

# Funciones de Green.

X superficie de Riemann compacta no-singular.

$C^{\omega}$  (resp  $C^{\infty}$ ) (1,1)-Forma  $\psi$  en X es localmente.

$$i f(z, \bar{z}) dz \wedge d\bar{z} \quad \text{cl } f \in C^{\omega} \text{ (resp } C^{\infty}\text{)}.$$

Notan que  $\begin{cases} dz = dx + i dy \\ d\bar{z} = dx - i dy \end{cases} \Rightarrow i dz \wedge d\bar{z} = 2 dx \wedge dy.$

$\psi$  es forma de volumen si  $f(z, \bar{z}) > 0 \quad \forall (z, \bar{z})$ .

$\psi$  es normalizada si  $\int_X \psi = 1$ .

X viene cl una forma canónica de volumen.  $\mu_X$  o  $\mu$

Una 1-forma holomorfa en X

es localmente  $\int h(z) dz$  cl  $\int h$  holomorfa

integral

conservando  $\langle \psi, \psi \rangle := \frac{i}{2} \int_X \psi \wedge \bar{\psi}$  producto Hermitiano

$\psi_1, \dots, \psi_g$  base ortonormal.

$$\Rightarrow \mu = \frac{i}{2g} \sum_{i=1}^g \psi_i \wedge \bar{\psi}_i \quad \text{es } \underline{\text{forma de volumen.}}$$

normalizada  $\wedge$  es la canónica.

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Si  $g=0$ , setona como forma de volumen

$$\varphi(z) = \frac{i}{2\pi} \frac{dz \wedge d\bar{z}}{(1+|z|^2)^2} \quad z \in \mathbb{C}$$

y corresponde a volumen esférico bajo proyección estereográfica

Si  $g=1$  por  $X = \mathbb{C}/\langle 1, \tau \rangle$  para cierto  $\tau = u + i v$   
 $u, v > 0$

$$\varphi(z) = \frac{i}{2v} dz \wedge d\bar{z} = \frac{1}{v} dx \wedge dy.$$

( donde  $dz = \pi_* dz$   $\pi: \mathbb{C} \rightarrow \mathbb{C}/\langle 1, \tau \rangle$  )

$$d\bar{z} = \pi_* d\bar{z}$$

Cálculo em S.R.

$$dz = dx + i dy$$

$$d\bar{z} = dx - i dy$$

$$\frac{\partial}{\partial \bar{z}} = \frac{1}{2} \left( \frac{\partial}{\partial x} - i \frac{\partial}{\partial y} \right)$$

$$\frac{\partial}{\partial z} = \frac{1}{2} \left( \frac{\partial}{\partial x} + i \frac{\partial}{\partial y} \right)$$

$$f \in X \rightarrow \mathbb{C} \quad \sigma \quad \omega = u \cdot dz + v \cdot d\bar{z}$$

$$\partial f = \frac{\partial f}{\partial z} dz$$

$$\bar{\partial} f = \frac{\partial f}{\partial \bar{z}} d\bar{z}$$

$$\partial \omega = \frac{\partial v}{\partial z} dz \wedge d\bar{z}$$

$$\bar{\partial} \omega = -\frac{\partial u}{\partial \bar{z}} dz \wedge d\bar{z}$$

$$f: X \rightarrow \mathbb{C} \quad df = \frac{\partial f}{\partial z} dz + \frac{\partial f}{\partial \bar{z}} d\bar{z} = (\partial + \bar{\partial}) f$$

$$\partial \bar{\partial} f = \left( \frac{\partial f}{\partial z} dz - \frac{\partial f}{\partial \bar{z}} d\bar{z} \right) \frac{i}{4\pi}$$

$$= \frac{i(\partial - \bar{\partial}) f}{4\pi}$$

}  $df \wedge d\bar{f}$   
}  $\frac{|\nabla f|^2}{4\pi} dx \wedge dy$

~~dd<sup>c</sup>f~~

$$dd^c f = \frac{i}{2\pi} \partial \bar{\partial} f = \frac{1}{4\pi} \Delta f dx \wedge dy$$

Función de Green para un divisor  $D$  c/a un volumen  $\psi$

$g_D : X \setminus \text{supp } D \rightarrow \mathbb{R}$  se llama  $\rightarrow$

si  $\textcircled{1} \forall (U, f)$  representando a  $D$  ie  $\sum [P] n_P U = D$ .

$$g_D(P) = -\log |F(P)|^2 + \alpha(P)$$

donde  $\alpha(P)$  es  $C^\infty(U)$

$$\textcircled{2} \quad dd^c g_D = \psi \quad \text{en } X \setminus \text{supp } D$$

$$\textcircled{3} \quad \int_X g_D \psi = 0$$

Obs  $\wedge \textcircled{1} \wedge \textcircled{2} \Rightarrow \tilde{g}_D - g_D = \alpha$  es suave  
 $\tilde{g}_D \neq g_D$  y  $\Delta \alpha = 0$  en  $X$

$$\Rightarrow \alpha = \text{armónica} \Rightarrow cte. \equiv \alpha.$$

si además  
~~para~~  $\textcircled{3} \quad \alpha \equiv 0.$

obs  $D = \sum [P_i] n_{P_i} \Rightarrow g_D = \sum n_i g_{P_i}$

(linealidad de  $g$  en  $D$ ).

Ejemplo en  $\mathbb{P}^1$

$$g_0(z) = \cancel{A \log |z|^2} - \log \frac{|z|^2}{1+|z|^2} + \text{cte}$$

$$dd^c g_0(z) = \underbrace{-\frac{1}{4\pi} \Delta \log |z|^2}_0 + \underbrace{dd^c (\log(1+|z|^2))}_\psi$$

$$g_w(z) = -\log \frac{|z-w|^2}{(1+|z|^2)} + \text{cte}(w)$$

$$dd^c g_w(z) = \underbrace{-\frac{1}{4\pi} \Delta \log |z-w|^2}_0 + \underbrace{dd^c \log(1+|z|^2)}_{\psi - dd^c \log(1+|w|^2)}$$

CALCULAR CTE !

## Propiedades

Ⓟ  $\{g_P\}_{P \in X}$  green <sup>cr.  $\psi$</sup>   $\Rightarrow g_P(Q) = g_Q(P) \quad \forall P \neq Q \text{ en } X.$

ie  $G : X \times X \setminus \text{Diag} \rightarrow \mathbb{R}$   $\leftarrow$  Funcion green  
 $(P, Q) \mapsto g_P(Q)$   $\leftarrow$  con  $\psi$ .

es simétrica.

obs  $\Rightarrow \bigwedge_{P, Q \in \mathbb{P}^1} g_P(Q) = g_Q(P) \quad \therefore g(z, w) = -\log \frac{|z-w|^2}{(1-z\bar{P})(1-w\bar{P})} + \text{cte}$   
 (indep de  $P, w$ )

Ⓟ  $\psi_1, \psi_2 \in C^\infty$  (resp  $C^\omega$ ) formas de volumen.

$G_1, G_2 : X \times X \setminus \text{Diag} \rightarrow \mathbb{R}$

$\Rightarrow \exists \beta \in C^\infty$  (resp  $C^\omega$ )  $\int g$

$G_1(P, Q) - G_2(P, Q) = \beta(P) + \beta(Q)$

Mas aún  $\beta$  es la única sol de

$d\beta = \psi_1 - \psi_2.$

$\int g \int \beta (\psi_1 + \psi_2) = 0.$

Prop 1

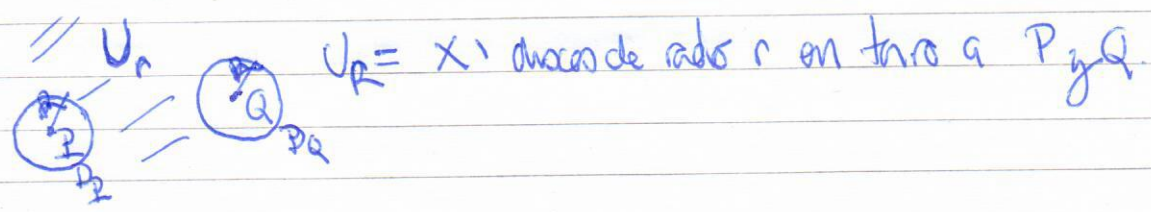
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Lema ① ^ ②

$$\Rightarrow \int_X \overbrace{g_P d^c g_Q - g_Q d^c g_P}^w = g_P(Q) - g_Q(P)$$

Dem



en  $U_r$  
$$d(\underbrace{g_P d^c g_Q - g_Q d^c g_P}_w) = w$$

ya que 
$$d\alpha \wedge d^c \beta = -d^c \alpha \wedge d\beta = d\beta \wedge d^c \alpha$$

Stokes 
$$\int_{\partial U_r} \tilde{w} = \int_{U_r} w = - \int_{\partial P} w - \int_{\partial Q} w$$

$$\int_{\partial P} w \approx \int_{\partial P} (f \log |z-P|^2 + \alpha(z)) \underbrace{d^c g_Q}_{\text{en polares } \frac{1}{2} r \frac{\partial g_Q}{\partial r} \frac{d\theta}{2\pi}} - g_Q (f \log |z-P|^2 + \alpha(P))$$

$$\approx \int_{\partial P} -r \log r d\theta \xrightarrow{0} 0 - g_Q(P) \frac{1}{2} r \left( -\frac{2}{r} + \alpha' \right) \frac{d\theta}{2\pi}$$

$\downarrow$   $\downarrow$   
 $0$   $g_Q(P)$

Prop 2  $\varphi \in C^\infty(\mathbb{C}^n)$

$\Delta$  Thm  $\int_X \varphi = 0 \Rightarrow \exists \rho \in C^\infty(\mathbb{C}^n)$   $dd^c \rho = \varphi$

$G_2(P, Q) + \rho(Q)$  satisfies sing log on  $\mathbb{P}^2$ .

$G_{\mathbb{P}^2}(Q)$   $dd^c G_{\mathbb{P}^2} = \varphi_2 + \varphi_1 - \varphi_2$

$\therefore G_2(P, Q) + \rho(Q) = G_{\mathbb{P}^2}(P, \cdot) + cte$

PD  $\int_X (G_2(P, Q) + \rho(Q)) \varphi_1(Q) = -\rho(P)$   
 $(S_{G_2} \varphi_2 = 0)$

$LI = \int_X G_2(P, Q) (\varphi_1(Q) - \varphi_2(Q)) - \int_X \rho(Q) \varphi_2(Q)$   
 $\underbrace{\hspace{10em}}_{dd^c \rho}$   $\underbrace{\hspace{10em}}_{\text{''}}$   
 $(-S_{\rho} \varphi_2(Q))$   
 $dd^c G_2$

$= \int_X G_2 dd^c \rho - \rho dd^c G_2 = -\int_X d(\rho dd^c G_2 - G_2 dd^c \rho)$

$= \int_{D_2} \rho dd^c G_2 - G_2 dd^c \rho = -\rho(P)$



Teo  $\exists$  ~~una~~ Función de Green  $G$  c/r a cualquier  
 $(L, \Gamma)$  forma  $\psi \in C^\infty$  (resp  $C^\omega$ ).  $G$  es  $C^\infty$  (resp  $C^\omega$ )  
 fuera de la diagonal.

Obs  $G$  invierte  $\Delta$  en sentido de distribuciones.

$$\begin{array}{l} \text{d.d.c.} \\ \text{[G]}: C^\infty(X) \rightarrow \mathbb{R} \\ f \mapsto \int_X G_{\mathbb{R}} dd^c f. \end{array}$$

Similantemente

$$\begin{array}{l} \text{[}\psi\text{]}: C^\infty(X) \rightarrow \mathbb{R} \\ f \mapsto \int_X f \psi. \end{array}$$

Teo ~~d.d.c.~~  $\text{[G]} = \text{[}\psi\text{]} - \delta_{\mathbb{R}}$

ie  $\int_{\mathbb{R}} dd^c f = \int f \psi - f(\mathbb{P}) \quad \forall f \in C^\infty$

Def  $\Delta$  c/r  $\psi$  es por definición

$$(\Delta_\psi f) \psi = -dd^c f.$$

Corolario  $\int_X f \psi = 0 \Rightarrow \int_X G_{\mathbb{R}} (\Delta f) \psi = f(\mathbb{P})$

Función de Green  $g=1$

$$X = \mathbb{C} / \langle 1, \tau \rangle$$

$$z = 0 + iv > 0$$

$$q_\tau = \exp(2\pi i z)$$

$$q_z = \exp(2\pi i z)$$

$$B_2(y) = y^2 - y + 1/6$$

$$P = 0 + \langle 1, \tau \rangle$$

Sea  $\lambda(z) = -\log |q_\tau \cdot (1-qz)| \prod_{n=1}^{\infty} (1-q_\tau^n qz) (1-q_\tau^n / qz)$

Notas  $|e^{2\pi i n z}| = |e^{-2\pi i n v}|^n$

$\therefore |\log(1-q_\tau^n qz)| \leq e^{-2\pi i n v} \therefore$  cv abs.

y  $\lambda(z)$  es  $C^\omega$ .

$$\lambda(z+1) = \lambda(z)$$

$$\lambda(z+\tau) - \lambda(z) = -\log \left| \frac{q_\tau}{qz} \cdot \frac{(1-qz q_\tau)}{(1-qz)} \cdot \frac{(1-qz)}{(1-qz q_\tau)} \right|$$

queremos  $\log |q_\tau(z+\tau)| - \log |q_\tau(z)| = \log |q_\tau|$

$$\log |q_\tau(z+\tau)| - \log |q_\tau(z)| = \log |q_\tau|$$

$$B_2(\tau/v) - B_2(v/v) = \tau/v$$

$$e^{2\pi i (v+\tau v)/v} = e^{2\pi i (v/v) + 2\pi i \tau}$$

$$B_2(y) = y^2 - y + 1/6$$

$$\frac{B_2(y/v+1) - B_2(y/v)}{2} = \frac{x}{v}$$

obs ①  $g(z) = z\lambda(z)$  tiene polo  $\log|z|^2$  cerca de 0 y trasladado.

$$\textcircled{2} \quad dd^c g = -dd^c \log \left| \frac{B_2(y/v)/2}{z} \right|^2$$

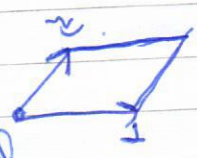
$$= +dd^c \log (2\pi v B_2(y/v))$$

$$= \frac{1}{4\pi} \Delta (2\pi v B_2(y/v))$$

$$= \frac{1}{4\pi} \cdot 2\pi v \cdot \frac{2}{v^2} dx \wedge dy = dx \wedge dy.$$

$$\textcircled{3} \quad \int_{\partial I} \psi = 0. \quad \psi = dx \wedge dy.$$

integraciones  $\lambda(z)$  en dominio fundamental



$\mathbb{Z}^2$

$$I_1 = \int_0^1 \int_0^1 \log |g_{\tau}^{B_2(s)/2}| dr ds.$$

$w = r + sv$   
 $x = r + sv$   
 $y = sv$

$$I_2 = \int_0^1 \int_0^1 \log |1 - g_{\tau}^n g_{r+\tau v}| dr ds$$

$$|g_{\tau}^n g_{r+\tau v}| = e^{-2\pi i v(n+sv)} < 1 \quad \forall n \geq 0$$

$$\therefore -\log(1 - g_{\tau}^n g_{r+\tau v}) = \sum_{m=1}^{\infty} \frac{w^m}{m} \quad | \quad w = g_{\tau}^n g_{r+\tau v}$$

$$y \operatorname{Re}(\log(1 - g_{\tau}^n g_{r+\tau v})) = \log |1 - g_{\tau}^n g_{r+\tau v}|$$

para cada  $s$ 

$$\int_0^1 w^m dr = \int_0^1 g_{\tau}^{nm} \cdot e^{2\pi i m(r+sv)} dr$$

$$= c \int_0^1 e^{2\pi i m r} dr = 0.$$

$$I_3 = \int_0^1 \int_0^1 \log |1 - g_{\tau}^n g_{r-\tau v}| dr ds = 0 \quad \uparrow \text{idem}$$