# The $p$-version of the boundary element method for a three-dimensional crack problem 

Alexei Bespalov * Norbert Heuer ${ }^{\dagger}$


#### Abstract

We study the $p$-version of the boundary element method for a crack problem in linear elasticity with Dirichlet boundary conditions. The unknown jump of the traction has strong edge singularities and is approximated by solving an integral equation of the first kind with weakly singular operator. We prove a quasi-optimal a priori error estimate in the energy norm. For sufficiently smooth given data this gives a convergence like $c p^{-1+\varepsilon}$ with $\varepsilon>0$. Here, $p$ denotes the polynomial degree of the piecewise polynomial functions used to approximate the unknown.


Key words: $p$-version, boundary element method, linear elasticity, singularities
AMS Subject Classification: 41A10, 65N15, 65N38

## 1 Introduction and formulation of the problem

We analyze the convergence of the $p$-version of the boundary element method (BEM) with weakly singular integral operator for problems in $\mathbf{R}^{3}$. That is we study approximation properties of piecewise polynomial functions on surfaces in a negative order Sobolev space (order $-1 / 2$ ). To the knowledge of the authors this is the first paper dealing with this case. The $p$-version of the finite element method and the $p$-version of the BEM on curves have been widely studied. For the $p$-version of the BEM dealing

[^0]with problems in three dimensions, however, there are very few results. The case of hypersingular operators on polyhedral surfaces (the energy space is $H^{1 / 2}$ ) is analyzed in [7]. There, using $H^{1}$-regularity of the solution, the optimal convergence of the $p$-version has been shown. In [2] we consider hypersingular operators on open surface, where no $H^{1}$-regularity can be assumed, and prove optimal a priori error estimates. The case of weakly singular integral operators on surfaces has been an open problem so far. Here we study this situation for the model problem of linear elasticity with a crack that has a smooth boundary. The solution exhibits in general strong edge singularities not being $L_{2}$-regular.

Let us recall the Sobolev spaces used. Then we formulate the model problem. Let $\Gamma$ be an open smooth surface in $\mathbf{R}^{3}$ with smooth boundary curve $\gamma$. Taking a closed smooth surface $\tilde{\Gamma}$ which contains $\Gamma$, we consider Sobolev spaces $H^{t}(\tilde{\Gamma})$ for $t>0$ being the restriction of $H^{t+1 / 2}\left(\mathbf{R}^{3}\right)$ to $\tilde{\Gamma}$ and for $t<0$ by duality: $H^{t}(\tilde{\Gamma})=\left(H^{-t}(\tilde{\Gamma})\right)^{\prime}$. Using these spaces we define the Sobolev spaces on the open surface $\Gamma: \tilde{H}^{t}(\Gamma)=\left\{u \in H^{t}(\tilde{\Gamma}) ; \operatorname{supp} u \subset \bar{\Gamma}\right\}$ and $H^{t}(\Gamma)=\left\{\left.u\right|_{\Gamma} ; u \in H^{t}(\tilde{\Gamma})\right\}$ for any real $t$. We use these notations for scalar functions as well as for vector functions, using the norms and inner products componentwise. In the sequel vector functions will be denoted by bold face symbols.

We consider the Dirichlet boundary value problem for the displacement field $\mathbf{u}=\left(u_{1}, u_{2}, u_{3}\right)$ of a homogeneous, isotropic, elastic material covering the domain $\Omega_{\Gamma}:=\mathbf{R}^{3} \backslash \bar{\Gamma}$ : For given $\mathbf{u}_{1}, \mathbf{u}_{2} \in H^{1 / 2}(\Gamma)$ with $\mathbf{u}_{1}-\mathbf{u}_{2} \in$ $\tilde{H}^{1 / 2}(\Gamma)$ find $\mathbf{u}$ satisfying

$$
\begin{gather*}
\mu \Delta \mathbf{u}+(\lambda+\mu) \operatorname{grad} \operatorname{div} \mathbf{u}=0 \quad \text { in } \Omega_{\Gamma},  \tag{1.1}\\
\left.\mathbf{u}\right|_{\Gamma_{1}}=\mathbf{u}_{1},\left.\mathbf{u}\right|_{\Gamma_{2}}=\mathbf{u}_{2},  \tag{1.2}\\
\mathbf{u}(x)=o(1), \frac{\partial}{\partial x_{j}} \mathbf{u}(x)=o\left(|x|^{-1}\right), j=1,2,3, \quad|x| \rightarrow \infty . \tag{1.3}
\end{gather*}
$$

Here, $\Gamma_{i}, i=1,2$, are the two sides of $\Gamma$ and $\mu>0, \lambda>-2 / 3 \mu$ are the given Lamé constants. The corresponding Neumann data of the linear elasticity problem are the tractions

$$
\mathbf{T}(\mathbf{u}):=\lambda(\operatorname{div} \mathbf{u}) \mathbf{n}+2 \mu \frac{\partial \mathbf{u}}{\partial \mathbf{n}}+\mu \mathbf{n} \times \operatorname{curl} \mathbf{u} \quad \text { on } \Gamma_{i}, i=1,2
$$

where $\mathbf{n}$ is the normal vector exterior to the bounded domain enclosed by $\tilde{\Gamma}$.
The problem (1.1)-(1.3) can be formulated as an integral equation of the first kind, see, e.g., [8, 4]: $\mathbf{u} \in H_{\mathrm{loc}}^{1}\left(\mathbf{R}^{3} \backslash \bar{\Gamma}\right)$ is the solution of the Dirichlet problem (1.1)-(1.3) if and only if the jump of the traction $\mathbf{t}:=$ $\left.\mathbf{T}(\mathbf{u})\right|_{\Gamma_{1}}-\left.\mathbf{T}(\mathbf{u})\right|_{\Gamma_{2}} \in \tilde{H}^{-1 / 2}(\Gamma)$ solves the weakly singular integral equation

$$
\begin{equation*}
\mathbf{V t}(x):=\int_{\Gamma} \mathbf{E}(y, x) \mathbf{t}(y) d s_{y}=\mathbf{g}(x), \quad x \in \Gamma \tag{1.4}
\end{equation*}
$$

where

$$
\mathbf{g}(x)=\frac{1}{2}\left(\mathbf{u}_{1}+\mathbf{u}_{2}\right)(x)+\int_{\Gamma} \mathbf{T}_{y} \mathbf{E}(y, x)\left(\mathbf{u}_{1}-\mathbf{u}_{2}\right)(y) d s_{y}
$$

Here,

$$
\mathbf{E}(y, x)=\frac{\lambda+3 \mu}{8 \pi \mu(\lambda+2 \mu)}\left(\frac{1}{|x-y|} I+\frac{\lambda+\mu}{\lambda+3 \mu} \frac{(x-y)(x-y)^{T}}{|x-y|^{3}}\right)
$$

denotes the fundamental solution of (1.1) with the identity matrix $I$. The solution $\mathbf{t}$ of (1.4) yields the solution to problem (1.1)-(1.3) via the representation or Betti's formula

$$
\mathbf{u}(x)=\int_{\Gamma}\left(\mathbf{E}(y, x) \mathbf{t}(y)-\left(\mathbf{T}_{y} \mathbf{E}(y, x)\right)^{T}\left(\mathbf{u}_{1}(y)-\mathbf{u}_{2}(y)\right)\right) d s_{y}, \quad x \notin \Gamma .
$$

In what follows, together with usual space coordinates $\left(x_{1}, x_{2}, x_{3}\right)=x \in \Gamma$ we will use surface coordinates $(s, \rho)$ in a small neighborhood of $\gamma$ on $\Gamma$ such that $s$ (respectively, $\rho$ ) varies in tangential (respectively, normal) direction to $\gamma$. Thus the boundary curve $\gamma$ is described by the equation $\rho=0$, and in a sufficiently small neighborhood of $\gamma$ one has $s=s(x)$ and $\rho=\rho(x)$. Throughout the paper we will specify this small neighborhood of $\gamma$ as the boundary strip $\Gamma_{\delta}$ of $\Gamma$ such that for small $\delta>0$,

$$
\Gamma_{\delta}=\{x \in \Gamma ; 0<\rho(x)<\delta\}
$$

Let us cite the following regularity result from [4].
Proposition 1.1 Let $|\sigma|<1 / 2$ and $\mathbf{u}_{j} \in H^{3 / 2+\sigma}(\Gamma), j=1,2$, with $\mathbf{u}_{1}-\mathbf{u}_{2} \in \tilde{H}^{3 / 2+\sigma}(\Gamma)$. Then the solution $\mathbf{t} \in \tilde{H}^{-1 / 2}(\Gamma)$ of the integral equation (1.4) has the form

$$
\begin{equation*}
\mathbf{t}=\boldsymbol{\beta}(s) \rho^{-1 / 2} \chi(\rho)+\mathbf{t}_{0} \tag{1.5}
\end{equation*}
$$

with vector functions $\boldsymbol{\beta} \in H^{1 / 2+\sigma}(\gamma)$ and $\mathbf{t}_{0} \in \tilde{H}^{1 / 2+\sigma^{\prime}}(\Gamma)$ for any $\sigma^{\prime}<\sigma$. Furthermore, $\chi \in C_{0}^{\infty}(\mathbf{R})$ denotes a cut-off function with $0 \leq \chi \leq 1$ and $\chi=1$ near zero.

In the next section we formulate the $p$-version of the BEM for the approximate solution of (1.4) and state the main result which proves an almost optimal convergence rate (Theorem 2.1). Technical details and the proof of Theorem 2.1 are given in Section 3.

## 2 The $p$-version of the BEM

Below $p$ will always denote a polynomial degree, and $C$ is a generic positive constant independent of $p$.

In order to define finite dimensional subspaces of $\tilde{H}^{-1 / 2}(\Gamma)$ we use a regular parameter representation $x=X(u), u \in U, U$ being a compact region in $\mathbf{R}^{2}$ whose boundary is mapped onto $\gamma$. On $U$ we use a fixed regular mesh $\mathcal{T}=\left\{U_{j} ; j=1, \ldots, J\right\}$ of quadrilaterals and triangles which are in general curvilinear such that $U$ is completely discretized. We assume that for each $j=1, \ldots, J$ there exists a smooth one-to-one mapping $M_{j}$ such that $\bar{U}_{j}=M_{j}(\bar{K})$ with $K=Q$ or $T$ (here, $Q=(-1,1)^{2}$ and $T=$ $\left\{\xi=\left(\xi_{1}, \xi_{2}\right) ; 0<\xi_{1}<1,0<\xi_{2}<\xi_{1}\right\}$ denote the reference square and triangle, respectively). The Jacobians of $M_{j}$ are assumed to be bounded above and below by positive constants independent of $j$.

Using the parameter representation $X$ we have a fixed regular mesh $\Delta=\left\{\Gamma_{j}=X\left(U_{j}\right) ; j=1, \ldots, J\right\}$ on $\Gamma$. The union of the elements of $\Delta$ touching the boundary curve $\gamma$ will be denoted by $A_{\gamma}$, i.e., $\bar{A}_{\gamma}=\cup\left\{\bar{\Gamma}_{j} ; \bar{\Gamma}_{j} \cap\right.$ $\gamma \neq \varnothing\}$. We assume that, close to the $\gamma$, the mesh is fine enough such that $\bar{A}_{\gamma} \subset\left(\Gamma_{\delta / 2} \cup \gamma\right)$. We also assume that the cut-off function $\chi$ in (1.5) is chosen such that $\operatorname{supp}\left(\boldsymbol{\beta}(s) \rho^{-1 / 2} \chi(\rho)\right) \subset \bar{A}_{\gamma}$.

Now for given integer $p$ we define the space $S_{p}(\Gamma)$ of piecewise polynomials on $\Gamma$. For $K=Q$ or $K=T$ let $\mathcal{Q}_{p}(K)$ be the set of polynomials of degree $\leq p$ (in each variable for $K=Q$ and of total degree $\leq p$ on $T$ ). Furthermore, for $K=I$ an interval, $\mathcal{Q}_{p}(I)$ denotes the set of polynomials of degree $\leq p$ on $I$. We will also use the set $\mathcal{R}_{p}\left(\Gamma_{j}\right)$ of polynomials of degree $\leq p$ in each variable $s$ and $\rho$ on the elements $\Gamma_{j} \subset A_{\gamma} \subset \Gamma_{\delta / 2}$. Then using the notation $\mathbf{v}_{j}=\left.\mathbf{v}\right|_{\Gamma_{j}}$ we define

$$
\begin{aligned}
& S_{p}(\Gamma):=\left\{\mathbf{v} ; \mathbf{v}_{j} \in\left[\mathcal{R}_{p}\left(\Gamma_{j}\right)\right]^{3} \text { if } \Gamma_{j} \subset A_{\gamma},\right. \text { and } \\
& \left.\qquad\left(\mathbf{v}_{j} \circ X \circ M_{j}\right) \in\left[\mathcal{Q}_{p}(K)\right]^{3}, K=Q \text { or } T, \text { if } \Gamma_{j} \subset\left(\Gamma \backslash A_{\gamma}\right)\right\}
\end{aligned}
$$

(here, we denote by $[\cdot]^{3}$ the sets of vector functions with corresponding polynomial components).

One has $S_{p}(\Gamma) \subset \tilde{H}^{-1 / 2}(\Gamma)$, and the $p$-version of the boundary element Galerkin method is as follows: For given $p$ find $\mathbf{t}_{p} \in S_{p}(\Gamma)$ such that

$$
\begin{equation*}
\left\langle\mathbf{V} \mathbf{t}_{p}, \mathbf{v}\right\rangle=\langle\mathbf{g}, \mathbf{v}\rangle \quad \forall \mathbf{v} \in S_{p}(\Gamma) . \tag{2.1}
\end{equation*}
$$

As it is well known, this method converges quasi-optimally, see [3], i.e., there exists a constant $C>0$ such that for all polynomial degrees $p$ there holds

$$
\begin{equation*}
\left\|\mathbf{t}-\mathbf{t}_{p}\right\|_{\tilde{H}^{-1 / 2}(\Gamma)} \leq C \inf \left\{\|\mathbf{t}-\mathbf{v}\|_{\tilde{H}^{-1 / 2}(\Gamma)} ; \mathbf{v} \in S_{p}(\Gamma)\right\} . \tag{2.2}
\end{equation*}
$$

We now present the main result giving an a priori error estimate.

Theorem 2.1 Let $|\sigma|<1 / 2$ and $\mathbf{u}_{j} \in H^{3 / 2+\sigma}(\Gamma), j=1,2$, with $\mathbf{u}_{1}-\mathbf{u}_{2} \in$ $\tilde{H}^{3 / 2+\sigma}(\Gamma)$. Then there holds the a priori error estimate

$$
\begin{equation*}
\left\|\mathbf{t}-\mathbf{t}_{p}\right\|_{\tilde{H}^{-1 / 2}(\Gamma)} \leq C p^{-\alpha}, \quad \alpha=1 / 2+\sigma-\varepsilon, \varepsilon>0 \tag{2.3}
\end{equation*}
$$

where $C>0$ depends on $\varepsilon$ but not on $p$. Here, $\mathbf{t}$ is the solution of (1.4) and $\mathbf{t}_{p}$ is the boundary element approximation to $\mathbf{t}$ given by (2.1).

This error estimate is quasi-optimal for sufficiently smooth given data. More precisely, if $\sigma$ is large enough then there exists for any $\varepsilon>0$ a constant $c>0$ such that the $p$-version converges like $c p^{-1+\varepsilon}$. A convergence like $c p^{-1}$ would be optimal, cf. the results in [7, 2]. The sub-optimality of (2.3) is due to Proposition 1.1 which states the regularity of the term $\boldsymbol{\beta}$ in the representation of the exact solution only in standard Sobolev spaces, which are not appropriate to obtain optimal results. For numerical results (dealing with the scalar version of the Laplace operator) which underline the a priori error estimate we refer to [6].

The proof of Theorem 2.1 is given in the next section.

## 3 Technical details

Before proving Theorem 2.1 we collect several auxiliary results.
Lemma 3.1 Let $\Omega \subset \mathbf{R}^{2}$ be a Lipschitz domain. If $u \in \tilde{H}^{t}(\Omega)$ with $0 \leq t \leq 1$, then for $i=1,2, \partial u / \partial x_{i} \in \tilde{H}^{t-1}(\Omega)$, and

$$
\left\|\partial u / \partial x_{i}\right\|_{\tilde{H}^{t-1}(\Omega)} \leq C\|u\|_{\tilde{H}^{t}(\Omega)},
$$

where $C>0$ is independent of $u$.
On an interval, this statement is proved in [9, Lemma 3.5]. In two dimensions the proof is similar and is skipped.

Lemma 3.2 Let $\Omega$, $\Omega_{1}$ be two Lipschitz domains in $\mathbf{R}^{n}(n=1,2,3)$, and $\Omega_{1} \subset \Omega$. Then, for $0 \leq t<1 / 2$, there holds

$$
\begin{equation*}
\|u\|_{\tilde{H}^{-t}\left(\Omega_{1}\right)} \leq C\|u\|_{\tilde{H}^{-t}(\Omega)} \quad \forall u \in \tilde{H}^{-t}(\Omega), \tag{3.1}
\end{equation*}
$$

where the constant $C>0$ is independent of $u$.
Proof. For $0 \leq t<1 / 2$, the identity $H_{0}^{t}\left(\Omega_{1}\right)=H^{t}\left(\Omega_{1}\right)$ holds (see, e.g., [5]). Let us consider the function $v \in H^{t}\left(\Omega_{1}\right)=H_{0}^{t}\left(\Omega_{1}\right)$ and denote by $\bar{v}$ the extension of $v$ by zero outside $\Omega_{1}$. Then $\bar{v} \in H^{t}(\Omega)=H_{0}^{t}(\Omega)$,

$$
\|\bar{v}\|_{H^{t}(\Omega)} \leq C\left(\|\bar{v}\|_{H^{t}\left(\Omega_{1}\right)}+\|\bar{v}\|_{H^{t}\left(\Omega \backslash \Omega_{1}\right)}\right)=C\|v\|_{H^{t}\left(\Omega_{1}\right)}
$$

and (3.1) follows from the definition of the norm in $\tilde{H}^{-t}\left(\Omega_{1}\right)$.

Lemma 3.3 Let $f \in H^{t}(K)$ for real $t>0$ with $K=I \subset \mathbf{R}$ (respectively, $K=Q$ or $K=T$ in $\left.\mathbf{R}^{2}\right)$. Then there exists a sequence $f_{p} \in \mathcal{Q}_{p}(K)$, $p=0,1,2, \ldots$, such that

$$
\left\|f-f_{p}\right\|_{L_{2}(K)} \leq C p^{-t}\|f\|_{H^{t}(K)}
$$

For a proof of Lemma 3.3 we refer to [1].
Lemma 3.4 [10, Lemma 3.3] Let $f(x) \in \tilde{H}^{-t_{1}}\left(I_{1}\right)$ and $g(y) \in \tilde{H}^{-t_{2}}\left(I_{2}\right)$ with $0 \leq t_{1}, t_{2} \leq 1$. Then $f(x) g(y) \in \tilde{H}^{-t_{1}-t_{2}}\left(I_{1} \times I_{2}\right)$ and

$$
\|f(x) g(y)\|_{\tilde{H}^{-t_{1}-t_{2}\left(I_{1} \times I_{2}\right)}} \leq c\|f(x)\|_{\tilde{H}^{-t_{1}}\left(I_{1}\right)}\|g(y)\|_{\tilde{H}^{-t_{2}\left(I_{2}\right)}}
$$

The constant $c$ is independent of $f$ and $g$.
To analyze the approximation of the singular part of $\mathbf{t}$ in (1.5) we first study singularities on an interval. Let us consider the singular function

$$
\begin{equation*}
\psi(x)=(1+x)^{\lambda-1} \chi(x), \quad x \in I=(-1,1), \tag{3.2}
\end{equation*}
$$

where $\lambda>0$ is real, $\chi \in C^{\infty}(I)$ is a cut-off function with $\chi(x)=1$ for $x \in(-1,-1+d]$ and $\chi(x)=0$ for $x \geq-1+2 d(0<d \leq 1 / 4)$.

Observe that $\psi \in \tilde{H}^{t}(I)$ for $-1 \leq t<\min \{0, \lambda-1 / 2\}$.
Theorem 3.1 Let $\psi(x)$ be given by (3.2) with $\lambda>0$. Then there exists a sequence $\psi_{p} \in \mathcal{Q}_{p}(I), p=1,2, \ldots$, such that for $-1 \leq t<\min \{0, \lambda-1 / 2\}$,

$$
\begin{equation*}
\left\|\psi-\psi_{p}\right\|_{\tilde{H}^{t}(\tilde{I})} \leq C p^{-2(\lambda-1 / 2-t)}, \quad \tilde{I}=(-1,0) \tag{3.3}
\end{equation*}
$$

Proof. Introducing a $C^{\infty}$ cut-off function $\tilde{\chi}(x)$ such that

$$
\begin{equation*}
\tilde{\chi}(x)=1 \text { for } x \in[-1,0] \text { and } \tilde{\chi}(x)=0 \text { for } x \geq 1 / 2, \tag{3.4}
\end{equation*}
$$

we define

$$
\Psi(x):=\tilde{\chi}(x) \int_{-1}^{x} \psi(\xi) d \xi, \quad \hat{\Psi}(x):=(1-x)^{-1} \Psi(x), \quad x \in I=(-1,1) .
$$

Then $\Psi(-1)=\hat{\Psi}(-1)=0, \Psi(x)=\hat{\Psi}(x)=0$ for $x \in[1 / 2,1]$, and

$$
\begin{equation*}
\Psi^{\prime}(x)=\psi(x) \text { for } x \in \tilde{I}=(-1,0) \tag{3.5}
\end{equation*}
$$

Further, using integration by parts we obtain

$$
\begin{equation*}
\hat{\Psi}(x)=\frac{(1+x)^{\lambda} \chi(x) \tilde{\chi}(x)}{\lambda(1-x)}-\frac{\tilde{\chi}(x)}{\lambda(1-x)} \int_{-1}^{x}(1+\xi)^{\lambda} \chi^{\prime}(\xi) d \xi=: F(x)-G(x) \tag{3.6}
\end{equation*}
$$

Referring to [2, Theorem 3.1] if $0<\lambda \leq 1 / 2$ and to [7, Theorem 5.1] if $\lambda>1 / 2$, we find a polynomial $F_{p} \in \mathcal{Q}_{p}(I)$ such that $F_{p}(-1)=F(-1)=0$ and

$$
\begin{equation*}
\left\|F-F_{p}\right\|_{H^{t}(I)} \leq C p^{-2(\lambda+1 / 2-t)}, \quad 0 \leq t<\min \{1, \lambda+1 / 2\} \tag{3.7}
\end{equation*}
$$

For the function $G \in C_{0}^{\infty}(I)$ there exists by Lemma 3.3 a polynomial $G_{p} \in \mathcal{Q}_{p}(I)$ such that $G_{p}( \pm 1)=G( \pm 1)=0$, and for arbitrary $\tau>0$,

$$
\begin{equation*}
\left\|G-G_{p}\right\|_{H^{t}(I)} \leq C p^{-\tau}, \quad 0 \leq t \leq 1 \tag{3.8}
\end{equation*}
$$

Let us define $\Psi_{p}(x):=(1-x)\left(F_{p}(x)-G_{p}(x)\right)$. Then $\Psi_{p} \in \mathcal{Q}_{p+1}(I)$, $\Psi_{p}( \pm 1)=0$, and for $0 \leq t<\min \{1, \lambda+1 / 2\}$ we deduce from (3.6)-(3.8)

$$
\begin{equation*}
\left\|\Psi-\Psi_{p}\right\|_{H^{t}(I)} \leq C\left\|\hat{\Psi}-\left(F_{p}-G_{p}\right)\right\|_{H^{t}(I)} \leq C p^{-2(\lambda+1 / 2-t)} \tag{3.9}
\end{equation*}
$$

Hence

$$
\begin{equation*}
\left\|\Psi-\Psi_{p}\right\|_{\tilde{H}^{t}(I)} \leq C p^{-2(\lambda+1 / 2-t)}, \quad t \in[0, \min \{1, \lambda+1 / 2\}) \backslash\{1 / 2\} \tag{3.10}
\end{equation*}
$$

because $\left(\Psi-\Psi_{p}\right) \in H_{0}^{t}(I)=\tilde{H}^{t}(I)$ for these values of $t$.
Now we set $\psi_{p}(x):=\Psi_{p}^{\prime}(x)$ for $x \in I$. Then $\psi_{p} \in \mathcal{Q}_{p}(I)$, and recalling (3.5) we have $\psi-\psi_{p}=\left(\Psi-\Psi_{p}\right)^{\prime}$ on $\tilde{I}$. Therefore, using sequentially the one-dimensional versions of Lemmas 3.2, 3.1, and then estimate (3.10) we obtain for any fixed $t^{\prime} \in(1 / 2, \min \{1, \lambda+1 / 2\})$

$$
\begin{align*}
\left\|\psi-\psi_{p}\right\|_{\tilde{H}^{t^{\prime}-1}(\tilde{I})} & =\left\|\left(\Psi-\Psi_{p}\right)^{\prime}\right\|_{\tilde{H}^{t^{\prime}-1}(\tilde{I})} \leq C\left\|\left(\Psi-\Psi_{p}\right)^{\prime}\right\|_{\tilde{H}^{t^{\prime}-1}(I)} \\
& \leq C\left\|\Psi-\Psi_{p}\right\|_{\tilde{H}^{t^{\prime}}(I)} \leq C p^{-2\left(\lambda+1 / 2-t^{\prime}\right)} \tag{3.11}
\end{align*}
$$

Thus we have proved (3.3) for $t \in(-1 / 2, \min \{0, \lambda-1 / 2\})$.
On the other hand, applying Lemma 3.1 and inequality (3.9) we have

$$
\left\|\psi-\psi_{p}\right\|_{\tilde{H}^{-1}(\tilde{I})}=\left\|\left(\Psi-\Psi_{p}\right)^{\prime}\right\|_{\tilde{H}^{-1}(\tilde{I})} \leq C\left\|\Psi-\Psi_{p}\right\|_{H^{0}(\tilde{I})} \leq C p^{-2(\lambda+1 / 2)}
$$

Since $-1 / 2<t^{\prime}-1<\min \{0, \lambda-1 / 2\}$ in (3.11), the interpolation between $\tilde{H}^{-1}(\tilde{I})$ and $\tilde{H}^{t^{\prime}-1}(\tilde{I})$ gives (3.3) for any $t \in[-1,-1 / 2]$.

Remark 3.1 When proving Theorem 3.1 we have also established the following inequality (see (3.9))

$$
\begin{equation*}
\left\|\Psi-\Psi_{p}\right\|_{L_{2}(I)} \leq C p^{-2(\lambda+1 / 2)} \tag{3.12}
\end{equation*}
$$

where $\Psi(x)=\tilde{\chi}(x) \int_{-1}^{x} \psi(\xi) d \xi, \Psi_{p}(x)=\int_{-1}^{x} \psi_{p}(\xi) d \xi$, the function $\psi(x)$ is given by (3.2), and $\psi_{p}(x)$ is a polynomial approximation to $\psi(x)$.

Moreover, $\Psi(x) \in L_{2}(I)$, and (3.12) yields

$$
\begin{equation*}
\left\|\Psi_{p}\right\|_{L_{2}(I)} \leq C \tag{3.13}
\end{equation*}
$$

Now we prove the main result of the paper.
Proof of Theorem 2.1. Due to the regularity result of Proposition 1.1 and the quasi-optimal convergence (2.2) of the BEM, one only needs to find a piecewise polynomial function that approximates $\mathbf{t}$ in (1.5) with the upper bound stated by (2.3).

For elements at the boundary $\gamma$ we need covering rectangles in surface coordinates. Let $\Gamma_{j} \subset A_{\gamma}$ be an element touching the boundary $\gamma$. Since $A_{\gamma} \subset\left(\Gamma_{\delta / 2} \cup \gamma\right)$, there exist two points on $\gamma$ with coordinates $\left(s_{1}, 0\right)$ and $\left(s_{2}, 0\right)$ such that

$$
\Gamma_{j} \subset Q_{j}=\left\{(s, \rho) \in \Gamma_{\delta / 2} ; s_{1}<s<s_{2}, 0<\rho<\delta / 2\right\}
$$

First, we define an approximation $\mathbf{t}_{0, p}$ to the vector function $\mathbf{t}_{0} \in \tilde{H}^{\alpha}(\Gamma) \subset$ $H^{\alpha}(\Gamma)$ (hereafter, $\alpha=1 / 2+\sigma-\varepsilon>0$ with sufficiently small $\varepsilon>0$ ). If $\Gamma_{j} \subset$ ( $\Gamma \backslash A_{\gamma}$ ), we apply Lemma 3.3 componentwise on the square (or triangle) $K$ such that $\Gamma_{j}=X\left(M_{j}(K)\right)$. However, if $\Gamma_{j} \subset A_{\gamma}$, we apply Lemma 3.3 on $Q_{j} \supset \Gamma_{j}$. Since $\Gamma$ is smooth, the function $\mathbf{t}_{0}$ on $\Gamma_{\delta} \supset A_{\gamma}$ has the same regularity in terms of coordinates $(s, \rho)$ as in terms of space variables $x=$ $X(u)$. Therefore, recalling the definition of $S_{p}(\Gamma)$ and applying Lemma 3.3 as indicated above, we find $\mathbf{t}_{0, p} \in S_{p}(\Gamma)$ such that

$$
\begin{equation*}
\left\|\mathbf{t}_{0}-\mathbf{t}_{0, p}\right\|_{\tilde{H}^{-1 / 2}\left(\Gamma_{j}\right)} \leq\left\|\mathbf{t}_{0}-\mathbf{t}_{0, p}\right\|_{L_{2}\left(\Gamma_{j}\right)} \leq C p^{-\alpha}\left\|\mathbf{t}_{0}\right\|_{H^{\alpha}\left(\Gamma_{j}\right)} \leq C p^{-\alpha} \tag{3.14}
\end{equation*}
$$

if $\Gamma_{j} \subset\left(\Gamma \backslash A_{\gamma}\right)$, and

$$
\begin{align*}
\left\|\mathbf{t}_{0}-\mathbf{t}_{0, p}\right\|_{\tilde{H}^{-1 / 2}\left(\Gamma_{j}\right)} & \leq\left\|\mathbf{t}_{0}-\mathbf{t}_{0, p}\right\|_{L_{2}\left(\Gamma_{j}\right)} \leq\left\|\mathbf{t}_{0}-\mathbf{t}_{0, p}\right\|_{L_{2}\left(Q_{j}\right)} \\
& \leq C p^{-\alpha}\left\|\mathbf{t}_{0}\right\|_{H^{\alpha}\left(Q_{j}\right)} \leq C p^{-\alpha} \tag{3.15}
\end{align*}
$$

if $\Gamma_{j} \subset A_{\gamma}$.
Now we consider the singular term $\boldsymbol{\beta}(s) \psi(\rho)=\boldsymbol{\beta}(s) \rho^{-1 / 2} \chi(\rho)$ in (1.5). Let $\Gamma_{j} \subset A_{\gamma}$, and $\Gamma_{j} \subset Q_{j}$ as above. Then using the one-dimensional version of Lemma 3.3 we approximate the function $\boldsymbol{\beta}(s) \in H^{1 / 2+\sigma}(\gamma)$ : there exists $\boldsymbol{\beta}_{p}(s) \in\left[\mathcal{Q}_{p}\left(s_{1}, s_{2}\right)\right]^{3}$ satisfying

$$
\begin{equation*}
\left\|\boldsymbol{\beta}-\boldsymbol{\beta}_{p}\right\|_{L_{2}\left(s_{1}, s_{2}\right)} \leq C p^{-(1 / 2+\sigma)}\|\boldsymbol{\beta}\|_{H^{1 / 2+\sigma}\left(s_{1}, s_{2}\right)} \leq C p^{-(1 / 2+\sigma)}\|\boldsymbol{\beta}\|_{H^{1 / 2+\sigma}(\gamma)} . \tag{3.16}
\end{equation*}
$$

For the singular function $\psi(\rho)$ we apply Theorem 3.1, scaled to the interval $(0, \delta)$, with $\lambda=1 / 2$ : there exists a polynomial $\psi_{p}(\rho) \in \mathcal{Q}_{p}(0, \delta)$ satisfying

$$
\begin{equation*}
\left\|\psi-\psi_{p}\right\|_{\tilde{H}^{-t}(0, \delta / 2)} \leq C p^{-2 t}, \quad 0<t \leq 1 . \tag{3.17}
\end{equation*}
$$

Since $\psi(\rho) \in \tilde{H}^{-t}(0, \delta / 2)$ with $t \in(0,1]$, we estimate by (3.17)

$$
\begin{equation*}
\left\|\psi_{p}\right\|_{\tilde{H}^{-t}(0, \delta / 2)} \leq C, \quad 0<t \leq 1 . \tag{3.18}
\end{equation*}
$$

Furthermore, introducing a $C^{\infty}$ cut-off function $\tilde{\chi}(\rho)$ such that (cf. (3.4))

$$
\tilde{\chi}(\rho)=1 \text { for } \rho \in[0, \delta / 2] \text { and } \tilde{\chi}(\rho)=0 \text { for } \rho \geq 3 \delta / 4
$$

and arguing as in the proof of Theorem 3.1 we obtain (cf. (3.12), (3.13))

$$
\begin{equation*}
\left\|\Psi-\Psi_{p}\right\|_{L_{2}(0, \delta)} \leq C p^{-2}, \quad\left\|\Psi_{p}\right\|_{L_{2}(0, \delta)} \leq C \tag{3.19}
\end{equation*}
$$

where $\Psi(\rho)=\tilde{\chi}(\rho) \int_{0}^{\rho} \psi(r) d r$ and $\Psi_{p}(\rho)=\int_{0}^{\rho} \psi_{p}(r) d r$.
Then making use of Lemma 3.2 (which remains valid with $\Omega_{1}=\Gamma_{j} \subset$ $\left.Q_{j}=\Omega\right)$, Lemma 3.4, the triangle inequality, and estimates (3.16)-(3.18), we derive for some fixed $t^{\prime} \in(0,1 / 2)$

$$
\begin{align*}
\| \boldsymbol{\beta} \psi & -\boldsymbol{\beta}_{p} \psi_{p} \|_{\tilde{H}^{-t^{\prime}}\left(\Gamma_{j}\right)} \\
& \leq C\left(\left\|\boldsymbol{\beta}\left(\psi-\psi_{p}\right)\right\|_{\tilde{H}-t^{\prime}\left(Q_{j}\right)}+\left\|\left(\boldsymbol{\beta}-\boldsymbol{\beta}_{p}\right) \psi_{p}\right\|_{\tilde{H}^{-t^{\prime}}\left(Q_{j}\right)}\right) \\
& \leq C\left(\|\boldsymbol{\beta}\|_{L_{2}\left(s_{1}, s_{2}\right)}\left\|\psi-\psi_{p}\right\|_{\tilde{H}^{-t^{\prime}}(0, \delta / 2)}+\left\|\boldsymbol{\beta}-\boldsymbol{\beta}_{p}\right\|_{L_{2}\left(s_{1}, s_{2}\right)}\left\|\psi_{p}\right\|_{\tilde{H}^{-t^{\prime}}(0, \delta / 2)}\right) \\
& \leq C p^{-\min \left\{1 / 2+\sigma, 2 t^{\prime}\right\}}\|\boldsymbol{\beta}\|_{H^{1 / 2+\sigma}(\gamma)} . \tag{3.20}
\end{align*}
$$

On the other hand, using the above notation for $\Psi(\rho)$ and $\Psi_{p}(\rho)$ we write

$$
\left\|\boldsymbol{\beta} \psi-\boldsymbol{\beta}_{p} \psi_{p}\right\|_{\tilde{H}^{-1}\left(\Gamma_{j}\right)}=\left\|\frac{\partial}{\partial \rho}\left(\boldsymbol{\beta}(s) \Psi(\rho)-\boldsymbol{\beta}_{p}(s) \Psi_{p}(\rho)\right)\right\|_{\tilde{H}^{-1}\left(\Gamma_{j}\right)}
$$

Then applying Lemma 3.1 in terms of coordinates $(s, \rho) \in \Gamma_{j}$ we have

$$
\begin{aligned}
\left\|\boldsymbol{\beta} \psi-\boldsymbol{\beta}_{p} \psi_{p}\right\|_{\tilde{H}^{-1}\left(\Gamma_{j}\right)} & \leq C\left\|\boldsymbol{\beta}(s) \Psi(\rho)-\boldsymbol{\beta}_{p}(s) \Psi_{p}(\rho)\right\|_{H^{0}\left(\Gamma_{j}\right)} \\
& \leq C\left\|\boldsymbol{\beta}(s) \Psi(\rho)-\boldsymbol{\beta}_{p}(s) \Psi_{p}(\rho)\right\|_{H^{0}\left(Q_{j}\right)} .
\end{aligned}
$$

Hence

$$
\begin{aligned}
&\left\|\boldsymbol{\beta} \psi-\boldsymbol{\beta}_{p} \psi_{p}\right\|_{\tilde{H}^{-1}\left(\Gamma_{j}\right)} \leq C\left(\|\boldsymbol{\beta}\|_{L_{2}\left(s_{1}, s_{2}\right)}\left\|\Psi-\Psi_{p}\right\|_{L_{2}(0, \delta / 2)}\right. \\
&\left.+\left\|\boldsymbol{\beta}-\boldsymbol{\beta}_{p}\right\|_{L_{2}\left(s_{1}, s_{2}\right)}\left\|\Psi_{p}\right\|_{L_{2}(0, \delta / 2)}\right)
\end{aligned}
$$

and we estimate by using (3.16), (3.19)

$$
\left\|\boldsymbol{\beta} \psi-\boldsymbol{\beta}_{p} \psi_{p}\right\|_{\tilde{H}^{-1}\left(\Gamma_{j}\right)} \leq C p^{-\min \{1 / 2+\sigma, 2\}}\|\boldsymbol{\beta}\|_{H^{1 / 2+\sigma}(\gamma)}
$$

Since $|\sigma|<1 / 2$, we may take $t^{\prime}$ in (3.20) such that $0<1 / 2+\sigma \leq 2 t^{\prime}<1$. Then interpolating between $\tilde{H}^{-1}\left(\Gamma_{j}\right)$ and $\tilde{H}^{-t^{\prime}}\left(\Gamma_{j}\right)$ we prove for any $\Gamma_{j} \subset A_{\gamma}$

$$
\begin{equation*}
\left\|\boldsymbol{\beta} \psi-\boldsymbol{\beta}_{p} \psi_{p}\right\|_{\tilde{H}^{-1 / 2}\left(\Gamma_{j}\right)} \leq C p^{-(1 / 2+\sigma)}\|\boldsymbol{\beta}\|_{H^{1 / 2+\sigma}(\gamma)} \tag{3.21}
\end{equation*}
$$

Now let us define the approximating function $\mathbf{v}_{p}$ on $\Gamma$ as follows: $\left.\mathbf{v}_{p}\right|_{\Gamma_{j}}=$ $\boldsymbol{\beta}_{p} \psi_{p}+\left.\mathbf{t}_{0, p}\right|_{\Gamma_{j}}$ if $\Gamma_{j} \subset A_{\gamma}$, and $\left.\mathbf{v}_{p}\right|_{\Gamma_{j}}=\left.\mathbf{t}_{0, p}\right|_{\Gamma_{j}}$ if $\Gamma_{j} \subset\left(\Gamma \backslash A_{\gamma}\right)$. Then $\mathbf{v}_{p} \in S_{p}(\Gamma)$, and, due to (3.14), (3.15), (3.21), the error $\left\|\mathbf{t}-\mathbf{v}_{p}\right\|_{\tilde{H}^{-1 / 2}(\Gamma)}$ satisfies the upper bound in (2.3). This proves the theorem.

## References

[1] Babuška, I., and Suri, M. (1987). The optimal convergence rate of the $p$-version of the finite element method. SIAM J. Numer. Anal. 24, 750-776.
[2] Bespalov, A., and Heuer, N. (2003). The $p$-version of the boundary element method for hypersingular operators on piecewise plane open surfaces. Preprint 03-20, Departamento de Ingeniería Matemática, Universidad de Concepción, Chile.
[3] Costabel, M. (1988). Boundary integral operators on Lipschitz domains: Elementary results. SIAM J. Math. Anal. 19, 613-626.
[4] Costabel, M., and Stephan, E. P. (1987). An improved boundary element Galerkin method for three-dimensional crack problems. Integral Equations Operator Theory 10, 467-504.
[5] Grisvard, P. (1985). Elliptic Problems in Nonsmooth Domains, Pitman Publishing Inc., Boston.
[6] Heuer, N., Maischak, M., and Stephan, E. P. (1999). Exponential convergence of the hp-Version for the Boundary Element Method on open surfaces. Numer. Math. 83, 641-666.
[7] Schwab, C., and Suri, M. (1996). The optimal p-version approximation of singularities on polyhedra in the boundary element method. SIAM J. Numer. Anal. 33, 729-759.
[8] Stephan, E. P. (1986). A boundary integral equation method for threedimensional crack problems in elasticity. Math. Methods Appl. Sci. 8, 609-623.
[9] Stephan, E. P., and Suri, M. (1991). The $h-p$ version of the boundary element method on polygonal domains with quasiuniform meshes. RAIRO Modél. Math. Anal. Numér. 25, 783-807.
[10] von Petersdorff, T. (1989). Randwertprobleme der Elastizitätstheorie für Polyeder - Singularitäten und Approximation mit Randelementmethoden. PhD thesis, Technische Hochschule Darmstadt, Germany.


[^0]:    *Computational Center, Far-Eastern Branch of the Russian Academy of Sciences, Khabarovsk, Russia. email: albespalov@yahoo.com. Financed by the FONDAP Program in Applied Mathematics, Chile.
    ${ }^{\dagger}$ BICOM, Department of Mathematical Sciences, Brunel University, Uxbridge, Middlesex UB8 3PH, U.K. email: norbert.heuer@brunel.ac.uk Supported in part by the FONDAP Program in Applied Mathematics and Fondecyt project no. 1040615, both Chile.

