rev. A **91**, 052113 (2015)
- [CTSR] Chiu, C.-K.; Teo, J. C. Y.; Schnyder, A. P.; Ryu, S.: *Classification of topological quantum matter with symmetries*. Rev. Mod. Phys. **88**, 035005 (2016)
- [DD] Dixmier, J.; Douady, A.: *Champs continus d'espaces hilbertiens et de C^* -algèbres*. Bull. Soc. Math. France **91**, 227-284 (1963)
- [DDT] Dorey, P.; Dunning, C.; Tateo, R.: *Spectral equivalences, Bethe Ansatz equations, and reality properties in \mathcal{PT} -symmetric quantum mechanics*. J. Phys. A **32**, 5679-5703 (2001)
- [DG1] De Nittis, G.; Gomi, K.: *Classification of "Real" Bloch-bundles: Topological Quantum Systems of type AI*. J. Geom. Phys. **86**, 303-338 (2014)
- [DG2] De Nittis, G.; Gomi, K.: *Differential geometric invariants for time-reversal symmetric Bloch-bundles: the "Real" case*. J. Math. Phys. **57**, 053506 (2016)
- [Dir] Dirac P. M. A.: *The physical interpretation of quantum mechanics*. Proc. R. Soc. London A **180**, 1-40 (1942)
- [DK] Donovan, D.; Karoubi, M.: *Graded Brauer groups and K-theory with local coefficients*. Inst. Hautes Études Sci. Publ. Math. **38**, 5-25 (1970)
- [DL1] De Nittis, G.; Lein, M.: *The Perturbed Maxwell Operator as Pseudodifferential Operator*. Documenta Math. **19**, 63-101 (2014)
- [DL2] De Nittis, G.; Lein, M.: *Effective Light Dynamics in Perturbed Photonic Crystals*. Commun. Math. Phys. **332**, 221-260 (2014)
- [DL3] De Nittis, G.; Lein, M.: *The Schrödinger Formalism of Electromagnetism and Other Classical Waves: How to Make Quantum-Wave Analogies Rigorous*. Accepted for publication in Ann. Phys., (2018)
- [Dup] Dupré M. J.: *Classifying Hilbert Bundles*. J. Funct. Anal. **15**, 244-278 (1974)
- [Du] Dupont, J. L.: *Symplectic Bundles and KR-Theory*. Math. Scand. **24**, 27-30 (1969)
- [EMCM] El-Ganainy, R.; Makris, K.; Christodoulides, D. N.; Musslimani, Z.: *Theory of coupled optical \mathcal{PT} -symmetric structures*. Opt. Lett. **32**, 2632-2634 (2007)
- [ESHK] Esaki, K.; Sato, M.; Hasebe, K.; Kohmoto, M.: *Edge states and topological phases in non-Hermitian systems*. Phys. Rev. B **84**, 205128 (2011)
- [EU] Espinoza, J.; Uribe, B.: *Topological properties of the unitary group*. JP Journal of Geometry and Topology **16**, 45-55 (2014)
- [FHT] Freed, D. S.; Hopkins, M. J.; Teleman, C.: *Loop groups and twisted K-theory. I*. J. Topol. **4**, 737-798 (2011)
- [FM] Freed, D. S.; Moore, G. W.: *Twisted Equivariant Matter*. Ann. Henri Poincaré **14**, 1927-2023 (2013)
- [FR] Fredkin, D. R.; Ron, A.: *Effectively left-handed (negative index) composite material*. Appl. Phys. Lett. **81**, 1753-1755 (2002)
- [Ghe] Gheondea, A.: *Canonical Forms of Unbounded Unitary Operators in Krein Spaces*. Publ. RIMS, Kyoto Univ. **24**, 205-224 (1988)
- [Gho] Ghosh, P. K.: *A note on the topological insulator phase in non-Hermitian quantum systems*. J. Phys.: Condens. Matter, **24**, 145302 (2012)
- [GI] Ginzburg, Yu. P.; Iokhvidov, I. S.: *The geometry of infinite-dimensional spaces with a bilinear metric*. Russian Math. Surveys **17**, 1-51 (1962)
- [Gom] Gomi, K.: *Freed-Moore K-theory*. E-print arXiv:1705.09134 (2017)
- [GV] Galda, A.; Vinokur, V. M.: *Parity-time symmetry breaking in magnetic systems*. Phys. Rev. B **94**, 020408 (2016)
- [HH] Hu, Y. C.; Hughes, T. L.: *Absence of topological insulator phases in non-Hermitian \mathcal{PT} -symmetric Hamiltonians*. Phys. Rev. B **84**, 153101 (2011)
- [HK2] Hasan, M. Z.; Kane, C. L.: *Colloquium: Topological insulators*. Rev. Mod. Phys. **82**, 3045-3067 (2010)
- [Hus] Husemoller, D.: *Fibre bundles*. Springer-Verlag, New York, 1994
- [IK1] Iokhvidov, I. S.; Krein, M. G.: *Spectral theory of operators in the spaces with indefinite metric I*. Trudy Mosk. Mat. Obshch. **5**, 367-432 (1956)
- [IK2] Iokhvidov, I. S.; Krein, M. G.: *Spectral theory of operators in the spaces with indefinite metric II*. Trudy Mosk. Mat. Obshch. **8**, 413-496 (1959)
- [JCHZ3] Jiang, H. T.; Chen, H.; Li, H. Q.; Zhang, Y. W.; Zi, J.; Zhu, S. Y.: *Properties of one-dimensional photonic crystals containing single-negative materials*. Phys. Rev. E **69**, 066607 (2004)

- [Jon] Jones, H. F.: *On pseudo-Hermitian Hamiltonians and their Hermitian counterparts*. J. Phys. A **38**, 1741-1746 (2005)
- [Kar] Karoubi, M.: *K-theory. An introduction*. Springer-Verlag, Berlin-New York, 1978.
- [Kat] Kato, T.: *Scattering Theory with Two Hilbert Spaces*. J. Funct. Anal **1**, 342-369 (1967)
- [Kit] Kitaev, A.: *Periodic table for topological insulators and superconductors*. AIP Conf. Proc. **1134**, 22-30 (2009)
- [KL] Krein, M. G.; Langer, H.: *Spectral function of a self-adjoint operator in a space with indefinite metric*. Dokl. Akad. Nauk SSSR **152**, 15-92 (1963)
- [KR] Kadison, R. V.; Ringrose, J. R.: *Fundamentals of the Theory of Operator Algebras I*. Graduate Studies in Mathematics, vol. 15. Academic Press, Inc., Boston (1983)
- [Kre] Krein, M. G.: *Introduction to the geometry of indefinite J -spaces and the theory of operators in these spaces*. Second Math. Summer School, Part I. Naukoav Dumka Kiev, 15-92 (1965)
- [KSh] Kissin, E.; Shulman, V.: *Representations on Krein spaces and derivations of C^* -algebras*. Addison Wesley Longman, London (1997)
- [KSu] Kuzhel, S.; Sudilovskaya, V.: *Towards theory of C -symmetries*. Opuscula Math. **37**, 65-80 (2017)
- [Ke] Kellendonk, J.: *On the C^* -algebraic approach to topological phases for insulators*. Ann. Henri Poincaré **18**, 2251-2300 (2017)
- [Kub] Kubota, Y.: *Controlled Topological Phases and Bulk-edge Correspondence* Commun. Math. Phys. **349**, 493-525 (2017)
- [Kuc] Kuchment, P.: *Floquet theory for partial differential equations*. Birkhäuser, Boston, 1993
- [Kui] Kuiper, N. H.: *The homotopy type of the unitary group of Hilbert space* Topology **3**, 19-30 (1965)
- [Kuz] Kuzhel, S.: *On Pseudo-Hermitian Operators with Generalized C -symmetries*. Operator Theory: Advances and Applications **190**, 375-385 (2009)
- [KWLWL] Ke, S.; Wang, B.; Long, H.; Wang, K.; Lu, P.: *Topological edge modes in non-Hermitian plasmonic waveguide arrays*. Opt. Express **25**, 11132-11143 (2017)
- [KYZ] Konotop, V. V.; Yang, J.; Zezyulin, D. A.: *Nonlinear waves in \mathcal{PT} -symmetric systems*. Rev. Mod. Phys. **88**, 035002 (2016)
- [Lan1] Langer, H.: *Zur Spektraltheorie J -selbstadjungierter Operatoren*. Math. Ann. **146**, 60-85 (1962)
- [Lan2] Langer, H.: *Spectral functions of definitizable operators in Krein spaces*. Lecture Notes in Mathematics **948**, 1-46 (1982)
- [Le] Lee, H. C.: *On Clifford Algebras and Their Representations*. Ann. of Math. **4**, 760-773 (1948)
- [Lie] Lieu, S.: *Topological phases in the non-Hermitian Su-Schrieffer-Heeger model*. Phys. Rev. B **97**, 045106 (2018)
- [LBHCN] Leykam, D.; Bliokh, K. Y.; Huang, C.; Chong, Y. D.; Nori, F.: *Edge Modes, Degeneracies, and Topological Numbers in Non-Hermitian Systems*. Phys. Rev. Lett. **118**, 040401 (2017)
- [LKS] Lee, J. M.; Kottos, T.; Shapiro, B.: *Macroscopic magnetic structures with balanced gain and loss*. Phys. Rev. B **91**, 094416 (2015)
- [Mos1] Mostafazadeh, A.: *Pseudo-Hermiticity versus \mathcal{PT} -symmetry: The necessary condition for the reality of the spectrum of a non-Hermitian Hamiltonian*. J. Math. Phys. **43**, 205-214 (2002)
- [Mos2] Mostafazadeh, A.: *Exact \mathcal{PT} -symmetry is equivalent to Hermiticity*. J. Phys. A **36**, 7081-7091 (2003)
- [Mos3] Mostafazadeh, A.: *Is Pseudo-Hermitian Quantum Mechanics an Indefinite-Metric Quantum Theory?*. A. Czech J. Phys. **53**, 1079-1084 (2003)
- [Mos4] Mostafazadeh, A.: *Krein-space formulation of \mathcal{PT} symmetry, $C\mathcal{PT}$ -inner products, and pseudo-Hermiticity*. A. Czech J. Phys. **56**, 919-933 (2006)
- [Mos5] Mostafazadeh, A.: *Pseudo-Hermitian Representation of Quantum Mechanics*. Int. J. Geom. Meth. Mod. Phys. **7**, 1191-1306 (2010)
- [Mos6] Mostafazadeh, A.: *Point interactions, metamaterials, and \mathcal{PT} -symmetry*. Ann. Phys. **368**, 56-69 (2016)
- [Nee] Neeb, K.-H.: *A note on a theorem of S. Banach*. J. Lie Theory **8**, 293-300 (1997)
- [Par] Parthasarathy, K. R.: *Projective unitary antiunitary representations of locally compact groups*. Commun. Math. Phys. **15**, 305-328 (1969)
- [Pau] Pauli, W.: *On Dirac's New Method of Field Quantization*. Rev. Mod. Phys. **15**, 175-207 (1943)
- [Pon] Pontryagin, L. S.: *Hermitian operators in spaces with indefinite metric*. Izv. Akad. Nauk. SSSR Ser. Mat. **8**, 243-280 (1944)
- [Pru] Prugovečki, E.: *Principles of Quantum General Relativity*. World Scientific Publishing, 1995
- [PS] Peano, V.; Schulz-Baldes, H.: *Topological edge states for disordered bosonic systems*. E-print arXiv:1708.08798 (2017)
- [QHZ] Qi, X.-L.; Hughes, T.; Zhang, S.-C.: *Topological field theory of time-reversal invariant insulators*. Phys. Rev. B **78**, 195424 (2008)

- [Ros] Rosenberg, J.: *Continuous-trace algebras from the bundle theoretic point of view*. J. Austral. Math. Soc. Ser. A **47**, 368-381 (1989)
- [RSM] Rubinstein, J.; Sternberg, P.; Ma, Q.: *Bifurcation Diagram and Pattern Formation of Phase Slip Centers in Superconducting Wires Driven with Electric Currents*. Phys. Rev. Lett. **99**, 167003 (2007)
- [Sch1] Schomerus, H.: *Topologically protected midgap states in complex photonic lattices*. Opt. Lett. **38**, 1912-1914 (2013)
- [Sch2] Schottenloher, M.: *The Unitary Group In Its Strong Topology*. E-print arXiv:1309.5891 (2013)
- [SGH] Scholtz, F. G.; Geyer, H. B.; Hahne, F. J. W.: *Quasi-Hermitian operators in quantum mechanics and the variational principle*. Ann. Phys. **213**, 74-101 (1992)
- [Shm] Shmul'yan, Yu.L.: *Extension theory operators and spaces with an indefinite metric*. Izv. Akad. Nauk, SSSR Ser. Mat. **38**, 895-907 (1974)
- [SHEK] Sato, M.; Hasebe, K.; Esaki, K.; Kohmoto, M.: *Time-Reversal Symmetry in Non-Hermitian Systems*. Prog. Theor. Phys. **127**, 937-974 (2012)
- [Sim] Simms, D. J.: *Topological aspects of the projective unitary group*. Proc. Camb. Phil. Soc. **68**, 57-60 (1970)
- [SLZEK] Schindler, J.; Li, A.; Zheng, M. C.; Ellis, F. M.; Kottos, T.: *Experimental study of active LRC circuits with \mathcal{PT} symmetries*. Phys. Rev. A **84**, 040101 (2011)
- [SRFL] Schnyder, A. P.; Ryu, S.; Furusaki, A.; Ludwig, A. W. W.: *Classification of topological insulators and superconductors in three spatial dimensions*. Phys. Rev. B **78**, 195125 (2008)
- [Tak] Takesaki, M.: *Theory of Operator Algebras I*. Springer-Verlag, 2002
- [Thi] Thiang, G. C.: *On the K-theoretic classification of topological phases of matter*. Ann. Henri Poincaré **17**, 757-794 (2016)
- [Tom] Tomita, M.: *Operators and operator algebras in Krein spaces, I Spectral Analysis in Pontrjagin spaces*. RIMS-Kōkyūroku **398**, 131-158 (1980)
- [TW] Tuynman, G. M.; Wiegerinck, W. A. J. J.: *Central extensions and physics*. J. Geom. Phys. **4**, 207-258 (1987)
- [Upm] Upmeyer, H.: *Symmetric Banach Manifolds and Jordan C^* -Algebras*. North-Holland Math. Stud. **104** Notas de Matemática 96, North-Holland, Amsterdam, 1985.
- [Wig] Wigner, E.P.: *Group theory and its application to the quantum, mechanics of atomic spectra* (translation from German by J. J. Griffin). Pure and Applied Physics, vol. 5, Academic Press, New York, 1959
- [Wil] Wilcox, C. H.: *Wave Operators and Asymptotic Solutions of Wave Propagation Problems of Classical Physics*. Arch. Ration. Mech. Anal. **22**, 37-78 (1966)
- [Yuc] Yuce, C.: *Topological phase in a non-Hermitian \mathcal{PT} -symmetric system*. Phys. Lett. A **379**, 1213-1218 (2015)
- [ZLC] Zhu, B.; Lu, R.; Chen, S.: *\mathcal{PT} -symmetry in the non-Hermitian Su-Schrieffer-Heeger model with complex boundary potentials*. Phys. Rev. A **89**, 062102 (2014)
- [ZSW] Zhao, K. F.; Schaden, M.; Wu, Z.: *Enhanced magnetic resonance signal of spin-polarized Rb atoms near surfaces of coated cells*. Phys. Rev. A **81**, 042903 (2010)

(De Nittis) FACULTAD DE MATEMÁTICAS & INSTITUTO DE FÍSICA, PONTIFICIA UNIVERSIDAD CATÓLICA, SANTIAGO, CHILE
E-mail address: gidenittis@mat.uc.cl

(Gomi) DEPARTMENT OF MATHEMATICAL SCIENCES, SHINSHU UNIVERSITY, NAGANO, JAPAN
E-mail address: kgomi@math.shinshu-u.ac.jp