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LOCAL EIGENVALUE ASYMPTOTICS OF THE PERTURBED KREIN LAPLACIAN

Vincent Bruneau¹, Georgi Raikov²

¹ UMR CNRS 5251, Institut de Mathématiques de Bordeaux, Université de Bordeaux, Talence, France

² Facultad de Matemáticas, Pontificia Universidad Católica de Chile, Chile

Local Eigenvalue Asymptotics of the Perturbed Krein Laplacian

Vincent Bruneau and Georgi Raikov

ABSTRACT. We consider the Krein Laplacian on a regular bounded domain $\Omega\subset\mathbb{R}^d,\ d\geq 2$, perturbed by a real-valued multiplier V vanishing on the boundary. Assuming that V has a definite sign, we investigate the asymptotics of the functions counting the eigenvalues of K+V which converge to the origin from below or from above. We show that the effective Hamiltonian that governs the main asymptotic term of these functions is the harmonic Toeplitz operator T_V with symbol V, unitarily equivalent to a pseudodifferential operator on the boundary. In the cases where V admits a power-like decay at $\partial\Omega$, or V is compactly supported in Ω , and Ω and supp V are radially symmetric, we obtain the main asymptotic term of the eigenvalue counting functions.

1. Introduction

In this article we study the spectral properties of the perturbed Krein Laplacian K+V in a bounded domain $\Omega\subset\mathbb{R}^d,\,d\geq 2$, with boundary $\partial\Omega\in C^\infty$. The Krein Laplacian $K:=-\Delta$ is defined on sufficiently regular functions $u:\Omega\to\mathbb{C}$ which satisfy the boundary condition

$$\frac{\partial u}{\partial \nu} = \mathcal{D}u \quad \text{on} \quad \partial \Omega,$$

where ν is the unit outer normal vector at $\partial\Omega$, and \mathcal{D} is the Dirichlet-to-Neumann operator, a first-order elliptic pseudodifferential operator (Ψ DO), self-adjoint in $L^2(\partial\Omega)$. Then $K\geq 0$ is self-adjoint in $L^2(\Omega)$, and one of its remarkable properties is that its essential spectrum is not empty. Namely, $\sigma_{\rm ess}(K)=\{0\}$, and the zero is an isolated eigenvalue of K of infinite multiplicity. Further, we assume that the perturbation of K is the multiplier by the function $V\in C(\overline{\Omega};\mathbb{R})$. Then, evidently the operator K+V, on the domain of K, is self-adjoint in $L^2(\Omega)$. Moreover,

$$\sigma_{\rm ess}(K+V) = V(\partial\Omega)$$

(see Theorem 2.1 below). Assuming that V vanishes identically on $\partial\Omega$, we get

$$\sigma_{\rm ess}(K+V) = \sigma_{\rm ess}(K) = \{0\}.$$

However, in contrast to the unperturbed operator K, the zero is an accumulation point of the discrete eigenvalues of the perturbed operator K + V. We suppose

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in addition that V has a definite (negative or positive) sign and investigate the asymptotic distribution of the eigenvalues of K+V adjoining the origin. First, in Theorem 2.2 we show that the effective Hamiltonian governing the eigenvalue counting functions for K+V is the Toeplitz operator $T_V:=PV_{|Ker\,K}$, where P is the orthogonal projection onto Ker K. That is why, in Section 3 we discuss the general spectral properties of T_V . Further, in Section 4 we assume that V admits a power-like decay at $\partial\Omega$, and examine the eigenvalue asymptotics for the compact operator T_V , unitarily equivalent to a classical Ψ DO on the boundary. We obtain the main asymptotic term and a sharp remainder estimate of the eigenvalue counting function for T_V . Finally, in Section 5, we analyze the case where V is compactly supported in Ω . More precisely, we suppose that Ω is the unit ball in \mathbb{R}^d while supp V is the concentric ball of radius $c \in (0, \infty)$, and obtain the main asymptotic term of the eigenvalue counting function for T_V .

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A more detailed exposition of some of the results of this paper can be found in the preprint [13].

2. The Krein Laplacian and its perturbations

Let $\Omega \subset \mathbb{R}^d$, $d \geq 2$, be a bounded domain with smooth boundary $\partial \Omega$. For $s \in \mathbb{R}$ denote by $H^s(\Omega)$ and $H^s(\partial \Omega)$ the Sobolev spaces on Ω and $\partial \Omega$ respectively. Moreover, as usual, we denote by $H^s_0(\Omega)$, s > 1/2, the closure of $C_0^\infty(\Omega)$ in $H^s(\Omega)$. Set also $H^2_D(\Omega) := H^2(\Omega) \cap H^1_0(\Omega)$. Define the minimal Laplacian

$$\Delta_{\min} := \Delta$$
, $\operatorname{Dom} \Delta_{\min} = H_0^2(\Omega)$.

The operator Δ_{\min} is symmetric but not self-adjoint in $L^2(\Omega)$, since we have

(2.1)
$$\Delta_{\min}^* =: \Delta_{\max} = \Delta, \quad \operatorname{Dom} \Delta_{\max} = \left\{ u \in L^2(\Omega) \mid \Delta u \in L^2(\Omega) \right\},$$

 Δu being the distributional Laplacian of $u \in L^2(\Omega)$. Note that

(2.2)
$$\operatorname{Ker} \Delta_{\max} = \mathcal{H}(\Omega) := \left\{ u \in L^2(\Omega) \, | \, \Delta u = 0 \text{ in } \Omega \right\}.$$

It is well known that $\mathcal{H}(\Omega)$ is a closed subspace of $L^2(\Omega)$ (see e.g. [23]). The Laplace equation $\Delta u = 0$ in (2.2) is understood a priori in the distributional sense. However, by the Weyl lemma, if u belongs to $\mathcal{D}'(\Omega)$, the class of distributions over $C_0^{\infty}(\Omega)$, and $\Delta u = 0$, then $u \in C^{\infty}(\Omega)$ (see the original work [30] for a proof in the case $u \in L^1_{\text{loc}}(\Omega)$, and the monograph [17] whose Chapter 10 contains an extension to general $u \in \mathcal{D}'(\Omega)$).

Lemma 2.1. The domain Dom Δ_{max} admits the direct-sum decomposition

(2.3)
$$\operatorname{Dom} \Delta_{\max} = \mathcal{H}(\Omega) + H_D^2(\Omega).$$

PROOF. Let us first show that the sum at the r.h.s. of (2.3) is direct. Assume that $u_1 \in \mathcal{H}(\Omega)$, $u_2 \in H_D^2(\Omega)$, and $u_1 + u_2 = 0$. Then u_2 satisfies the homogeneous

boundary-value problem

$$\begin{cases} \Delta u_2 = 0 & \text{in } \Omega, \\ u_2 = 0 & \text{on } \partial \Omega. \end{cases}$$

Hence, $u_2 = 0$, and $u_1 = 0$. Evidently, if $u_1 \in \mathcal{H}(\Omega)$, $u_2 \in H_D^2(\Omega)$, then $u_1 + u_2 \in \text{Dom } \Delta_{\text{max}}$. Pick now $u \in \text{Dom } \Delta_{\text{max}}$, and define the Dirichlet Laplacian

$$\Delta_D := \Delta$$
, $\operatorname{Dom} \Delta_D := H_D^2(\Omega)$.

Then u_1 and u_2 defined by $u_2 := \Delta_D^{-1} \Delta u$, $u_1 := u - u_2$ clearly satisfy

$$u_1 \in \mathcal{H}(\Omega), \quad u_2 \in H_D^2(\Omega), \quad u = u_1 + u_2.$$

Introduce the Krein Laplacian

$$K := -\Delta$$
, $\operatorname{Dom} K = \mathcal{H}(\Omega) \dotplus H_0^2(\Omega)$.

The operator $K \geq 0$, self-adjoint in $L^2(\Omega)$, is the von Neumann - Krein "soft" extension of $-\Delta_{\min}$, remarkable for the fact that any other self-adjoint extension $S \geq 0$ of $-\Delta_{\min}$ satisfies

$$(S+I)^{-1} \le (K+I)^{-1}$$

(see [29, 24]). Evidently, Ker $K = \mathcal{H}(\Omega)$. The domain Dom K admits a more explicit description in the terms of the Dirichlet-to-Neumann operator \mathcal{D} . For $f \in C^{\infty}(\partial\Omega)$, $\mathcal{D}f$ is defined by

$$\mathcal{D}f = \frac{\partial u}{\partial \nu|_{\partial\Omega}},$$

where u is the solution of the boundary-value problem

$$\left\{ \begin{array}{ll} \Delta u = 0 & \text{in } \Omega, \\ u = f & \text{on } \partial \Omega. \end{array} \right.$$

The operator \mathcal{D} is a first-order elliptic operator; by the elliptic regularity, it extends to a bounded operator form $H^s(\partial\Omega)$ into $H^{s-1}(\partial\Omega)$, $s \in \mathbb{R}$. Then we have

$$\operatorname{Dom} K = \left\{ u \in \operatorname{Dom} \Delta_{\max} \; \left| \; \frac{\partial u}{\partial \nu}_{|\partial \Omega} = \mathcal{D} \left(u_{|\partial \Omega} \right) \; \right\} \right.$$

(see [18, Theorem III.1.2]).

Denote by L the restriction of K onto Dom $K \cap \mathcal{H}(\Omega)^{\perp}$ where $\mathcal{H}(\Omega)^{\perp} := L^2(\Omega) \ominus \mathcal{H}(\Omega)$. Then, L is self-adjoint in the Hilbert space $\mathcal{H}(\Omega)^{\perp}$.

PROPOSITION 2.1. ([24], [4, Theorem 5.1]) The spectrum of L is purely discrete and positive, and, hence, L^{-1} is a compact operator on $\mathcal{H}(\Omega)^{\perp}$. As a consequence, $\sigma_{\rm ess}(K) = \{0\}$, and the zero is an isolated eigenvalue of K of infinite multiplicity.

The Krein Laplacian K arises naturally in the so called abstract buckling problem

$$\left\{ \begin{array}{l} \Delta^2 u = -\lambda \Delta u, \\ u_{|\partial\Omega} = \frac{\partial u}{\partial \nu}_{|\partial\Omega} = 0, \\ u \in \operatorname{Dom} \Delta_{\max}. \end{array} \right.$$

Then the admissible values of $\lambda \neq 0$ for which the buckling problem has a non-trivial solution, coincide with the non-zero eigenvalues of K (see e.g. [19, 5]).

Let $V \in C(\overline{\Omega}; \mathbb{R})$. Then the operator K + V with domain Dom K is self-adjoint in $L^2(\Omega)$. In this article, we investigate the spectral properties of K + V.

Remark: The perturbations K_V of the Krein Laplacian K discussed in [6] are of different nature than the perturbations K+V considered here. Namely, the authors of [6] assume that $V \geq 0$, define the maximal operator $K_{V,\max}$ as

$$K_{V,\max} := -\Delta + V$$
, Dom $K_{V,\max} := \text{Dom } \Delta_{\max}$,

and set

$$K_V := -\Delta + V$$
, $\operatorname{Dom} K_V := \operatorname{Ker} K_{V,\max} + H_0^2(\Omega)$.

Thus, if $V \neq 0$, then the operators K_V and $K_0 = K$ are self-adjoint on different domains, while the operators K + V introduced here are self-adjoint on the same domain Dom K. It is shown in [6] that for any $0 \leq V \in L^{\infty}(\Omega)$ we have $K_V \geq 0$, $\sigma_{\rm ess}(K_V) = \{0\}$, and the zero is an isolated eigenvalue of K_V of infinite multiplicity. As we will see in what follows, the spectral properties of K + V could be quite different.

In Theorem 2.1 below we locate the essential spectrum of the operator K+V. For its proof we need some additional notations and definitions. Let $P: L^2(\Omega) \to L^2(\Omega)$ be the orthogonal projection onto $\mathcal{H}(\Omega)$. Assume that $V \in C(\overline{\Omega})$, and introduce the harmonic Toeplitz operator $T_V := PV: \mathcal{H}(\Omega) \to \mathcal{H}(\Omega)$. Evidently, T_V is bounded, and if V is real-valued, then T_V is self-adjoint. Note that T_V could be well defined as a bounded and even compact operator for a much wider class of symbols V which are locally integrable in Ω and satisfy certain regularity properties near $\partial\Omega$.

We start our analysis with the location of the essential spectrum $\sigma_{\rm ess}(T_V)$ of T_V , and a criterion for the compactness of T_V in the case $V \in C(\overline{\Omega})$.

PROPOSITION 2.2. Let $\Omega \subset \mathbb{R}^d$, $d \geq 2$, be a bounded domain with boundary $\partial \Omega \in C^{\infty}$. Let $V \in C(\overline{\Omega})$.

- (i) [14, Theorem 4.5] We have $\sigma_{\text{ess}}(T_V) = V(\partial \Omega)$.
- (ii) [14, Corollary 4.7] The operator T_V is compact in $\mathcal{H}(\Omega)$ if and only if V = 0 on $\partial\Omega$.

Now we are in position to locate $\sigma_{ess}(K+V)$.

THEOREM 2.1. Let $V \in C(\overline{\Omega}; \mathbb{R})$. Then we have

(2.4)
$$\sigma_{\rm ess}(K+V) = V(\partial\Omega).$$

In particular, $\sigma_{\text{ess}}(K+V) = \{0\}$ if and only if $V_{|\partial\Omega} = 0$.

PROOF. First, we will show that the operator

$$(2.5) (K+V-i)^{-1} - (K+PVP-i)^{-1}$$

is compact. Set Q := I - P. Then

$$(K + V - i)^{-1} - (K + PVP - i)^{-1} =$$

(2.6)
$$-(K+V-i)^{-1}(K-i)(K-i)^{-1}(QVQ+PVQ+QVP)(K-i)^{-1}(K-i)(K+PVP-i)^{-1}$$
. Evidently, the operators $(K+V-i)^{-1}(K-i)$, $(K-i)(K+PVP-i)^{-1}$, P , and V , are bounded in $L^2(\Omega)$. Moreover, using the orthogonal decomposition $L^2(\Omega) = \frac{1}{2} \frac{$

 $\mathcal{H}(\Omega) \oplus \mathcal{H}(\Omega)^{\perp}$, and bearing in mind Proposition 2.1, we find that the operators $Q(K-i)^{-1}$ and $(K-i)^{-1}Q$ are compact in $L^2(\Omega)$. Now the compactness of the operator defined in (2.5) follows from (2.6). Therefore,

(2.7)
$$\sigma_{\rm ess}(K+V) = \sigma_{\rm ess}(K+PVP).$$

Further, we have $K + PVP = T_V \oplus L$ in $L^2(\Omega) = \mathcal{H}(\Omega) \oplus \mathcal{H}(\Omega)^{\perp}$, and, hence,

(2.8)
$$\sigma_{\rm ess}(K + PVP) = \sigma_{\rm ess}(T_V) \cup \sigma_{\rm ess}(L).$$

By Proposition 2.2 (i), we have $\sigma_{\rm ess}(T_V) = V(\partial\Omega)$, and by Proposition 2.1, $\sigma_{\rm ess}(L) = \emptyset$. Thus, (2.7) and (2.8) imply (2.4).

In the rest of the section we assume that $0 \leq V \in C(\overline{\Omega})$ with $V_{|\partial\Omega} = 0$, and investigate the asymptotic distribution of the discrete spectrum of the operators $K \pm V$, adjoining the origin. In particular, we show that the harmonic Toeplitz operator T_V is the effective Hamiltonian governing the main asymptotic term of the corresponding eigenvalue counting functions (see (2.11) - (2.12) below).

Let $\lambda_0 := \inf \sigma(L)$. By Proposition 2.1, we have $\lambda_0 > 0$. Introduce the eigenvalue counting functions

$$\mathcal{N}_{-}(\lambda) := \operatorname{Tr} \mathbb{1}_{(-\infty, -\lambda)}(K - V), \quad \lambda > 0,$$

$$\mathcal{N}_{+}(\lambda) := \operatorname{Tr} \mathbb{1}_{(\lambda, \lambda_0)}(K + V), \quad \lambda \in (0, \lambda_0).$$

Here and in the sequel $\mathbb{1}_S$ denotes the characteristic function of the set S; thus $\mathbb{1}_{\mathcal{I}}(T)$ is the spectral projection of the operator $T=T^*$ corresponding to the interval $\mathcal{I}\subset\mathbb{R}$. Let $T=T^*$ be a compact operator in a Hilbert space, and s>0. Set

(2.9)
$$n_{\pm}(s;T) := \operatorname{Tr} \mathbf{1}_{(s,\infty)}(\pm T).$$

Thus, $n_{+}(s;T)$ (resp., $n_{-}(s;T)$) is just the number of the eigenvalues of T larger than s (resp., smaller than -s), counted with their multiplicities.

If $T_j = T_j^*$, j = 1, 2, are two compact operators, then the Weyl inequalities

$$(2.10) n_{\pm}(s_1 + s_2; T_1 + T_2) \le n_{\pm}(s_1; T_1) + n_{\pm}(s_2; T_2)$$

hold for $s_i > 0$, j = 1, 2, (see e.g. [10, Theorem 9, Section 9.2]).

Theorem 2.2. Assume that $0 \leq V \in C(\overline{\Omega})$ and $V_{|\partial\Omega} = 0$. Then for any $\varepsilon \in (0,1)$ we have

$$(2.11) n_{+}(\lambda; T_{V}) \leq \mathcal{N}_{-}(\lambda) \leq n_{+}((1-\varepsilon)\lambda; T_{V}) + O(1),$$

$$(2.12) n_+((1+\varepsilon)\lambda; T_V) + O(1) \le \mathcal{N}_+(\lambda) \le n_+((1-\varepsilon)\lambda; T_V) + O(1),$$

as $\lambda \downarrow 0$.

PROOF. By the Birman-Schwinger principle [7, Lemma 1.1], we have 2.13)

$$\mathcal{N}_{-}(\lambda) = n_{+}(1; (K+\lambda)^{-1/2}V(K+\lambda)^{-1/2}) = n_{+}(1; V^{1/2}(K+\lambda)^{-1}V^{1/2}), \ \lambda > 0.$$

It follows from the mini-max principle that

$$n_{+}(1;(K+\lambda)^{-1/2}V(K+\lambda)^{-1/2}) \ge$$

$$n_{+}(1;P(K+\lambda)^{-1/2}V(K+\lambda)^{-1/2}P) = n_{+}(\lambda;PVP) = n_{+}(\lambda;T_{V}),$$

which, combined with the first equality in (2.13), implies the lower bound in (2.11). Further, by the Weyl inequalities (2.10), we have

$$n_{+}(1; V^{1/2}(K+\lambda)^{-1}V^{1/2}) \le$$

$$(2.14) \ n_+((1-\varepsilon)\lambda; V^{1/2}PV^{1/2}) + n_+(\varepsilon; V^{1/2}Q(K+\lambda)^{-1}V^{1/2}), \ \lambda > 0, \ \varepsilon \in (0,1),$$

where, as above, Q = I - P. Evidently,

$$(2.15) n_{+}(s; V^{1/2}PV^{1/2}) = n_{+}(s; PVP) = n_{+}(s; T_{V}), \quad s > 0,$$

while Proposition 2.1 easily implies that for any $\varepsilon > 0$ we have

(2.16)
$$n_{\pm}(\varepsilon; V^{1/2}Q(K+\lambda)^{-1}V^{1/2}) = O(1), \quad \lambda \to 0.$$

Putting together (2.13) and (2.14) - (2.16), we obtain the upper bound in (2.11). In order to prove (2.12), we recall that the generalized Birman-Schwinger principle (see e.g. [2, Theorem 1.3]), easily entails

(2.17)
$$\mathcal{N}_{+}(\lambda) = n_{-}(1; V^{1/2}(K - \lambda)^{-1}V^{1/2}) + O(1), \quad \lambda \downarrow 0.$$

By the Weyl inequalities, the estimates

$$n_{+}((1+\varepsilon)\lambda; V^{1/2}PV^{1/2}) - n_{+}(\varepsilon; V^{1/2}Q(K-\lambda)^{-1}V^{1/2}) \le$$

$$n_{-}(1; V^{1/2}(K - \lambda)^{-1}V^{1/2}) \le$$

(2.18)
$$n_{+}((1-\varepsilon)\lambda; V^{1/2}PV^{1/2}) + n_{-}(\varepsilon; V^{1/2}Q(K-\lambda)^{-1}V^{1/2})$$

hold true for every $\varepsilon \in (0,1)$. Now (2.17), (2.18), (2.15), and (2.16), imply (2.12).

П

3. General properties of harmonic Toeplitz operators

In this section we establish sufficient conditions which guarantee $T_V \in S_p$, the pth Schatten-von Neumann class, or $T_V \in S_{p,w}$, the weak counterpart of S_p .

We first introduce the notations we need. Let X and Y be separable Hilbert spaces. We denote by $\mathcal{L}(X,Y)$ (resp., $S_{\infty}(X,Y)$) the class of linear bounded (resp., compact) operators $T:X\to Y$. Let $T\in S_{\infty}(X,Y)$. Next, $S_p(X,Y)$, $p\in (0,\infty)$, is the class of compact operators $T:X\to Y$ for which the functional

$$||T||_p := \left(\operatorname{Tr} \left(T^*T\right)^{p/2}\right)^{1/p}$$

is finite.

Let $T \in S_{\infty}(X,Y)$. Pick s > 0, and bearing in mind notation (2.9), set

(3.1)
$$n_*(s;T) := n_+(s^2; T^*T).$$

Thus, $n_*(s;T)$ is the number of the singular values of the operator T, larger than s, and counted with their multiplicities. Evidently, if $T = T^* \in S_{\infty}(X,X)$, then

$$(3.2) n_+(s;T) \le n_*(s;T), \quad s > 0.$$

Further, $S_{p,\mathbf{w}}(X,Y)$, $p \in (0,\infty)$, is the class of operators $T \in S_{\infty}(X,Y)$ for which the functional

(3.3)
$$||T||_{p,\mathbf{w}} := \sup_{s>0} s \, n_*(s;T)^{1/p}$$

is finite.

If X = Y, we write $\mathcal{L}(X)$, $S_p(X)$, and $S_{p,w}(X)$, instead of $\mathcal{L}(X,X)$, $S_p(X,X)$, and $S_{p,w}(X,X)$, respectively. Moreover, whenever appropriate, we omit X and Y in the notations \mathcal{L} , S_p , and $S_{p,w}$.

Let us now turn to the study of the spectral properties of the harmonic Toeplitz operators $T_V = PV$. It is well known that the projection P onto $\mathcal{H}(\Omega)$ (see (2.2)) admits an integral kernel $\mathcal{R} \in C^{\infty}(\Omega \times \Omega)$, called the reproducing kernel of P (see e.g. [23, 14]). Thus

$$(Pu)(x) = \int_{\Omega} \mathcal{R}(x, y)u(y)dy, \quad x \in \Omega, \quad u \in L^{2}(\Omega).$$

Let $\{\varphi_j\}_{j\in\mathbb{N}}$ be an orthogonal basis in $\mathcal{H}(\Omega)$. Then

(3.4)
$$\mathcal{R}(x,y) = \sum_{j \in \mathbb{N}} \varphi_j(x) \overline{\varphi_j(y)}, \quad x, y \in \Omega,$$

the series being locally uniformly convergent in $\Omega \times \Omega$. Evidently, $\mathcal{R}(x,y)$ is independent of the choice of the basis $\{\varphi_j\}_{j\in\mathbb{N}}$. Moreover, the kernel \mathcal{R} is real-valued and symmetric. For $x\in\Omega$ put

$$\rho(x) := \mathcal{R}(x, x).$$

Then, (3.4) implies that

$$|\mathcal{R}(x,y)| \le \rho(x)^{1/2} \rho(y)^{1/2}, \quad x, y \in \Omega.$$

For $x, y \in \Omega$, set

(3.5)
$$r(x) := \operatorname{dist}(x, \partial\Omega), \quad \delta(x, y) := |x - y| + r(x) + r(y).$$

LEMMA 3.1. [23, Theorem 1.1] For any multiindices $\alpha, \beta \in \mathbb{Z}_+^d$ there exists a constant $C_{\alpha,\beta} \in (0,\infty)$ such that

$$(3.6) \left| D_x^{\alpha} D_y^{\beta} \mathcal{R}(x, y) \right| \le \frac{C_{\alpha, \beta}}{\delta(x, y)^{d + |\alpha| + |\beta|}}, \quad x, y \in \Omega.$$

Moreover, there exists a constant $C \in (0, \infty)$ such that

(3.7)
$$\varrho(x) \ge Cr(x)^{-d}, \quad x \in \Omega.$$

For a Borel set $A \subset \Omega$ set $\rho(A) := \int_{\mathcal{A}} \varrho(x) dx$. By (3.6) with $\alpha = \beta = 0$, and (3.7), ρ is an infinite σ -finite measure on Ω which is absolutely continuous with respect to the Lebesgue measure.

Our next goal is to establish conditions which guarantee $T_V \in S_p(\mathcal{H}(\Omega))$, $p \in [1, \infty)$, or $T_V \in S_{p,w}(\mathcal{H}(\Omega))$, $p \in (1, \infty)$. For $p \in (0, \infty)$ define $L^p_w(\Omega; d\rho)$ as the class of ρ -measurable functions $u : \Omega \to \mathbb{C}$ for which the quasinorm

$$||u||_{L^p_{\mathbf{w}}(\Omega;d\rho)} := \sup_{t>0} t\rho \left(\left\{ x \in \Omega \, | \, |u(x)| > t \right\} \right)^{1/p}$$

is finite.

Proposition 3.1. Let $\Omega \subset \mathbb{R}^d$, $d \geq 2$ be a bounded domain with boundary $\partial \Omega \in C^{\infty}$.

(i) Assume
$$V \in L^p(\Omega; d\rho)$$
, $p \in [1, \infty)$. Then $T_V \in S_p(\mathcal{H}(\Omega))$ and

$$||T_V||_p \le ||V||_{L^p(\Omega;d\rho)}.$$

(ii) Assume
$$V \in L^p_{\mathbf{w}}(\Omega; d\rho), p \in (1, \infty)$$
. Then $T_V \in S_{p, \mathbf{w}}(\mathcal{H}(\Omega))$ and

(3.9)
$$||T_V||_{p,\mathbf{w}} \le ||V||_{L^p_{\mathbf{w}}(\Omega;d\rho)}.$$

PROOF. Let us consider the operator PVP as defined on $L^2(\Omega)$. Evidently,

$$(3.10) ||T_V||_p = ||PVP||_p, ||T_V||_{p,w} = ||PVP||_{p,w}, p \in (0,\infty).$$

We have $PVP=F^*e^{i\arg V}F$ where $F:L^2(\Omega)\to L^2(\Omega)$ is the operator with integral kernel

$$|V(x)|^{1/2}\mathcal{R}(x,y), \quad x,y \in \Omega.$$

Assume $V \in L^1(\Omega; d\rho)$. Then

$$(3.11) ||PVP||_1 \le ||F^*||_2 ||e^{i\arg V}|| ||F||_2 = ||F||_2^2 = ||V||_{L^1(\Omega;d\rho)}.$$

Assume now $V \in L^{\infty}(\Omega; d\rho)$. Since ||P|| = 1 and $d\rho$ is absolutely continuous with respect to the Lebesgue measure,

$$(3.12) ||PVP|| \le ||V||_{L^{\infty}(\Omega)} = ||V||_{L^{\infty}(\Omega;d\rho)}.$$

Interpolating between (3.11) and (3.12), and applying [8, Theorem 3.1], we find that

$$||PVP||_p \le ||V||_{L^p(\Omega;d\rho)}, \quad p \in [1,\infty),$$

 $||PVP||_{p,w} \le ||V||_{L^p_w(\Omega;d\rho)}, \quad p \in (1,\infty),$

which combined with (3.10), implies (3.8) and (3.9).

Remark: Let $\mu \geq 0$ be a finite Borel measure on Ω . In this case, the harmonic Toeplitz operator T_{μ} is defined by

$$(T_{\mu}u)(x) := \int_{\Omega} \mathcal{R}(x,y)u(y)d\mu(y), \quad u \in \mathcal{H}(\Omega), \quad x \in \Omega.$$

If $d\mu(x) = V(x)dx$ with $0 \le V \in L^1(\Omega)$, then, of course, $T_{\mu} = T_V$. Criteria on μ which guarantee the boundedness of T_{μ} , the compactness of T_{μ} , or the inclusion $T_{\mu} \in S_p(\mathcal{H}(\Omega)), p \in [1, \infty)$, are contained in [14]. These criteria are formulated in terms of the Berezin transform $\tilde{\mu}$ of the measure μ , defined by

$$\tilde{\mu}(x) := \varrho(x)^{-1} \int_{\Omega} \mathcal{R}(x, y)^2 d\mu(y), \quad x \in \Omega.$$

The combination of Theorem 2.2 with Proposition 3.1 entails the following

Corollary 3.1. Let Ω be a bounded domain with boundary $\partial \Omega \in C^{\infty}$. Assume that $V \in C(\overline{\Omega}; \mathbb{R}) \cap L^p_{\mathrm{w}}(\Omega; d\rho)$, $p \in (1, \infty)$. Then for any $\varepsilon \in (0, 1)$ we have

(3.13)
$$\mathcal{N}_{\pm}(\lambda) \le (1 - \varepsilon)^{-p} \lambda^{-p} \|V\|_{L_{w}^{p}(\Omega;d\rho)}^{p} + O(1),$$

for sufficiently small $\lambda > 0$.

PROOF. Estimate (3.13) follows immediately from (2.11) - (2.12), combined with (3.2) and (3.3). \Box

At the end of this section we show that if the symbol V is compactly supported in Ω , then $T_V \in S_p$ for any $p \in (0, \infty)$, even if the behavior of V is quite irregular. In fact, we will replace in this case V by $\phi \in \mathcal{E}'(\Omega)$, the class of distributions over $\mathcal{E}(\Omega) := C^{\infty}(\Omega)$. We recall that $\phi \in \mathcal{D}'(\Omega)$ is in $\mathcal{E}'(\Omega)$, if and only if supp ϕ is compact in Ω . If $\phi \in \mathcal{E}'(\Omega)$, we define $T_{\phi} : \mathcal{H}(\Omega) \to \mathcal{H}(\Omega)$ as the operator with integral kernel

$$K_{\phi}(x,y) := (\phi, \mathcal{R}(x,\cdot)\mathcal{R}(\cdot,y)), \quad x,y \in \Omega,$$

where (\cdot,\cdot) denotes the pairing between $\mathcal{E}'(\Omega)$ and $\mathcal{E}(\Omega)$. Of course, if $\phi = \mu$ and $\mu \geq 0$ is a finite Borel measure such that supp μ is compact in Ω , then $T_{\phi} = T_{\mu}$.

PROPOSITION 3.2. Let $\Omega \subset \mathbb{R}^d$, $d \geq 2$, be a bounded domain with boundary $\partial \Omega \in C^{\infty}$. Assume that $\phi \in \mathcal{E}'(\Omega)$. Then we have $T_{\phi} \in S_p(\mathcal{H}(\Omega))$ for any $p \in (0,\infty)$, and, hence,

$$(3.14) n_*(\lambda; T_\phi) = O(\lambda^{-\alpha}), \quad \lambda \downarrow 0,$$

for any $\alpha \in (0, \infty)$.

PROOF. Since supp ϕ is compact in Ω , we have $K_{\phi} \in C^{\infty}(\overline{\Omega} \times \overline{\Omega})$. Therefore, (3.14) follows immediately from, say, [8, Proposition 2.1].

Remarks: (i) In Section 5 we will show that if Ω is the unit ball in \mathbb{R}^d , and $V \geq 0$ is compactly supported, and possesses a partial radial symmetry, then the eigenvalues of T_V decay exponentially fast.

(ii) Harmonic Toeplitz operators T_{ϕ} with $\phi \in \mathcal{E}'(\Omega)$ were considered in [3] where, in particular, it was proved that rank $T_{\phi} < \infty$, if and only if supp ϕ is finite.

4. Spectral asymptotics of T_V for V of power-like decay at $\partial\Omega$

4.1. Statement of the main results. In this section we assume that V is smooth and positive near the boundary, and admits a power-like decay at $\partial\Omega$, and investigate the asymptotic behavior of the discrete spectrum of T_V near the origin. We obtain the main asymptotic term of $n_+(\lambda; T_V)$ as $\lambda \downarrow 0$, and give a sharp estimate of the remainder (see Theorems 4.1 and 4.2 below).

In what follows we consider $\partial\Omega$ as a compact (d-1)-dimensional Riemannian manifold with metric tensor $g(y):=\{g_{jk}(y)\}_{j,k=1}^{d-1},\ y\in\partial\Omega$, generated by the Euclidean metrics in \mathbb{R}^d , and denote by dS the measure induced by g on $\partial\Omega$.

Let $a, \tau \in C^{\infty}(\overline{\Omega})$ satisfy a > 0 on $\overline{\Omega}$, $\tau > 0$ on Ω , and $\tau = r := \operatorname{dist}(\cdot, \partial\Omega)$ (see (3.5)) in a neighborhood of $\partial\Omega$. Assume that

$$(4.1) V(x) = \tau(x)^{\gamma} a(x), \quad \gamma \ge 0, \quad x \in \Omega.$$

Set $a_0 := a_{|\partial\Omega}$

THEOREM 4.1. Assume that V satisfies (4.1) with $\gamma > 0$. Then we have

(4.2)
$$n_{+}(\lambda; T_{V}) = \mathcal{C} \lambda^{-\frac{d-1}{\gamma}} \left(1 + O(\lambda^{\frac{1}{\gamma}}) \right), \quad \lambda \downarrow 0,$$

where

(4.3)
$$\mathcal{C} := \omega_{d-1} \left(\frac{\Gamma(\gamma+1)^{\frac{1}{\gamma}}}{4\pi} \right)^{d-1} \int_{\partial\Omega} a_0(y)^{\frac{d-1}{\gamma}} dS(y),$$

and $\omega_n = \pi^{n/2}/\Gamma(1+n/2)$ is the Lebesgue measure of the unit ball in \mathbb{R}^n , $n \ge 1$.

The proof of Theorem 4.1 can be found in Subsection 4.2.

The assumption of Theorem 4.1 that V satisfies (4.1) in the whole domain Ω is rather restrictive. In the following theorem we consider Toeplitz operators with symbols which satisfy (4.1) with $\gamma > 0$ only in a neighborhood of $\partial \Omega$ while they can have quite an irregular behavior on a compact subset of Ω ; in particular, away from the boundary, these symbols are not obliged to be smooth and positive.

THEOREM 4.2. Let V satisfy the assumptions of Theorem 4.1, and $\phi \in \mathcal{E}'(\Omega; \mathbb{R})$. Then we have

$$(4.4) n_{+}(\lambda; T_{V+\phi}) = \mathcal{C} \,\lambda^{-\frac{d-1}{\gamma}} \left(1 + O(\lambda^{\frac{\varepsilon}{\gamma}}) \right), \quad \lambda \downarrow 0,$$

where $T_{V+\phi} := T_V + T_{\phi}$, C is the constant defined in (4.3), $\varepsilon = 1$ if $d \geq 3$, and $\varepsilon < 1$ is arbitrary if d = 2.

The proof of Theorem 4.2 is contained in Subsection 4.3.

Remark: Let $d \geq 3$. Then Theorem 4.2 implies that (4.2) remains true if we replace T_V by $T_{V+\phi}$ with $\phi \in \mathcal{E}'(\Omega; \mathbb{R})$. In particular, (4.2) is valid if $V \in L^1_{\text{loc}}(\Omega; \mathbb{R})$ satisfies (4.1) with $\gamma > 0$ only in a neighborhood of $\partial \Omega$. If d = 2, the prize we have to pay for the substitution of T_V by $T_{V+\phi}$ is that the remainder estimate in (4.2) is better than in (4.4). However, as mentioned in the remark at the end of Subsection 4.3, if d = 2 and the distribution $\phi \in \mathcal{E}'(\Omega)$ is non-negative, then (4.4) holds true also with $\varepsilon = 1$.

Combining Theorems 2.1 and 4.2, we obtain the following

COROLLARY 4.1. Assume that $0 \le V \in C(\overline{\Omega})$, and (4.1) with $\gamma > 0$ holds true in a neighborhood of $\partial\Omega$. Then we have

(4.5)
$$\lim_{\lambda \downarrow 0} \lambda^{\frac{d-1}{\gamma}} \mathcal{N}_{\pm}(\lambda) = \mathcal{C},$$

C being the constant defined in (4.3).

Remark: Assume the hypotheses of Corollary 4.1. Then, recalling (3.6) with $\alpha = \beta = 0$, and (3.7), we find that $V \in L^p_{\mathbf{w}}(\Omega; d\rho)$ if and only if $p = \frac{d-1}{\gamma}$. Thus, if $\gamma < d-1$, then (4.5) implies that the order of our estimate (3.13) is sharp.

4.2. Proof of Theorem 4.1. Assume that $f \in H^s(\partial\Omega)$, $s \in \mathbb{R}$. Then the boundary-value problem

(4.6)
$$\begin{cases} \Delta u = 0 & \text{in } \Omega, \\ u = f & \text{on } \partial\Omega, \end{cases}$$

admits a unique solution $u \in H^{s+1/2}(\Omega)$, we have

(4.7)
$$||u||_{H^{s+1/2}(\Omega)} \times ||f||_{H^s(\partial\Omega)},$$

and, therefore, the mapping $f \mapsto u$ defines an isomorphism between $H^s(\partial\Omega)$ and $H^{s+1/2}(\Omega)$ (see [25, Sections 5, 6, 7, Chapter 2]).

If s = 0, we set

$$(4.8) u = Gf.$$

By (4.7) with s=0, and the compactness of the embedding of $H^{1/2}(\Omega)$ into $L^2(\Omega)$, we find that the operator $G:L^2(\partial\Omega)\to L^2(\Omega)$ is compact. By [16, Theorem 12, Section 2.2], we have

(4.9)
$$u(x) = \int_{\partial\Omega} \mathcal{K}(x, y) f(y) dS(y), \quad x \in \Omega,$$

where

(4.10)
$$\mathcal{K}(x,y) := -\frac{\partial \mathcal{G}}{\partial \nu_y}(x,y), \quad x \in \Omega, \quad y \in \partial \Omega,$$

 \mathcal{G} is the Dirichlet Green function associated with Ω . Note that

$$(4.11) \mathcal{K} \in C^{\infty}(\Omega \times \partial \Omega).$$

Lemma 4.1. We have

(4.12)
$$\operatorname{Ker} G = \{0\},\$$

$$(4.13) \overline{\operatorname{Ran} G} = \mathcal{H}(\Omega).$$

PROOF. Relation (4.12) follows from (4.7) with s=0. Let us check (4.13). Pick $u\in\mathcal{H}(\Omega)$. Then, by (4.6) with s=-1/2, we have $f:=u_{|\partial\Omega}\in H^{-1/2}(\partial\Omega)$. Let $f_n\in L^2(\partial\Omega),\, n\in\mathbb{N}$, and

(4.14)
$$\lim_{n \to \infty} ||f_n - f||_{H^{-1/2}(\partial\Omega)} = 0.$$

Set $u_n := Gf_n$. Then $u_n \in \operatorname{Ran} G$, $n \in \mathbb{N}$, and by (4.7) with s = -1/2, and (4.14), we have $\lim_{n \to \infty} \|u_n - u\|_{L^2(\Omega)} = 0$ which implies (4.13).

Set $J:=G^*G$. Then the operator $J=J^*\geq 0$ is compact in $L^2(\partial\Omega)$. Due to (4.12), we have $\operatorname{Ker} J=\{0\}$. Let $\{\lambda_j\}_{j\in\mathbb{N}}$ be the non-increasing sequence of the eigenvalues $\lambda_j>0$ of J, and let $\{\phi_j\}_{j\in\mathbb{N}}$ be the corresponding orthonormal eigenbasis in $L^2(\partial\Omega)$ with $J\phi_j=\lambda_j\phi_j,\,j\in\mathbb{N}$. Define the operator J^{-1} , self-adjoint in $L^2(\partial\Omega)$, by (4.15)

$$J^{-1}u := \sum_{j \in \mathbb{N}} \lambda_j^{-1} \langle u, \phi_j \rangle \phi_j, \quad \text{Dom } J^{-1} := \left\{ u \in L^2(\partial\Omega) \left| \sum_{j \in \mathbb{N}} \lambda_j^{-2} \left| \langle u, \phi_j \rangle \right|^2 < \infty \right\},$$

 $\langle \cdot, \cdot \rangle$ being the scalar product in $L^2(\partial \Omega)$. Evidently, $\overline{JJ^{-1}} = J^{-1}J = I$.

Further, introduce the polar decomposition of the operator $G = U|G| = UJ^{1/2}$ where $U: L^2(\partial\Omega) \to L^2(\Omega)$ is an isometric operator. By Lemma 4.1, we have $\text{Ker } U = \{0\}$ and $\text{Ran } U = \mathcal{H}(\Omega)$. Thus, we obtain the following

PROPOSITION 4.1. The orthogonal projection P onto $\mathcal{H}(\Omega)$ satisfies

$$(4.16) P = GJ^{-1}G^* = UU^*.$$

Assume that V satisfies (4.1) with $\gamma \geq 0$, and set $J_V := G^*VG$; from this point of view, we have $J = J_1$.

PROPOSITION 4.2. Let V satisfy (4.1) with $\gamma \geq 0$. Then the operator T_V is unitarily equivalent to (the closure of) the operator $J^{-1/2}J_VJ^{-1/2}$.

PROOF. By (4.16), we have

$$PVP = UJ^{-1/2}G^*VGJ^{-1/2}U^* = UJ^{-1/2}J_VJ^{-1/2}U^*.$$

and the operator U maps unitarily $L^2(\partial\Omega)$ onto $\mathcal{H}(\Omega)$.

Let $y \in \partial \Omega$ and $\eta \in T_{\eta}^* \partial \Omega$. Set

$$|\eta| = |\eta|_y := \left(\sum_{j,k=1}^{d-1} g^{jk}(y)\eta_j\eta_k\right)^{1/2},$$

where $\{g^{jk}(y)\}_{j,k=1}^{d-1}$ is the matrix inverse to g(y).

Proposition 4.3. Under the assumptions of Proposition 4.2 the operator $J^{-1/2}J_VJ^{-1/2}$ is a ΨDO with principal symbol

$$(4.17) 2^{-\gamma}\Gamma(\gamma+1)|\eta|^{-\gamma}a_0(y), \quad (y,\eta) \in T^*\partial\Omega.$$

PROOF. Using the pseudo-differential calculus due to L. Boutet de Monvel (see [11, 12]), M. Engliš showed recently in [15, Sections 6, 7] that if V satisfies (4.1) with $\gamma \geq 0$, then the operator J_V is a Ψ DO with principal symbol

$$2^{-\gamma-1}\Gamma(\gamma+1)|\eta|^{-\gamma-1}a_0(y), \quad (y,\eta) \in T^*\partial\Omega.$$

In particular, $J = J_1$ is a Ψ DO with principal symbol $2^{-1}|\eta|^{-1}$. Then the pseudo-differential calculus (see e.g. [28, Chapters I, II]) easily implies that $J^{-1/2}$ is a Ψ DO with principal symbol $2^{1/2}|\eta|^{1/2}$, and $J^{-1/2}J_VJ^{-1/2}$ is a Ψ DO with principal symbol defined in (4.17).

Now we are in position to prove Theorem 4.1. It is easy to see that under its assumptions we have Ker $J^{-1/2}J_VJ^{-1/2}=\{0\}$. Using the spectral theorem, define the operator

$$A := \left(J^{-1/2} J_V J^{-1/2}\right)^{-1/\gamma}$$

(cf. (4.15)). Then, by the pseudo-differential calculus, A is a $\Psi {\rm DO}$ with principal symbol

$$2\Gamma(\gamma+1)^{-1/\gamma}|\eta|a_0(y)^{-1/\gamma}, \quad (y,\eta) \in T^*\partial\Omega.$$

By Proposition 4.2 and the spectral theorem, we have

(4.18)
$$n_{+}(\lambda; T_{V}) = n_{+}(\lambda; J^{-1/2}J_{V}J^{-1/2}) = \operatorname{Tr} \mathbb{1}_{(-\infty, \lambda^{-1/\gamma})}(A), \quad \lambda > 0.$$

A classical result of L. Hörmander $[\mathbf{22}]$ easily implies that

$$\text{Tr } \mathbf{1}_{(-\infty,E)}(A) =$$

$$(2\pi)^{-d+1} \left| \left\{ (y,\eta) \in T^* \partial \Omega \mid 2\Gamma(\gamma+1)^{-1/\gamma} |\eta| a_0(y)^{-1/\gamma} < E \right\} \right| + O(E^{-(d-2)}) =$$

$$(4.19) \qquad \mathcal{C}E^{d-1}(1 + O(E^{-1})), \quad E \to \infty,$$

where $|\cdot|$ is the Lebesgue measure on $T^*\partial\Omega$, and \mathcal{C} is the constant defined in (4.3). Combining (4.18) and (4.19), we arrive at (4.2).

Remark: The natural idea to parametrize the functions $u \in \mathcal{H}(\Omega)$ by their restrictions on $\partial\Omega$ has been used in the theory of harmonic Toeplitz operators and related areas by various authors; it could be traced back at least to the classical work [11], and has been recently applied in [15] in order to obtain a suitable representation of the operator J_V . We would like also to mention here the article [9] where the authors consider the operator generated by the ratio of two quadratic differential forms defined on the solutions of a homogeneous elliptic equation. The order of the numerator is lower than the order of the denominator, and, since the

domain considered is supposed to be bounded and to have a regular boundary, the operator generated by the ratio is compact.

The harmonic Toeplitz operator T_V could be interpreted as the operator generated by the quadratic-form ratio

(4.20)
$$\frac{\int_{\Omega} V|u|^2 dx}{\int_{\Omega} |u|^2 dx}, \quad u \in \mathcal{H}(\Omega).$$

Note that both the numerator and the denominator in (4.20) are of zeroth order, and the compactness of T_V is now due to the fact that V vanishes at $\partial\Omega$.

In spite of the differences between the operators considered in [9], and the harmonic Toeplitz operators studied here, the unitary equivalence of T_V and $J^{-1/2}J_VJ^{-1/2}$ established in our Proposition 4.2 has much in common with the reduction to a Ψ DO on $\partial\Omega$, performed in [9].

4.3. Proof of Theorem **4.2.** The Weyl inequalities (2.10) imply

$$n_{+}(\lambda(1+\lambda^{\theta});T_{V}) - n_{-}(\lambda^{1+\theta};T_{\phi}) \le n_{+}(\lambda;T_{V+\phi}) <$$

(4.21)
$$n_{+}(\lambda(1-\lambda^{\theta});T_{V}) + n_{+}(\lambda^{1+\theta};T_{\phi}),$$

for $\lambda \in (0,1)$ and $\theta > 0$. By (4.2),

$$n_{+}(\lambda(1\pm\lambda^{\theta});T_{V})=$$

$$(4.22) \quad \mathcal{C}\left(\lambda(1\pm\lambda^{\theta})\right)^{-\frac{d-1}{\gamma}} + O\left(\lambda^{-\frac{d-2}{\gamma}}\right) = \mathcal{C}\lambda^{-\frac{d-1}{\gamma}} + O\left(\lambda^{-\frac{d-2}{\gamma}}\right), \quad \lambda \in (0,1),$$

provided that $\theta > 1/\gamma$. Next, by estimate (3.14), we have

(4.23)
$$n_{\pm}(\lambda^{1+\theta}; T_{\phi}) = O(\lambda^{-\alpha(1+\theta)}), \lambda > 0,$$

for any $\alpha \in (0, \infty)$. Assume $d \geq 3$ and choose $\alpha \in \left(0, \frac{d-2}{\gamma(1+\theta)}\right)$. Then (4.4) follows from (4.21) - (4.23). If d=2, then we can pick any $\varepsilon > 0$ and choose $\alpha \in \left(0, \frac{1-\varepsilon}{\gamma(1+\theta)}\right)$, in order to check that in this case (4.21) – (4.23) again imply (4.4).

Remark: Arguing as in the proof of Theorem 4.1 (see Propositions 4.2 and 4.3), we can show that $T_{V+\phi}$ with $\phi \in \mathcal{E}'(\Omega;\mathbb{R})$ is unitarily equivalent to a self-adjoint Ψ DO with principal symbol defined in (4.17). The only problem to extend in a straightforward manner our proof of Theorem 4.1 to $T_{V+\phi}$ is that this operator may have a non trivial kernel unless, for example, $\phi \geq 0$. In particular, if d=2 and $\phi \in \mathcal{E}'(\Omega;\mathbb{R})$ satisfies $\phi \geq 0$, then (4.4) holds also for $\varepsilon = 1$.

5. Spectral asymptotics of T_V for compactly supported V

In this section we consider the asymptotics of $n_+(\lambda; T_V)$ as $\lambda \downarrow 0$ in the case where V is compactly supported in Ω , i.e. when V vanishes identically in a neighborhood of $\partial\Omega$. In this case T_V admits an integral kernel which is in the class $C^{\infty}(\overline{\Omega} \times \overline{\Omega})$, and T_V can be considered as a Ψ DO of order $-\infty$.

Set

$$B_R := \{ x \in \mathbb{R}^d \mid |x| < R \}, \quad d \ge 2, \quad R \in (0, \infty).$$

Since we are still unable to handle the case of general bounded Ω and compactly supported V, we suppose that Ω is the unit ball B_1 in \mathbb{R}^d while supp V coincides with B_c with $c \in (0,1)$. Using the known fact that if V is proportional to $\mathbb{1}_{B_c}$, then the eigenvalues of T_V can be calculated explicitly, we obtain the main asymptotic term of $n_+(\lambda; T_V)$ as $\lambda \downarrow 0$, for generic T_V such that supp $V = B_c$.

Let $\Omega = B_1$. Thus, $\partial \Omega = \mathbb{S}^{d-1} := \{x \in \mathbb{R}^d \mid |x| = 1\}$. The space $\mathcal{H}(B_1)$ admits an explicit orthonormal eigenbasis which we are now going to describe. Recall that $k(k+d-2), k \in \mathbb{Z}_+$, are the eigenvalues of the Beltrami-Laplace operator $-\Delta_{\mathbb{S}^{d-1}}$, self-adjoint in $L^2(\mathbb{S}^{d-1})$ (see e.g. [28, Section 22]). Moreover,

$$\dim \mathrm{Ker} \; (-\Delta_{\mathbb{S}^{d-1}} - k(k+d-2)I) =: m_k = \binom{d+k-1}{d-1} - \binom{d+k-3}{d-1}$$

where $\binom{m}{n} = \frac{m!}{(m-n)! \, n!}$ if $m \ge n$, and $\binom{m}{n} = 0$ if m < n (see e.g. [28, Theorem 22.1]). Set

$$M_k := {d+k-1 \choose d-1} + {d+k-2 \choose d-1}, \quad k \in \mathbb{Z}.$$

Evidently,

(5.1)
$$M_k = \frac{2k^{d-1}}{(d-1)!} \left(1 + O\left(k^{-1}\right) \right), \quad k \to \infty,$$

(see e.g. [1, Eq. 6.1.47]). By induction, we easily find that

(5.2)
$$\sum_{j=0}^{k} m_j = M_k, \quad k \in \mathbb{Z}_+.$$

Let $\psi_{k,\ell}$, $\ell = 1, \ldots, m_k$, be an orthonormal basis in Ker $(-\Delta_{\mathbb{S}^{d-1}} - k(k+d-2)I)$, $k \in \mathbb{Z}_+$. It is well known that $\psi_{k,\ell}$ are restrictions on \mathbb{S}^{d-1} of homogeneous polynomials of degree k, harmonic in \mathbb{R}^d (see e.g [28, Section 22]). Then the functions

$$\phi_{k,\ell}(x) := \sqrt{2k+d} |x|^k \psi_{k,\ell}(x/|x|), \quad x \in B_1, \quad \ell = 1, \dots, m_k, \quad k \in \mathbb{Z}_+,$$

form an orthonormal basis in $\mathcal{H}(B_1)$. Let $\mathcal{H}_k(B_1)$, $k \in \mathbb{Z}_+$, be the subspace of $\mathcal{H}(B_1)$ generated by $\phi_{k,\ell}$, $\ell = 1, \ldots, m_k$.

Further, let V(x) = v(|x|), $x \in B_1$, and let $v : [0,1) \to \mathbb{R}$ satisfy $\lim_{r \uparrow 1} v(r) = 0$, $v \in L^1((0,1); r^{d-1}dr)$. Then T_V is self-adjoint and compact in $\mathcal{H}(B_1)$, and

$$(5.3) T_V u = \mu_k u, \quad u \in \mathcal{H}_k(B_1),$$

where

(5.4)
$$\mu_k(v) := (2k+d) \int_0^1 v(r) r^{2k+d-1} dr, \quad k \in \mathbb{Z}_+.$$

Set

$$\xi(s; v) = \# \{k \in \mathbb{Z}_+ \mid \mu_k(v) > s\}, \quad s > 0.$$

Let us calculate the eigenvalues of T_V in the simple model situation where $v(r) = b \, \mathbb{1}_{[0,c]}(r), \, r \in [0,1)$, with b > 0 and $c \in (0,1)$. Then (5.4) implies

$$\mu_k(v) = b c^{2k+d}, \quad k \in \mathbb{Z}_+.$$

Evidently, the sequence $\{\mu_k(v)\}_{k\in\mathbb{Z}_+}$ is decreasing. Setting $V(x):=v(|x|),\,x\in\mathbb{R}^d$, we get

(5.6)
$$n_{+}(\lambda; T_{V}) = M_{\xi(\lambda; v)-1}, \quad \lambda > 0.$$

Let us discuss the asymptotics of $n_{+}(\lambda; T_{V})$ as $\lambda \downarrow 0$. By (5.5),

(5.7)
$$\xi(\lambda; v) = \frac{1}{2} \frac{|\ln \lambda|}{|\ln c|} + O(1), \quad \lambda \downarrow 0.$$

By (5.6), (5.1), and (5.7), we get

(5.8)
$$n_{+}(\lambda; T_{V}) = \frac{2^{-d+2}}{(d-1)! |\ln c|^{d-1}} |\ln \lambda|^{d-1} + O(|\ln \lambda|^{-d+2}), \quad \lambda \downarrow 0.$$

Remark: The fact that the basis $\{\phi_{k,\ell}\}$ diagonalizes the operator T_V with radially symmetric symbol V, acting in $\mathcal{H}(B_1)$, was noted in [27, Part 2.3.2], and was used there, in particular, to obtain asymptotic relations of type (5.8). The fact that the Toeplitz operators with radially symmetric symbols, acting in the holomorphic Fock-Segal-Bargmann space, are diagonalized in a certain canonic basis, was utilized already in [26, 21]. A similar result concerning Toeplitz operators with radially symmetric symbols, acting in the holomorphic Bergman space, can be found in [20].

Next, we use (5.8) in order to study the spectral asymptotics for Toeplitz operators with symbols V which possess partial radial symmetry.

THEOREM 5.1. Let $\Omega = B_1$. Assume that $V: B_1 \to [0, \infty)$ satisfies $V \in L^{\infty}(B_1)$ and supp $V = \overline{B_c}$ for some $c \in (0,1)$. Suppose moreover that for any $\delta \in (0,c)$ we have ess $\inf_{x \in B_{\delta}} V(x) > 0$. Then

(5.9)
$$\lim_{\lambda \downarrow 0} |\ln \lambda|^{-d+1} n_{+}(\lambda; T_{V}) = \frac{2^{-d+2}}{(d-1)! |\ln c|^{d-1}}.$$

PROOF. Pick $\delta \in (0, c)$. Then for almost every $x \in B_1$ we have

$$b_{-} \mathbf{1}_{B_{\delta}}(x) \leq V(x) \leq b_{+} \mathbf{1}_{B_{\delta}}(x),$$

where

$$b_{-} := \operatorname{ess\,inf}_{x \in B_{\delta}} V(x), \quad b_{+} := \operatorname{ess\,sup}_{x \in B_{1}} V(x).$$

Then the mini-max principle and (5.8) imply

$$\frac{2^{-d+2}}{(d-1)!|\ln\delta|^{d-1}} \le \lim\inf_{\lambda\downarrow 0} |\ln\lambda|^{-d+1} \, n_+(\lambda;T_V) \le \limsup_{\lambda\downarrow 0} |\ln\lambda|^{-d+1} \, n_+(\lambda;T_V) \le$$

$$\frac{2^{-d+2}}{(d-1)!|\ln c|^{d-1}}.$$

Letting $\delta \uparrow c$, we obtain (5.9).

Remark: Hopefully, in a future work we will extend the result of Theorem 5.1 to more general domains Ω , and more general compactly supported V.

Putting together Theorems 2.2 and 5.1, we obtain the following

COROLLARY 5.1. Let $\Omega = B_1 \subset \mathbb{R}^d$, $d \geq 2$, $0 \leq V \in C(\overline{B}_1)$. Assume that $\sup V = \overline{B_c}$ for some $c \in (0,1)$, and that for any $\delta \in (0,c)$ we have $\inf_{x \in B_\delta} V(x) > 0$. Then

$$\lim_{\lambda \downarrow 0} |\ln \lambda|^{-d+1} \mathcal{N}_{\pm}(\lambda) = \frac{2^{-d+2}}{(d-1)! |\ln c|^{d-1}}.$$

References

- M.Abramowitz, I.Stegun, Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables, National Bureau of Standards, Applied Mathematics Series 55, 1964
- S. Alama, P. A. Deift, R. Hempel, Eigenvalue branches of the Schrödinger operator H λW in a gap of σ(H), Comm. Math. Phys. 121 (1989), 291-321.
- A. ALEXANDROV, G. ROZENBLUM, Finite rank Toeplitz operators: some extensions of D. Luecking's theorem, J. Funct. Anal. 256 (2009), 2291–2303.
- A. Alonso, B. Simon, The Birman-Krein-Vishik theory of self-adjoint extensions of semibounded operators, J. Operator Theory 4 (1980), 251–270.
- M. S. ASHBAUGH, F. GESZTESY, M. MITREA, R. SHTERENBERG, G. TESCHL, The Krein-von Neumann extension and its connection to an abstract buckling problem, Math. Nachr. 283 (2010), 165–179.
- M. S. ASHBAUGH, F. GESZTESY, M. MITREA, G. TESCHL, Spectral theory for perturbed Krein Laplacians in nonsmooth domains, Adv. Math. 223 (2010), 1372–1467.
- M. Sh. Birman, On the spectrum of singular boundary-value problems, Mat. Sb. (N.S.) 55 (1961), 125-174 (Russian); English translation in: Eleven Papers on Analysis, AMS Transl. 53, 23-80, AMS, Providence, R.I., 1966.
- 8. M. Š. BIRMAN, M. Z. SOLOMJAK, Estimates for the singular numbers of integral operators, (Russian) Uspehi Mat. Nauk **32** (1977), 17–84; English translation in: Russ. Math. Surveys **32**, (1977), 15–89.
- M. Š. BIRMAN, M. Z. SOLOMJAK, Asymptotic behavior of the spectrum of variational problems on solutions of elliptic equations, (Russian) Sibirsk. Mat. Zh. 20 (1979), 3–22; English translation in: Siberian Math. J. 20 (1979), 1–15.
- M.Š.BIRMAN, M.Z.SOLOMJAK, Spectral Theory of Self-Adjoint Operators in Hilbert Space,
 D. Reidel Publishing Company, Dordrecht, 1987.
- 11. L. BOUTET DE MONVEL, Boundary problems for pseudo-differential operators, Acta Math. 126 (1971), 11–51.
- L. BOUTET DE MONVEL, V. GUILLEMIN, The Spectral Theory of Toeplitz Operators, Annals
 of Mathematics Studies, 99 Princeton University Press, Princeton, NJ; University of Tokyo
 Press, Tokyo, 1981.
- V. Bruneau, G. Raikov, Spectral properties of harmonic Toeplitz operators and applications to the perturbed Krein Laplacian, Preprint: arXiv:1609.08229 (2016).
- B. R. CHOE, Y. J. LEE, K. NA, Toeplitz operators on harmonic Bergman spaces, Nagoya Math. J. 174 (2004), 165–186.
- M. Engliš, Boundary singularity of Poisson and harmonic Bergman kernels, J. Math. Anal. Appl. 429 (2015), 233–272.
- L. C. EVANS, Partial Differential Equations, Second edition. Graduate Studies in Mathematics, 19 American Mathematical Society, Providence, RI, 2010.
- L. GÅRDING, Some points of analysis and their history, University Lecture Series, 11 AMS, Providence, RI; Higher Education Press, Beijing, 1997.
- G. GRUBB, A characterization of the non-local boundary value problems associated with an elliptic operator, Ann. Scuola Norm. Sup. Pisa 22 (1968), 425–513.
- 19. G. Grubb, Spectral asymptotics for the "soft" selfadjoint extension of a symmetric elliptic differential operator, J. Operator Theory 10 (1983), 9–20.
- 20. S. GRUDSKY, A. KARAPETYANTS, N. VASILEVSKI, Toeplitz operators on the unit ball in \mathbb{C}^n with radial symbols, J. Operator Theory 49 (2003), 325–346.
- S. N. GRUDSKY, N. L. VASILEVSKI, Toeplitz operators on the Fock space: radial component effects, Integral Equations Operator Theory 44 (2002), 10–37.

- L. HÖRMANDER, The spectral function of an elliptic operator, Acta Math. 121 (1968), 193– 218
- H. KANG, H. KOO, Estimates of the harmonic Bergman kernel on smooth domains, J. Funct. Anal. 185 (2001), 220–239.
- M. Krein, The theory of self-adjoint extensions of semi-bounded Hermitian transformations and its applications, I, (Russian) Mat. Sb. N.S. 20 (1947), 431–495.
- J.-L LIONS, E. MAGENES, Non-homogeneous Boundary Value Problems and Applications.
 I, Vol. I. Die Grundlehren der mathematischen Wissenschaften, 181 Springer-Verlag, New York-Heidelberg, 1972.
- G. RAIKOV, S. WARZEL, Quasi-classical versus non-classical spectral asymptotics for magnetic Schrödinger operators with decreasing electric potentials, Rev. Math. Phys. 14 (2002), 1051–1072.
- G. ROZENBLUM, On lower eigenvalue bounds for Toeplitz operators with radial symbols in Bergman spaces, J. Spectr. Theory 1 (2011), 299–325.
- 28. M.A.Shubin, *Pseudodifferential Operators and Spectral Theory*, Second Edition, Berlin etc.: Springer-Verlag (2001).
- J. VON NEUMANN, Allgemeine Eigenwerttheorie Hermitescher Funktionaloperatoren, Math. Ann. 102 (1930), 49–131.
- H, Weyl, The method of orthogonal projection in potential theory, Duke Math. J. 7 (1940), 411–444.

UMR CNRS 5251, Institut de Mathématiques de Bordeaux, Université de Bordeaux, 351, Cours de la Libération, 33405 Talence, France

E-mail address: vbruneau@math.u-bordeaux.fr

FACULTAD DE MATEMÁTICAS, PONTIFICIA UNIVERSIDAD CATÓLICA DE CHILE, VICUÑA MACKENNA 4860, SANTIAGO DE CHILE

E-mail address: graikov@mat.uc.cl